

# Distillation of Qubits through Zeno-like Measurements

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## References:

- H. Nakazato, T. Takazawa, and K. Y., Phys. Rev. Lett. **90**, 060401 (2003);
- H. Nakazato, M. Unoki, and K. Y., Phys. Rev. A **70**, 012303 (2004);
- G. Compagno *et al.*, quant-ph/0405074 (2004).
- L.-A. Wu, D. A. Lidar, and S. Schneider, quant-ph/0402209 (2004).

## Introduction/Motivation

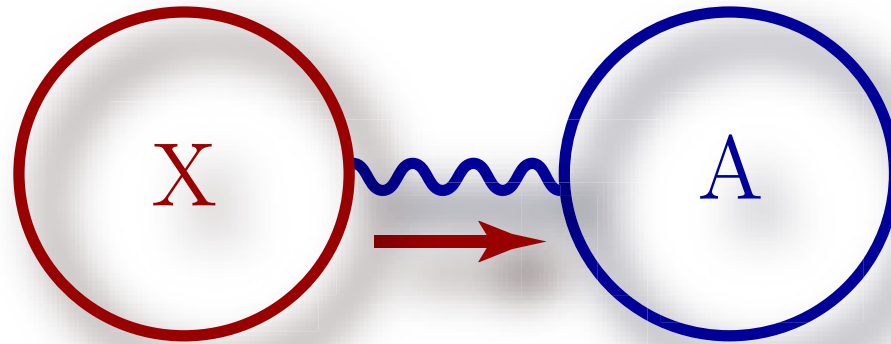
- How to prepare a quantum state  
*with high quantum coherence (entangled state, ...)*
- Purification/Distillation
  - ... Extraction of a pure state *from an arbitrary mixed state*
  - ▷ in case a state preparation *via a direct projective measurement*  
is not possible

## Abstract

- A novel method of purification/distillation  
*based on “Zeno-like measurements”*
  - ▷ general scheme   ▷ mechanism   ▷ optimization
- Possible Applications to Qubit-Systems
  - ▷ *Initialization of Multiple Qubits*    $\rho \rightarrow |\downarrow\downarrow\dots\rangle\langle\downarrow\downarrow\dots|$
  - ▷ *Entanglement Distillation*    $\rho \rightarrow |\Psi^-\rangle\langle\Psi^-|$
  - ▷ *Distillation of Entanglement*  
*between Spatially Separated Qubits*

# Purification through Zeno-like Measurements

H. Nakazato, T. Takazawa, and K. Y., Phys. Rev. Lett. **90**, 060401 (2003);  
H. Nakazato, M. Unoki, and K. Y., Phys. Rev. A **70**, 012303 (2004).



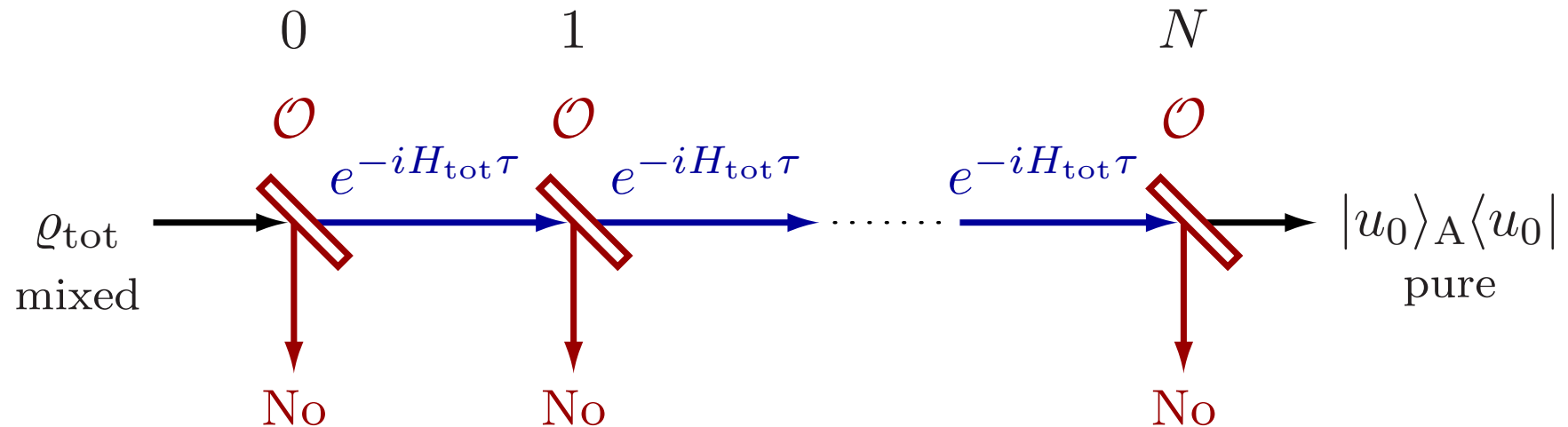
Repeated Measurements on X  
(Zeno-like Measurements)



