Noncommutative geometry and time-frequency analysis

Franz Luef

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Hermitian structure:

Let \mathcal{A} be a unital C^* -algebra. Then a vector space V is a **left Hilbert** \mathcal{A} -module, i.e. $(\mathcal{A}, g) \mapsto \mathcal{A} \cdot g$ is a map from $V \times \mathcal{A} \to \mathcal{A}$, with a pairing $\mathcal{A} \langle ., . \rangle$ such that for all $f, g, h \in V$:

• $_{\mathcal{A}}\langle A \cdot f, g \rangle = A_{\mathcal{A}}\langle f, g \rangle$ for all $A \in \mathcal{A}$;

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Suppose compatible connections ∇₁, ∇₂ are given on V and ∂₁, ∂₂ are derivations on A. Then a gauge connection on V is defined by these two compatible connections and its curvature is defined by

$$F_{12} = \nabla_1 \nabla_2 - \nabla_2 \nabla_1 - \nabla_{[\partial_1, \partial_2]}.$$

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Twisted convolution of **a** and **b** is defined by

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Spectrally invariant subalgebras of noncommutative tori

$$\mathcal{A}_{s}^{1}(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c) = \{A = \sum_{\lambda} a(k, l) \pi(\alpha k, \beta l) : \|\mathbf{a}\|_{\ell_{s}^{1}} < \infty\}$$

with $\mathbf{a}\|_{\ell_{\mathbf{s}}^{1}} = \sum_{k,l} |a(k,l)| (1+|k|^{2}+|l^{2}|)^{s/2}$ smooth noncommutative torus $\mathcal{A}^{\infty}(\alpha\mathbb{Z} \times \beta\mathbb{Z}, c) = \bigcap_{s \geq 0} \mathcal{A}_{s}^{1}(\alpha\mathbb{Z} \times \beta\mathbb{Z}, c)$

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Theorem:

 $\mathcal{A}^1_s(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$ and $\mathcal{A}^{\infty}(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$ are spectrally invariant subalgebras of the **noncommutative torus** $C^*(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$.

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Wireless communication – OFDM

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Basic Idea:

A large number of parallel narrow-band subchannels is used instead of a single wide-band channel to transport information, i.e. sending discrete symbols on a continuous channel (acoustive wave).

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Idealization: Transmission of an infinite number of symbols

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- Transmitter: Assuming N suchannels, a bandwidth of W Hz, symbol length of a_T seconds, and subchannel separation $b_F := W/N$, the transmitter of a general OFDM system uses the following waveforms $g_l(t) = g(t)e^{2\pi i l b_F t}$ for l = 0, ..., N 1.
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Transmitter sends a superposition of individual symbols:

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- translation $T_x f(t) = f(t x)$ for $x \in \mathbb{R}$, modulation $M_{\omega}f(t) = e^{2\pi i t \cdot \omega} f(t)$ for $\omega \in \widehat{\mathbb{R}}$
- time-frequency shift $\pi(x,\omega)f(t) = M_{\omega}T_{x}f(t)$ for $(x,\omega) \in \mathbb{R} \times \widehat{\mathbb{R}}$

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Time-frequency representations

■ Short-Time Fourier Transform (STFT)

$$V_g f(x,\omega) = \int_{\mathbb{R}} f(t)\overline{g}(t-x)e^{-2\pi i t\omega}dt = \langle f, \pi(x,\omega)g \rangle$$

For $\varphi(t) = e^{-\pi t^2}$ we have $V_{\varphi}\varphi(x,\omega) = e^{-\pi i x \cdot \omega} e^{-\frac{1}{2}(x^2 + \omega^2)}$.

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Time-frequency localization

Fix a "nice" window function g, e.g. Gauss function, and look at the decay or summability properties of a time-frequency representation, e.g. STFT.

Feichtinger proposed to impose integrability conditions on a time-frequency representation and showed how this allows one to define Banach spaces of functions and distributions:

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Modulation spaces:

Suppose g is a Schwartz function. Then $f \in S'(\mathbb{R})$ is in the **modulation space** $M_s^{p,q}(\mathbb{R})$ if

$$\|f\|_{M^{p,q}_s} = \left(\int_{\mathbb{R}} \left(\int_{\mathbb{R}} |V_g f(x,\omega)|^p (1+|x|+|\omega|)^{sp} dx\right)^{q/p} d\omega\right)^{1/q}$$

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Resolution of Identity

Matrix coefficients of the Schrödinger representation satisfy an orthogonality relation:

Moyal's Identity

For f, g, h, k are in $L^2(\mathbb{R})$ we have

$$\langle V_g f, V_h k \rangle_{L^2(\mathbb{R}^2)} = \langle f, k \rangle_{L^2} \overline{\langle h, k \rangle_{L^2}},$$

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Resolution of identity Suppose $||g||_2 = 1$

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The Gabor frame operator is a Hilbert module operator: $S_{g,h}f = {}_{\Lambda}\langle g,h \rangle \cdot f.$

Left action of
$$\mathcal{A}_{s}^{1}(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$$
 on $M_{s}^{1}(\mathbb{R}^{d})$ by
 $D_{\mathbf{a}}g = \pi_{\Lambda}(\mathbf{a}) \cdot g = \left[\sum_{k,l} a(\alpha k, \beta l)\pi(\lambda)\right]g$ for $\mathbf{a} \in \ell_{s}^{1}(\alpha \mathbb{Z} \times \beta \mathbb{Z})$

$$_{\Lambda}\langle f,g\rangle = \sum_{k,l} \langle f,\pi(\alpha k,\beta l)g\rangle\pi(\alpha k,\beta l)$$

• For
$$f, g \in M_s^1(\mathbb{R})$$
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 $_{\Lambda}\langle f, g \rangle = \pi_{\Lambda}(V_g f) = \sum_{k,l} \langle f, \pi(\alpha k, \beta l)g \rangle \pi(\alpha k, \beta l)$
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Hilbert $C^*(\alpha^{-1}\mathbb{Z}, \beta^{-1}\mathbb{Z}, \overline{c})$ -modules

For
$$\mathbf{b} \in \ell_s^1(lpha^{-1}\mathbb{Z}, eta^{-1}\mathbb{Z})$$
 and $g \in M_s^1(\mathbb{R})$ we define

$$g \cdot \pi_{\Lambda^{\circ}}(\mathbf{b}) = \sum_{m,n} \pi(\frac{m}{\beta}, \frac{n}{\alpha})^* g \overline{b}(\frac{m}{\beta}, \frac{n}{\alpha})$$

$$\langle f,g\rangle_{\Lambda^{\circ}} = \sum_{m,n} \pi(\frac{m}{\beta},\frac{n}{\alpha})^* \langle g,\pi(\frac{m}{\beta},\frac{n}{\alpha})f\rangle \ f,g \in M^1_s(\mathbb{R}^d)$$

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 $M^1_s(\mathbb{R}^d)$ is an equivalence bimodule between $\mathcal{A}^1_s(\alpha\mathbb{Z}\times\beta\mathbb{Z},c)$ and $\mathcal{A}^1_s(\alpha^{-1}\mathbb{Z},\beta^{-1}\mathbb{Z},\overline{c})$,

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Theorem:

 $M^1_{\mathcal{S}}(\mathbb{R})$ is a finitely generated projective right $\mathcal{A}^1_{\mathcal{S}}(\alpha^{-1}\mathbb{Z}, \beta^{-1}\mathbb{Z}, \overline{c}).$

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In terms of noncommutative geometry this statement translates into the following:

Theorem:

A Gabor frame $\mathcal{G}(g, \alpha \mathbb{Z} \times \beta \mathbb{Z})$ ia a line bundle over the twisted group algebra $\mathcal{A}^{1}(\alpha^{-1}\mathbb{Z}, \beta^{-1}\mathbb{Z}, \overline{c})$.

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Frames for Hilbert C^* -modules

Let A be a unital C*-algebra. A sequence
 {g_j : j = 1, ..., n} in a (left) Hilbert A-module AV is called
 a standard module frame if there are positive reals C, D
 such that

$$C_{\mathcal{A}}\langle f, f \rangle \leq \sum_{j=1}^{n} {}_{\mathcal{A}}\langle f, g_j \rangle_{\mathcal{A}} \langle g_j, f \rangle \leq D_{\mathcal{A}}\langle f, f \rangle$$

for each $f \in {}_{\mathcal{A}}V$.

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Existence of multi-window Gabor frames

Theorem:

Then there exist $g_1, ..., g_n$ in $M^1_s(\mathbb{R})$ such that for all f in $L^2(\mathbb{R})$

$$\|f\|_2^2 = \sum_{i=1}^n \sum_{k,l} |\langle f, \pi(\alpha k, \beta l) g_i \rangle|^2.$$

 $\mathcal{G}(g_1, ..., g_n, \alpha \mathbb{Z} \times \beta \mathbb{Z})$ is a multi-window Gabor frame for $L^2(\mathbb{R})$).

