# **Resource-Efficient Linear Optics Quantum Computation**

via the Cluster State Approach

Dan Browne and Terry Rudolph

# Introduction

- Photons make excellent carriers of quantum information and single qubit operations can be achieved with linear optical elements – (polarising) beam splitters and phase shifters.
- *Two-qubit gates* cannot be achieved deterministically by linear optics alone.
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- Need strong non-linear materials or implement *non-deterministic* gates via photo-detection.
- Schemes<sup>a</sup> for scalable (near)-deterministic gates are complicated and need a large amount of resources (entangled modes, feed-forward detection) to implement even the simplest gate.
- Here we describe a new scheme which employs the measurement-based *cluster state quantum computation* approach and achieves significant gains in resource efficiency.
- <sup>*a*</sup> E.g. Knill-Laflamme-Milburn (KLM), Nature (2001).

#### **Cluster States**

A cluster state<sup>a</sup> is an entangled multi-qubit state which may be represented by a graph.



- Vertices represent qubits prepared in state<sup>b</sup>  $|+\rangle = |0\rangle + |1\rangle$ .
- Edges represent the application of the entangling quantum CPHASE gate

 $|0\rangle^a \langle 0|\mathbb{1}^b + |1\rangle^a \langle 1|\sigma_z^b$ 

between the connected qubits.

- We will refer to the graph edges as **bonds**.
- Known extra Pauli's on any cluster qubit can be accounted for.

<sup>&</sup>lt;sup>a</sup>Briegel and Raussendorf, PRL, **86**, 910 (2001)

<sup>&</sup>lt;sup>b</sup>Normalisation factors will be omitted.

#### Some Properties of Cluster States

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- Redundant encoding is to encode a logical qubit in several qubits. The 2-qubit redundant encoding is  $|0\rangle_{log.} \equiv |0\rangle|0\rangle$ ,  $|1\rangle_{log.} \equiv |1\rangle|1\rangle$ .
- A  $\sigma_x$  measurement  $\{(|0\rangle + |1\rangle), (|0\rangle |1\rangle)\}$  on a qubit in a *linear* cluster combines the neighbouring qubits into a *single logical qubit* (redundantly encoded).



Single logical cluster qubit encoded redundantly in two physical qubits

## **Cluster State Quantum Computation**

On a cluster state with sufficient size and bond layout, an arbitrary quantum network can be simulated by *adaptive single-qubit measurements* alone<sup>a</sup>. The following cluster state layout<sup>b</sup>



<sup>a</sup>Raussendorf and Briegel, PRL 86, 5188; Raussendorf, Browne and Briegel, PRA 68, 022312

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with the above measurements, simulates the quantum network:



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#### Outline of our scheme

- Qubits will be single photon polarisation  $|0\rangle \equiv |H\rangle$ ,  $|1\rangle \equiv |V\rangle$ .
- Polarisation measurement in arbitrary bases is trivial. The main part of our scheme is the cluster state generation.
- Instead of using CPHASE gates between qubits, we (probabilistically) *fuse* clusters<sup>a</sup>.

<sup>a</sup>The same idea underlies the "valence bond model" of Verstraete and Cirac, quant-ph/0311130 <sup>b</sup>Can be generated e.g. via non-linear processes or linear optics and feed-forward.

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- Let us introduce *fusion* operator  $|0\rangle\langle00| + |1\rangle\langle11|$ .
- This replaces two qubits with a single one while *retaining* all cluster state bonds on each qubit.



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- The initial resource will be **photon Bellpairs**<sup>b</sup>  $|H\rangle(|H\rangle + |V\rangle) + |V\rangle(|H\rangle - |V\rangle)$ . These are 2-qubit cluster states.



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## The fusion operation

• The fusion can be seen explicitly if we write out each bond CPHASE:



which, after the fusion operation  $|0\rangle\langle00| + |1\rangle\langle11|$ , becomes



$$|0\rangle \prod_{i=1}^{m} \mathbb{1}^{(i)} |\psi\rangle \prod_{i=1}^{n} \mathbb{1}^{(i)} |\psi'\rangle + |1\rangle \prod_{i=1}^{m} \sigma_z^{(i)} |\psi\rangle \prod_{i=1}^{n} \sigma_z^{(i)} |\psi'\rangle$$

Polarising Beam Splitter - (PBS)

The key component for realising the fusion operation is the PBS.



#### **Building Linear Clusters - Type-I Fusion**

- The fusion operation can be realised *nondeterministically* using the illustrated setup:<sup>*a*</sup>
- With a photon incident in each port, there are 4 possible outcomes, each with probability 25%.



