Discontinuous Collocation Methods for DAEs in Mechanics

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Differential-Algebraic Equations

We consider overdetermined mixed index 2 and 3 DAEs:

$$\dot{y} = v(t, y, z)$$

$$\dot{z} = f(t, y, z, \psi) + r(t, y, \lambda)$$

$$0 = g(t, y)$$

$$0 = g_t(t, y) + g_y(t, y)v(t, y, z)$$

$$0 = k(t, y, z)$$

The functions

$$\begin{aligned} \mathbf{v} &: \mathbb{R} \times \mathbb{R}^{n_y} \times \mathbb{R}^{n_z} \to \mathbb{R}^{n_y} \\ f &: \mathbb{R} \times \mathbb{R}^{n_y} \times \mathbb{R}^{n_z} \times \mathbb{R}^{n_k} \to \mathbb{R}^{n_z} \\ r &: \mathbb{R} \times \mathbb{R}^{n_y} \times \mathbb{R}^{n_g} \to \mathbb{R}^{n_z} \\ g &: \mathbb{R} \times \mathbb{R}^{n_y} \to \mathbb{R}^{n_g} \\ k &: \mathbb{R} \times \mathbb{R}^{n_y} \times \mathbb{R}^{n_z} \to \mathbb{R}^{n_k} \end{aligned}$$

The two matrices

$$g_{y}(t,y)v_{z}(t,y,z)r_{\lambda}(t,y,\lambda)$$

$$\begin{bmatrix} g_{y}v_{z}(t,y,z)\\ k_{z}(t,y,z) \end{bmatrix} \begin{bmatrix} r_{\lambda}(t,y,\lambda) & f_{\psi}(t,y,z) \end{bmatrix}$$

are assumed invertible. These assumptions allow for the system of DAEs to be expressed as ODEs. These are the so called *underlying ODEs*.

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Well-known formulations of classical mechanics are *Hamiltonian Mechanics* and *Lagrangian Mechanics*. If the evolution of the system is constrained, the result is a DAE.

In the context of classical mechanics,

- Holonomic constraints are restrictions on the coordinates of a system. For example, the position of a pendulum in Cartesian coordinates is constrained to a circle or sphere.
- Nonholonomic constraints are (nonintegrable) restrictions on the velocities of a system. For example, an ice skate is constrained to move in the direction the blade is pointing.

Examples from Mechanics

Lagrangian System

$$\dot{q} = v$$

$$\frac{d}{dt} \nabla_{v} L(t, q, v) = \nabla_{q} L(t, q, v) - g_{q}(t, q)^{T} \lambda - k_{v}(t, q, v)^{T} \psi$$

$$0 = g(t, q)$$

$$0 = g_{t}(t, q) + g_{q}(t, q)v$$

$$0 = k(t, q, v)$$

Hamiltonian System

$$\begin{split} \dot{q} &= \nabla_{p} H(t, q, p) \\ \dot{p} &= -\nabla_{q} H(t, q, p) - g_{q}(t, q)^{T} \lambda - k_{p}(t, q, p)^{T} \psi \\ 0 &= g(t, q) \\ 0 &= g_{t}(t, q) + g_{q}(t, q) \nabla_{p} H(t, q, p) \\ 0 &= k(t, q, p) \end{split}$$

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We introduce two classes of Runge-Kutta type methods for solving overdetermined mixed index 2 and 3 DAEs.

- Specialized partitioned additive Runge-Kutta (SPARK) methods were introduced by Jay in 1998, and can be applied to systems with nonholonomic constraints. Methods were developed in 2007 and can be applied to systems with holonomic constraints.
- Murua's partitioned Runge-Kutta methods were proposed by Murua in 1996 for index 2 DAEs, and can be applied to systems with nonholonomic constraints.

We consider in this presentation an extension of the SPARK methods for systems of mixed index 2 and 3 DAEs. We consider also the *extended Murua's partitioned Runge-Kutta (EMPRK) methods* to systems of mixed index 2 and 3 DAEs.