Extraction of a Pure State in A

$$\rho_A \xrightarrow{\text{mixed}} |u_0\rangle_A \langle u_0| \text{ pure}$$

## Zeno-like Measurements on X



We retain only those events where  $|\phi\rangle_X$  is found at *every* measurement.

▷ projection op.  $\mathcal{O} = |\phi\rangle_X \langle \phi|$

Parameters:

$\tau$  ... interval between measurements

$|\phi\rangle_X$  ... measuring state

$H_{\text{tot}}$  ... parameters in the Hamiltonian

(coupling constants  $g, \dots$ )

## Evolution under Zeno-like Measurements

State after  $N$  Successful Confirmations

$$\begin{aligned}\varrho_{\text{tot}}^{(\tau)}(N) &= (\mathcal{O}e^{-iH_{\text{tot}}\tau})^N \mathcal{O} \varrho_{\text{tot}} \mathcal{O} (e^{iH_{\text{tot}}\tau} \mathcal{O})^N / P^{(\tau)}(N) \\ &= |\phi\rangle_{\text{X}} \langle \phi| \otimes \varrho_{\text{A}}^{(\tau)}(N)\end{aligned}$$

$$\varrho_{\text{A}}^{(\tau)}(N) = (V_{\phi}(\tau))^N \tilde{\varrho}_{\text{A}} (V_{\phi}^{\dagger}(\tau))^N / P^{(\tau)}(N) \xrightarrow{N \text{ increases}} ?$$

Success Probability

$$P^{(\tau)}(N) = \text{Tr}_{\text{A}} [(V_{\phi}(\tau))^N \tilde{\varrho}_{\text{A}} (V_{\phi}^{\dagger}(\tau))^N]$$

$$V_{\phi}(\tau) = {}_{\text{X}}\langle \phi| e^{-iH_{\text{tot}}\tau} |\phi\rangle_{\text{X}}, \quad \tilde{\varrho}_{\text{A}} = {}_{\text{X}}\langle \phi| \varrho_{\text{tot}} |\phi\rangle_{\text{X}} \cdots \text{ operators in } \mathcal{H}_{\text{A}}$$

## Purification, Conditions & Optimization

$$(V_\phi(\tau))^N = \sum_n \lambda_n^N |u_n\rangle_A \langle v_n| \xrightarrow{N \text{ increases}} \lambda_0^N |u_0\rangle_A \langle v_0| \quad (0 \leq |\lambda_n| \leq 1)$$

- $\lambda_0 \dots$  largest (in magnitude) eigenvalue (*unique, discrete, and nondegenerate*)

$$\varrho_A^{(\tau)}(N) = (V_\phi(\tau))^N \tilde{\varrho}_A (V_\phi^\dagger(\tau))^N / P^{(\tau)}(N) \xrightarrow{N \text{ increases}} |u_0\rangle_A \langle u_0| \dots \text{Pure State!}$$

- $|u_0\rangle_A$  is the right eigenvector of  $V_\phi(\tau) = \text{X} \langle \phi | e^{-iH_{\text{tot}}\tau} | \phi \rangle_{\text{X}}$  belonging to  $\lambda_0$ .
  - $|u_0\rangle_A$  is irrespective of the initial mixed state  $\varrho_{\text{tot}}$ .