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• **Two** outcomes give us the desired fusion operators  $|0\rangle\langle00| + |1\rangle\langle11|$  or  $|0\rangle\langle00| - |1\rangle\langle11|$ , (one and only one photon, H or V). The second of these adds an extra  $\sigma_z$  but this is naturally accounted for. Thus, overall success probability is 50%.

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- If 0 or 2 photons are detected, this is a *failure*, equivalent to measuring both qubits in the  $(\sigma_z)$  computational basis. The qubits are thus both cut from their respective clusters.

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#### **Linear Clusters**



• On average, the length of the cluster does not increase<sup>a</sup>.

<sup>a</sup>Actually due to the "reflective boundary" at length 2, one reaches the desired length in quadratic steps.

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• On average, cluster length increases by 1/2 qubit. Therefore, to add one qubit to the cluster you need  $2 \times 4 - 1 = 7$  Bell Pairs.

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- The best protocol we have found uses 5-photon clusters as building blocks. This gives a rate: 6.5 Bell pairs per added qubit.

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# Joining Clusters into the 2-D Pattern

 We now need to join these linear clusters into the desired 3-dimensional layout.

• With a deterministic fusion, this pattern of fusions produces the cluster state layout required.



• However, our Type-I fusion is only successful half the time. Failure is equivalent to  $\sigma_z$  measurement, which would break up hard won existing bonds!

- Recall that a  $\sigma_x$  measurement on a cluster state does not cut the qubits bonds, but merges neighbouring qubits into a single redundantly encoded qubit.
- If we modify our fusion operation by introducing extra 45° rotations to each qubit, the failure outcomes will be  $\sigma_x$  measurements.
- However, the "success" projection is then no longer diagonal in the computational basis and does not perform the required fusion.



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- If we modify our fusion operation by introducing extra 45° rotations to each qubit, the failure outcomes will be  $\sigma_x$  measurements.
- However, the "success" projection is then no longer diagonal in the computational basis and does not perform the required fusion.
- We get round this by measuring *both* outputs.
- This leads to projections onto states  $|++\rangle+|--\rangle = |00\rangle+|11\rangle$ or  $|++\rangle-|--\rangle = |01\rangle+|10\rangle$ .



- If one of the qubits this is applied to is *redundantly encoded* this gives us the desired fusion! We call this a Type-II fusion.
- Again, the success probability of this step is 50%.

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Apply a  $\sigma_x$  measurement to prepare redundantly encoded qubit



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- Note that failures only "use up" cluster qubits to the right of the fusion.
- Thus, failures *cannot* propagate back through the cluster as in the Type-I fusion and other schemes.

# Quantifying Resource Requirements

- In our cluster state measurement pattern, there are the same number of:
  - simulated 2-qubit gates:



T-shaped units in the cluster state:



This means the resources required to build the T-shape are a measure of the resources per two-qubit gate.

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- If we use the method above, the construction of a T-shape consumes on average 8 bonds from the linear clusters used.
- Thus, per general 2-qubit gate the resource requirements are:

 $8 \times 6.5 = 52$  Bell Pairs.

# Other Schemes: Rough Comparison of Resources

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Approx. entanglement resources required per (general) 2-qubit gate:

Knill-Laflamme-Milburn (KLM), Nature 401, 46 (2001)

**Yoran and Reznik**, PRL 91, 037903 (2003). (Measurement based "chain state" q.c., uses KLM gates)

**Nielsen**, accepted PRL (2004) (Cluster state scheme using KLM gates)

Our scheme, quant-ph/0405157

(for 92.5% gate success prob.)  $\sim$ 100-photon "KLM state"<sup>a</sup>

 $\sim$ 23 12-photon "KLM states"

 ${\sim}54$  8-photon "KLM states"

52 2-photon Bell states

<sup>*a*</sup>Note that the KLM resource states require a complicated linear optical network conditional on several / many measurements for their generation.

# Summary

- We have presented a scheme for *linear optics quantum computation* based on the cluster state approach, that is very resource efficient compared to other schemes.
- For the shorter-term, the work provides a recipe for the generation of interesting new entangled states.
- The procedures at the heart of the scheme have already been implemented experimentally.
- The scheme has other advantages. For example, the absence of concatenated beam-splitters, unavoidable in other schemes, makes the mode-matching requirements much less strict.

# **Future Directions**

- Cluster state layout can be optimised for specific algorithms how much more gains in resource efficiency are possible?
- Can the general scheme be optimised to further reduce its experimental complexity?
- What about fault-tolerance?

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This work can be found in pre-print quant-ph/0405157.

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