(1,1)-SPARK/EMPRK Midpoint-Trapezoidal Method

$$Y_{1} = y_{0} + \frac{h}{2}v(t_{0} + \frac{1}{2}h, Y_{1}, Z_{1})$$

$$Z_{1} = z_{0} + \frac{h}{2}f(t_{0} + \frac{1}{2}h, Y_{1}, Z_{1}, \Psi_{1}) + \frac{h}{2}r(t_{0}, y_{0}, \Lambda_{0})$$

$$y_{1} = y_{0} + h v(t_{0} + \frac{1}{2}h, Y_{1}, Z_{1})$$

$$z_{1} = z_{0} + h f(t_{0} + \frac{1}{2}h, Y_{1}, Z_{1}, \Psi_{1}) + \frac{h}{2}r(t_{0}, y_{0}, \Lambda_{0}) + \frac{h}{2}r(t_{1}, y_{1}, \Lambda_{1})$$

$$0 = g(t_{1}, y_{1})$$

$$0 = g_{t}(t_{1}, y_{1}) + g_{y}(t_{1}, y_{1})v(t_{1}, y_{1}, z_{1})$$

$$0 = k(t_{1}, y_{1}, z_{1})$$

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The (s, s)-SPARK/EMPRK Methods

The internal stages

$$\begin{split} Y_{i} &= y_{0} + h \sum_{j=1}^{s} a_{ij} v(t_{0} + c_{j}h, Y_{j}, Z_{j}), \quad i = 1, \dots, s \\ Z_{i} &= z_{0} + h \sum_{j=1}^{s} a_{ij} f(t_{0} + c_{j}h, Y_{j}, Z_{j}, \Psi_{j}) + h \sum_{j=0}^{s} \widetilde{a}_{ij} r(t_{0} + \widetilde{c}_{j}h, \widetilde{Y}_{j}, \Lambda_{j}), \\ i &= 1, \dots, s \\ \widetilde{Y}_{i} &= y_{0} + h \sum_{j=1}^{s} \overline{a}_{ij} v(t_{0} + c_{j}h, Y_{j}, Z_{j}), \quad i = 0, \dots, s \\ \widetilde{Z}_{i} &= z_{0} + h \sum_{j=1}^{s} \overline{a}_{ij} f(t_{0} + c_{j}h, Y_{j}, Z_{j}, \Psi_{j}) + h \sum_{j=0}^{s} \breve{a}_{ij} r(t_{0} + \widetilde{c}_{j}h, \widetilde{Y}_{j}, \Lambda_{j}), \end{split}$$

 $i = 0, \ldots, s$

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The constraints

$$0 = g(t_0 + \tilde{c}_i h, \tilde{Y}_i), \quad i = 0, ..., s$$

$$0 = g(t_1, y_1)$$

$$0 = g_t(t_1, y_1) + g_y(t_1, y_1)v(t_1, y_1, z_1)$$

$$0 = \sum_{j=1}^{s} b_j c_j^{i-1} k(t_0 + c_j h, Y_j, Z_j), \quad i = 1, ..., s - 1$$

$$0 = k(t_0 + \tilde{c}_i h, \tilde{Y}_i, \tilde{Z}_i), \quad i = 0, ..., s$$

$$0 = k(t_1, y_1, z_1)$$

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The numerical solution

$$y_{1} = y_{0} + h \sum_{j=1}^{s} b_{j} v(t_{0} + c_{j}h, Y_{j}, Z_{j})$$

$$z_{1} = z_{0} + h \sum_{j=1}^{s} b_{j} f(t_{0} + c_{j}h, Y_{j}, Z_{j}, \Psi_{j}) + h \sum_{j=0}^{s} \widetilde{b}_{j} r(t_{0} + \widetilde{c}_{j}h, \widetilde{Y}_{j}, \Lambda_{j})$$

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An example of a class of methods is the (s, s)-Gauss-Lobatto methods.

- Use *s*-stage Gauss coefficients for the a_{ij} , b_i , c_i
- Use (s + 1)-stage Lobatto coefficients for the \widetilde{b}_i , \widetilde{c}_i
- Use (s + 1)-stage Lobatto-IIIA for \check{a}_{ij}
- Define *a_{ij}* and *a_{ij}* by

$$\sum_{j=1}^{s} \overline{a}_{ij} c_j^{k-1} = \frac{\widetilde{c}_i^k}{k}, \quad k = 1, \dots, s$$
$$\overline{a}_{0j} = 0, \quad j = 1, \dots, s$$
$$\sum_{j=0}^{s} \widetilde{a}_{ij} \widetilde{c}_j^{k-1} = \frac{c_i^k}{k}, \quad k = 1, \dots, s$$
$$\widetilde{a}_{i0} = \widetilde{b}_0, \quad i = 1, \dots, s.$$

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Theorem

Suppose that $y_0 = y_0(h)$, $z_0 = z_0(h)$, $\lambda_0 = \lambda_0(h)$, $\psi_0 = \psi_0(h)$ satisfy

$$\begin{split} o(h^2) &= g(t_0, y_0) \\ o(h) &= \frac{d}{dt}(g(t, y))(t_0, y_0, z_0) \\ o(1) &= \frac{d^2}{dt^2}(g(t, y))(t_0, y_0, z_0, \lambda_0, \psi_0) \\ o(h) &= k(t_0, y_0, z_0) \\ o(1) &= \frac{d}{dt}(k(t, y, z))(t_0, y_0, z_0, \lambda_0, \psi_0) \end{split}$$

Then the SPARK and EMPRK methods possess a unique solution for h sufficiently small.