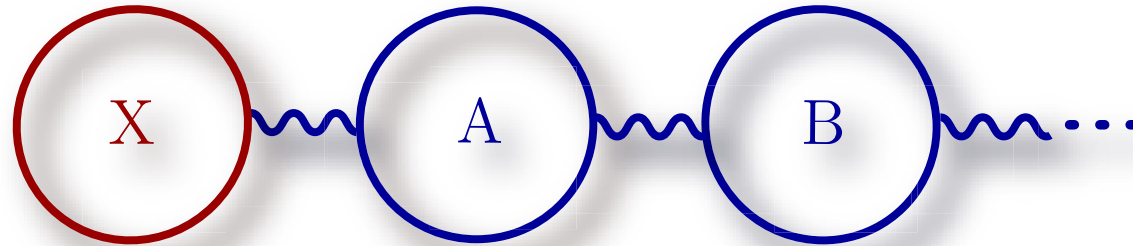
$$P^{(\tau)}(N) \sim |\lambda_0|^{2N} \text{XA} \langle \phi v_0 | \varrho_{\text{tot}} | \phi v_0 \rangle_{\text{XA}} \quad (0 \leq |\lambda_n| \leq 1)$$

### Optimization

- $|\lambda_0| = 1 \quad \rightarrow$  suppression of the decay of  $P^{(\tau)}(N)$
- $|\lambda_n/\lambda_0| \ll 1 \quad \rightarrow$  faster purification

- adjust  $\tau, |\phi\rangle_{\text{X}}, H_{\text{tot}}$

## Possible Applications to Qubit-Systems



$$H_0 = \Omega \frac{1 + \sigma_3^X}{2} + \Omega \frac{1 + \sigma_3^A}{2} + \Omega \frac{1 + \sigma_3^B}{2} + \dots$$

$$H'_{XA} = g(\sigma_+^X \sigma_-^A + \sigma_-^X \sigma_+^A), \quad H'_{AB} = g(\sigma_+^A \sigma_-^B + \sigma_-^A \sigma_+^B), \quad \dots$$

Fidelity

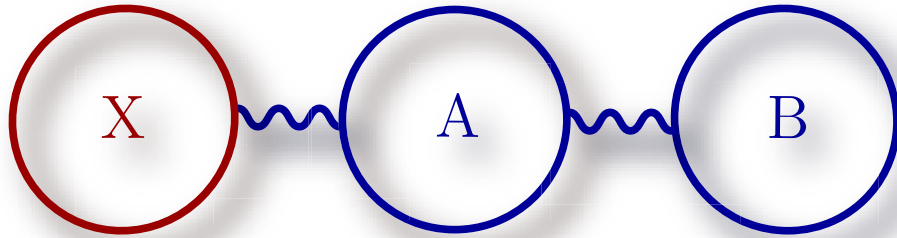
$$F^{(\tau)}(N) \equiv \text{tr}_A \langle u_0 | \rho_A^{(\tau)}(N) | u_0 \rangle_A$$

Success Probability

$$P^{(\tau)}(N) \sim |\lambda_0|^{2N} \text{tr}_{XA} \langle \phi v_0 | \rho_{\text{tot}} | \phi v_0 \rangle_{XA}$$



# Initialization of Qubits



Repeated Confirmations of  $|\downarrow\rangle_X$   
(Zeno-like Measurements)



Initialization of Qubits A and B

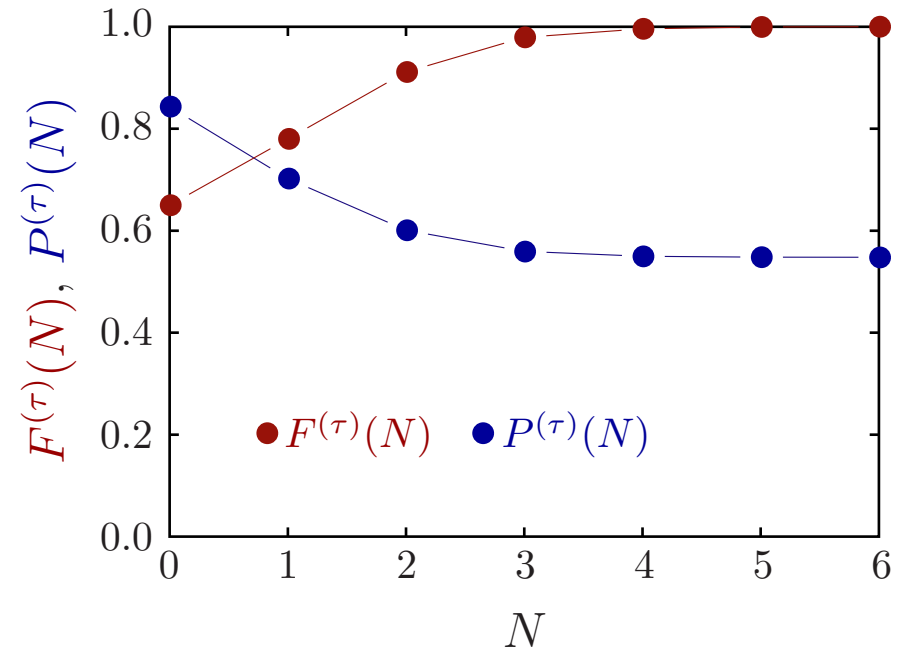
$$\rho_{AB} \xrightarrow{\text{mixed}} |\downarrow\downarrow\rangle_{AB}\langle\downarrow\downarrow| \text{ initialized}$$

## Tuning

1.  $|\phi\rangle_X = |\downarrow\rangle_X$   
 $\Leftarrow$  for  $|\lambda_0| = 1$
2.  $\sqrt{2}g\tau \neq n\pi$  ( $n = 1, 2, \dots$ )  
 $\Leftarrow$  for the uniqueness of  $\lambda_0$
3.  $\sqrt{2}g\tau = 2n\pi \pm \zeta$ ,  $\zeta = \tan^{-1}(2\sqrt{2})$   
 $\Leftarrow$  for a quick initialization

Initialization with the “optimal” prob.

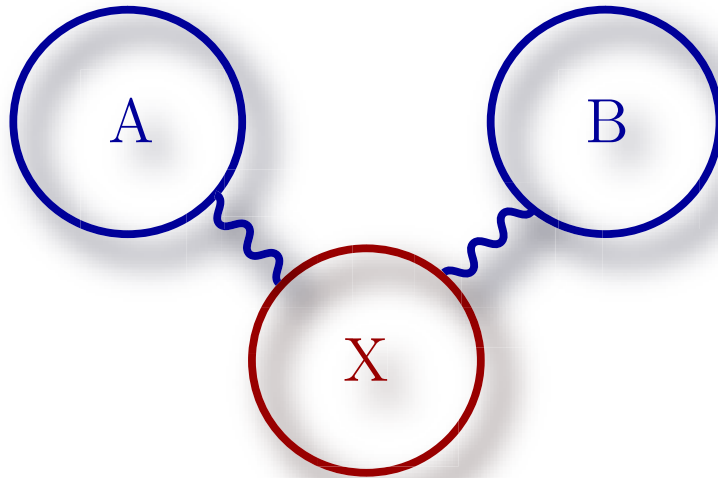
$$X_{AB}\langle\downarrow\downarrow\downarrow|\rho_{\text{tot}}|\downarrow\downarrow\downarrow\rangle_{XAB}$$



$$e^{-\beta H_{\text{tot}}} \rightarrow |\downarrow\downarrow\downarrow\rangle_{XAB}\langle\downarrow\downarrow\downarrow|$$

$\Omega = 2, \quad g = 1, \quad \tau \simeq 1.73, \quad \beta = 1$

# Entanglement Purification



Repeated Measurements on X  
(Zeno-like Measurements)



Entanglement between A and B

$$\rho_{AB} \longrightarrow |\Psi^-\rangle_{AB}\langle\Psi^-|$$

mixed                      Bell state

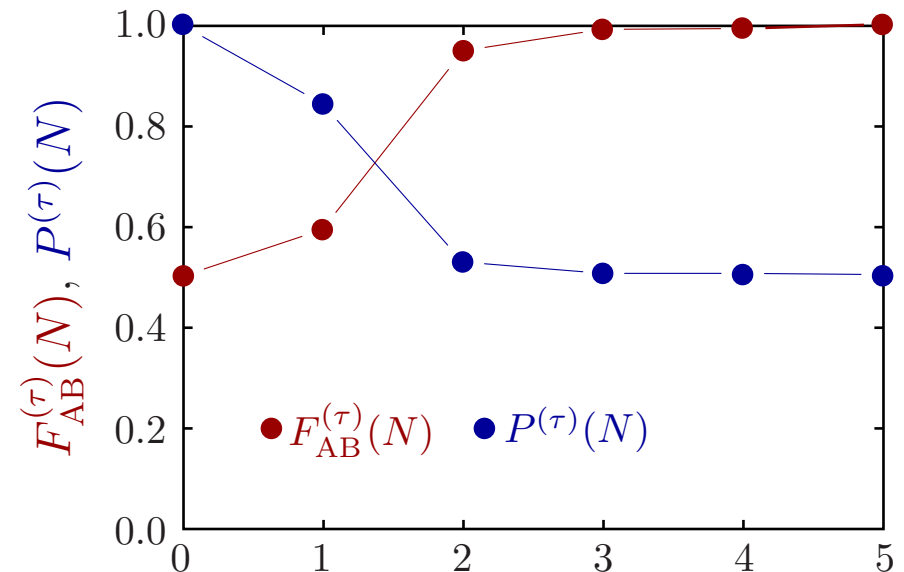
$$|\Psi^-\rangle_{AB} = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle_{AB} - |\downarrow\uparrow\rangle_{AB})$$