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- Local order is defined through series expansions of the numerical and exact solutions. However, using series expansions to determine local order is tedious.
- Collocation methods (and discontinuous collocation methods), by contrast, have a much cleaner derivation of their local order.

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We are thus interested in showing the equivalence of the SPARK and EMPRK methods to a class of discontinuous collocation type methods.

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Collocation and Discontinuous Collocation



Hairer, Nørsett, Wanner. Solving Ordinary Differential Equations I, Nonstiff Problems. 2000

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Discontinuous Collocation Type Methods

Let c_1, \ldots, c_s be distinct real numbers, and $\tilde{c}_0, \ldots, \tilde{c}_s$ also be distinct real numbers, with $\tilde{c}_0 = 0$ and $\tilde{c}_s = 1$. Assume also that \tilde{b}_0 , \tilde{b}_s are positive real numbers. We then define the *s*-degree polynomials Y(t), $\Lambda(t)$, $\Psi(t)$, $Z^f(t)$, the (s + 1)-degree polynomials Z(t), $Z^r(t)$, and the (s + 2)-degree polynomials $\tilde{Z}(t)$, $\tilde{Z}^r(t)$ as the polynomials satisfying the initial conditions

$$\begin{split} Y(t_0) &= y_0, \\ Z^f(t_0) &= z_0, \quad Z^r(t_0) = -h \widetilde{b}_0 \widetilde{\mu}(t_0), \\ \widetilde{Z}^r(t_0) &= 0, \\ Z(t_0) &= Z^f(t_0) + Z^r(t_0) = z_0 - h \widetilde{b}_0 \widetilde{\mu}(t_0) \\ \widetilde{Z}(t_0) &= Z^f(t_0) + \widetilde{Z}^r(t_0) = z_0, \end{split}$$

where

$$\widetilde{\mu}(t) := \dot{Z}^{r}(t) - r(t, Y(t), \Lambda(t)),$$

and the conditions

$$\begin{split} \dot{Y}(t_{0}+c_{i}h) &= v(t_{0}+c_{i}h,Y(t_{0}+c_{i}h),Z(t_{0}+c_{i}h)), \quad i=1,\ldots,s \\ \dot{Z}^{f}(t_{0}+c_{i}h) &= f(t_{0}+c_{i}h,Y(t_{0}+c_{i}h),Z(t_{0}+c_{i}h),\Psi(t_{0}+c_{i}h)), \\ \quad i=1,\ldots,s \\ \dot{Z}^{r}(t_{0}+\widetilde{c}_{i}h) &= r(t_{0}+\widetilde{c}_{i}h,Y(t_{0}+\widetilde{c}_{i}h),\Lambda(t_{0}+\widetilde{c}_{i}h)), \quad i=1,\ldots,s-1 \\ Z(t) &= Z^{f}(t) + Z^{r}(t) \\ \dot{\widetilde{Z}}^{r}(t_{0}+\widetilde{c}_{i}h) &= r(t_{0}+\widetilde{c}_{i}h,Y(t_{0}+\widetilde{c}_{i}h),\Lambda(t_{0}+\widetilde{c}_{i}h)), \quad i=0,\ldots,s \\ \widetilde{Z}(t) &= Z^{f}(t) + \widetilde{Z}^{r}(t) \end{split}$$

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$$0 = g(t_0 + \tilde{c}_i h, Y(t_0 + \tilde{c}_i h)), \quad i = 0, ..., s$$

$$0 = g_t(t_1, y_1) + g_y(t_1, y_1)v(t_1, y_1, z_1)$$

$$0 = k(t_0 + \tilde{c}_i h, Y(t_0 + \tilde{c}_i h), \tilde{Z}(t_0 + \tilde{c}_i h)), \quad i = 0, ..., s$$

$$0 = \sum_{j=1}^{s} b_j c_j^{i-1} k(t_0 + c_j h, Y(t_0 + c_j h), Z(t_0 + c_j h)), \quad i = 1, ..., s - 1$$

$$0 = k(t_1, y_1, z_1).$$

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The numerical solution at $t_1 = t_0 + h$ is taken to be

$$y_1 := Y(t_1)$$

 $z_1 := Z(t_1) - h\widetilde{b}_s \widetilde{\mu}(t_1).$

These discontinuous collocation type methods can be shown to be equivalent to the Gauss-Lobatto SPARK and EMPRK methods.

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Discontinuous Collocation Type Methods

Theorem

A SPARK/EMPRK method with distinct values for c_j and for \tilde{c}_j is a discontinuous collocation method iff the coefficients satisfy

$$\sum_{j=1}^{s} a_{ij} c_j^{k-1} = \frac{c_i^k}{k}, \qquad \sum_{j=1}^{s} b_j c_j^{k-1} = \frac{1}{k}, \qquad k = 1, \dots, s$$

$$\sum_{j=0}^{s} \tilde{a}_{ij} \tilde{c}_j^{k-1} = \frac{c_i^k}{k}, \qquad \sum_{j=0}^{s} \tilde{b}_j \tilde{c}_j^{k-1} = \frac{1}{k}, \qquad k = 1, \dots, s - 1$$

$$\tilde{a}_{i0} = \tilde{b}_0, \qquad \tilde{a}_{is} = 0$$

$$\sum_{j=1}^{s} \tilde{a}_{ij} c_j^{k-1} = \frac{\tilde{c}_i^k}{k}, \qquad k = 1, \dots, s - 1$$

$$\sum_{j=0}^{s} \tilde{a}_{ij} \tilde{c}_j^{k-1} = \frac{\tilde{c}_i^k}{k}, \qquad k = 1, \dots, s + 1.$$

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Theorem

The (s, s)-Gauss-Lobatto SPARK and EMPRK methods with consistent initial values have local order 2s, i.e., for $|h| \le h_0$,

$$y_1 - y(t_1) = \mathcal{O}(h^{2s+1}), \qquad z_1 - z(t_1) = \mathcal{O}(h^{2s+1}).$$

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Theorem

Consider the (s, s)-Gauss-Lobatto SPARK / EMPRK methods with consistent initial conditions (y_0, z_0) at time t_0 . Then the (s, s)-Gauss-Lobatto SPARK / EMPRK methods are convergent of order 2s, i.e.

$$y_n - y(t_n) = \mathcal{O}(h^{2s}), \quad z_n - z(t_n) = \mathcal{O}(h^{2s}),$$

where y_n and z_n are the numerical solution at time $t_n := t_0 + nh$, for $nh \leq Const$.

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Consider a simple pendulum of mass m and length ℓ .

$$L(t,q,v) := T - U, \qquad T := \frac{1}{2}m(v_1^2 + v_2^2), \qquad U := -m\gamma q_2,$$

$$0=g(t,q)=rac{1}{2}(q_1^2+q_2^2-\ell^2).$$

- m is the mass (m = 1)
- ℓ is the length of the bob ($\ell = 1$)
- γ is the acceleration due to gravity ($\gamma = 1$)

The initial conditions are

$$q_0 = (1 \quad 0)^T, \qquad v_0 = (0 \quad 0)^T.$$

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The system is integrated from 0 to 100.



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(2,2)-Gauss-Lobatto, h = .05



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(2,2)–Gauss-Lobatto, h = .01



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Consider a skate sliding on an inclined plane without friction.

$$L(t, q, v) := T - U$$

$$T := \frac{1}{4}m(v_1^2 + v_2^2 + v_3^2 + v_4^2) \qquad U := -\frac{1}{2}m\gamma\sin(\beta)(q_1 + q_3)$$

$$0 = g(t, q) = \frac{1}{2}((q_3 - q_1)^2 + (q_4 - q_2)^2 - \ell^2)$$

$$0 = k(t, q, v) = -(q_4 - q_2)(v_1 + v_3) + (q_3 - q_1)(v_2 + v_4)$$

- m is the mass of the skate (m = 1)
- ℓ is the length of the skate ($\ell = 2$)
- γ is the acceleration due to gravity, β is the incline of the plane $(\gamma \sin(\beta) = 1)$

The initial conditions are

$$q_0 = egin{pmatrix} -1/2 & 0 & 1/2 & 0 \end{pmatrix}^T \qquad v_0 = egin{pmatrix} 0 & -1/2 & 0 & 1/2 \end{pmatrix}^T.$$

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The system is integrated from 0 to 10.



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(2,2)–Gauss-Lobatto SPARK Method, h = .1



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(2,2)-Gauss-Lobatto EMPRK Method, h = .1



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- SPARK and EMPRK methods are Runge-Kutta style methods for solving overdetermined mixed index 2 and 3 DAEs.
- These methods have a unique solution.
- These methods can be expressed as discontinuous collocation methods.
- For the Gauss-Lobatto coefficients, the SPARK and EMPRK methods are of order 2*s*.

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