## Tuning

1.  $|\Omega|\tau = 2n\pi$  ( $n = 0, 1, \dots$ )  
 $\Leftarrow$  for  $|\lambda_{\Psi^-}| = 1$
2.  $|\phi\rangle_X \neq |\uparrow\rangle_X, |\downarrow\rangle_X$  and  $\sqrt{2}g\tau \neq n\pi$   
 $\Leftarrow$  for the uniqueness of  $\lambda_0$
3.  $\sqrt{2}g\tau = 2n\pi \pm \zeta$ ,  $\zeta = \tan^{-1}(2\sqrt{2})$   
 $\Leftarrow$  for a quick purification

with the “optimal” prob.

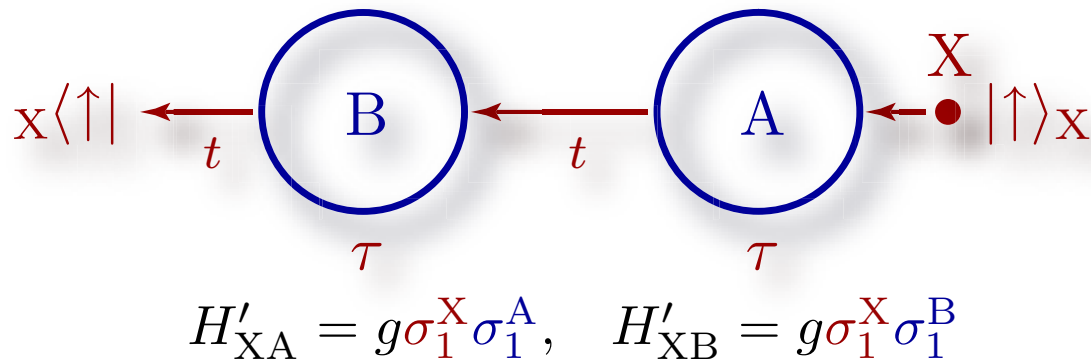
$${}_{XAB}\langle\phi\Psi^-|\rho_{\text{tot}}|\phi\Psi^-\rangle_{XAB}$$



$$|\rightarrow\rangle_X \otimes |\uparrow\rangle_A \otimes |\downarrow\rangle_B \rightarrow |\rightarrow\rangle_X \otimes |\Psi^-\rangle_{AB}$$

$$\Omega = 0, \quad g\tau \simeq 0.5$$

# Entanglement between *Separated* Qubits I



Repeatedly Throwing X



Entanglement between A and B

$\rho_{AB} \longrightarrow |\Psi\rangle_{AB}\langle\Psi|$   
 mixed                      entangled state

$$|\Psi\rangle_{AB} = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle_{AB} - e^{i\chi}|\downarrow\uparrow\rangle_{AB})$$

Projected Time-Evolution Op.

$$V_{\uparrow\uparrow} = \mathbf{x}\langle\uparrow| e^{-iH_0 t} e^{-i(H_0 + H'_{XB})\tau} e^{-iH_0 t} e^{-i(H_0 + H'_{XA})\tau} |\uparrow\rangle_{\mathbf{X}}$$

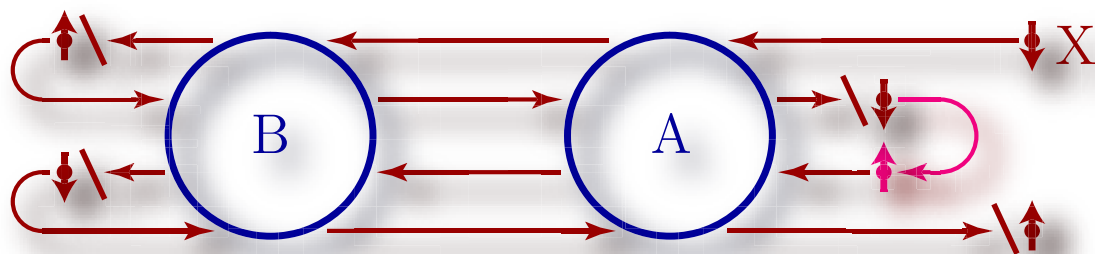
Tuning

$$\cos\sqrt{\Omega^2 + g^2}\tau - i\frac{\Omega}{\sqrt{\Omega^2 + g^2}}\sin\sqrt{\Omega^2 + g^2}\tau = -e^{i\Omega t}\cos g\tau$$

for  $|\lambda_\Psi| = 1$  and the uniqueness of  $\lambda_0$

cf., A. Messina, Eur. Phys. J. D **18**, 379 (2002) and  
 D. E. Browne and M. B. Plenio, Phys. Rev. A **67**, 012325 (2003),  
 where the initial states  $\rho_{tot}$  should be carefully prepared.

## Entanglement between *Separated* Qubits II



$$H'_{XA} = g_A(\sigma_+^X \sigma_-^A + \sigma_-^X \sigma_+^A),$$

$$H'_{XB} = g_B(\sigma_+^X \sigma_-^B + \sigma_-^X \sigma_+^B)$$

Throwing X, 4 Times



Entanglement between A and B

$$\rho_{AB} \text{ mixed} \longrightarrow |\Psi^\pm\rangle_{AB} \langle \Psi^\pm| \text{ Bell state}$$

$$|\Psi^\pm\rangle_{AB} = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle_{AB} \pm |\downarrow\uparrow\rangle_{AB})$$

Projected Time-Evolution Op.

$$V_c = V_{\uparrow\downarrow\uparrow} V_{\downarrow\uparrow\downarrow}$$

$$V_{\downarrow\uparrow\downarrow} = \mathbf{x} \langle \downarrow | U_I^{(A)}(\tau_A) U_I^{(B)}(\tau_B) | \uparrow \rangle_{\mathbf{x}} \langle \uparrow | U_I^{(B)}(\tau_B) U_I^{(A)}(\tau_A) | \downarrow \rangle_{\mathbf{x}}$$

$$V_{\uparrow\downarrow\uparrow} = \mathbf{x} \langle \uparrow | U_I^{(A)}(\tau_A) U_I^{(B)}(\tau_B) | \downarrow \rangle_{\mathbf{x}} \langle \downarrow | U_I^{(B)}(\tau_B) U_I^{(A)}(\tau_A) | \uparrow \rangle_{\mathbf{x}}$$

— Tuning —

$$\cos g_A \tau_A = 0 \quad \text{and} \quad \sin g_B \tau_B = \pm \frac{1}{\sqrt{2}}$$

for  $|\lambda_{\Psi^-}| = 1$  and the uniqueness of  $\lambda_0$

- Success Prob.  $P = {}_{AB} \langle \Psi^\pm | \rho_{AB} | \Psi^\pm \rangle_{AB} \cdots$  “optimal”

## Summary

- Simple (just repeat one and the same measurement).  
( $\Leftrightarrow$  combination of rotation, CNOT operation, measurement, ...)
- High fidelity with high success probability after a few steps.
  - ▷  $|\lambda_0| = 1, \quad |\lambda_1| \ll 1$
- “Optimal” success probability  $P^{(\tau)}(N) \rightarrow {}_{XA}\langle\phi\Psi|\varrho_{\text{tot}}|\phi\Psi\rangle_{XA}$  is possible.
  - ▷ The target state  $|\phi\Psi\rangle_{XA}$  contained in the initial state  $\varrho_{\text{tot}}$  is fully extracted.  
( $\Leftrightarrow$  decays to zero)
- General & Flexible
  - ▷ It accepts various experimental settings.
  - ▷ Many potential generalizations.  
( $\Leftrightarrow$  special design for a specific system)
- Robustness ?
  - ▷ Precise tuning of parameters ?
  - ▷ Errors in Measurements ?
- Experimentally feasible setups ?
  - ▷ Instead of repeating measurements;  
Continuous application of an external field.
  - ▷ “One-photon” scheme  $\Rightarrow$  “Multi-photon” scheme, using coherent light