By a result of Feichtinger and Gröchenig this implies that $\mathcal{G}(g_1, ..., g_n, \alpha \mathbb{Z} \times \beta \mathbb{Z})$ is a multi-window Gabor frame for the class of modulation spaces $M_m^{p,q}(\mathbb{R})$.

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G(g, αZ × βZ) is a Gabor frame if and only if M¹_s(ℝ) is a singly-generated projective right A¹_s(Λ°, c̄)-module.

Theorem:

Given a lattice Λ . Then there exists a $g \in M^1_s(\mathbb{R}^d)$ if and only if $\operatorname{vol}(\Lambda) < 1$.

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$$\mathsf{A} = \iint_{\mathbb{R}^2} \mathsf{a}(z) \pi(z) dz$$

for $a \in L^1(\mathbb{R}^2)$, i.e. integrated representation of the Schrödinger representation.

An equivalent way to express the Weyl quantization of an operator, **spreading representation**, pseudodifferential operators: **Moyal plane**.

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We denote all operators A with $a \in S(\mathbb{R}^2)$ with $\mathcal{A}^{\infty}(\mathbb{R}^2, c)$ and call it the **smooth Moyal plane**.

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Stone-von Neumann:

 $\mathcal{A}^\infty(\mathbb{R}^2,c)$ is Morita equivalent to the complex numbers $\mathbb{C}.$

Define a Hermitian structure on $\mathcal{A}^{\infty}(\mathbb{R}^2, c)$ via

$$_{\mathbb{R}^{2}}\langle f,g
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and as left action:

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Connections on noncommutative tori and Moyal plane

Derivations ∂_1 and ∂_2 on $\mathcal{A}^{\infty}(\mathbb{R}^2, c)$:

$$\partial_1 A = 2\pi i \iint_{\mathbb{R}^2} xa(x,\omega)\pi(x,\omega)dxd\omega$$

$$\partial_2 A = 2\pi i \iint_{\mathbb{R}^2} \omega a(x,\omega) \pi(x,\omega) dx d\omega$$

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We define covariant derivatives ∇_1 and ∇_2 on $\mathcal{S}(R)$:

$$(\nabla_1 g)(t) = 2\pi i t g(t) \text{ and } (\nabla_2 g)(t) = g'(t).$$

 $\nabla_2 (A \cdot g) = (\partial_2 A) \cdot g + A \cdot (\nabla_2 g)$
 $\partial_2 (A \cdot g) = 2\pi i \omega \iint a(z)\pi(z)gdz + \iint a(z)\pi(z)g'dz.$

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For example,

$$2\pi i\omega V_g f(x,\omega) = V_g f'(x,\omega) + V_{g'} f(x,\omega)$$

covariant derivatives on S(R):

 $(\nabla_1 g)(t) = 2\pi i \alpha t g(t)$ and $(\nabla_2 g)(t) = \beta g'(t)$

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$$\partial_1(A) = 2\pi i\alpha \sum_{k,l} ka_{kl}\pi(\alpha k,\beta l)$$
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Connections on $C^*(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$

 $abla_1 g(t) = 2\pi i lpha t g(t)$ and $abla_2 g(t) = eta g'(t)$

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Connections

$$abla_1^\circ g(t) = 2\pi i eta^{-1} t g(t)$$
 and $abla_2^\circ g(t) = lpha^{-1} g'(t)$

curvature F_{12} on $C^*(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$:

$$F_{12} = -2\pi i\alpha\beta I$$

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Let's consider the case $A = {}_{\Lambda} \langle f, g \rangle$ in more detail:

 $\nabla_i(\Lambda\langle f,g\rangle \cdot h) = \delta_i(\Lambda\langle f,g\rangle) \cdot h + \Lambda\langle f,g\rangle \cdot \nabla_i h$ provides a relation between a Gabor system $\mathcal{G}(g,\Lambda)$ and $\mathcal{G}(\nabla_i g,\Lambda)$.

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Complex structures on noncommutative tori and Moyal plane

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Balian-Low:

Let $\mathcal{G}(g, \alpha \mathbb{Z} \times \beta \mathbb{Z})$ be a Riesz basis for its closed span \mathcal{H} in $L^2(\mathbb{R})$. Then $\nabla_i g$ or $\nabla_i h$ is not in \mathcal{H} , where h denotes the canonical dual Gabor atom $h = S_{g,g}^{-1}$.

Proof is based on an observation of G. Battle, which uses the left Leibniz property for $A = \pi(\alpha k, \beta l)$ implies:

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$$\begin{split} \langle \nabla_1 g, \nabla_2 h \rangle &= \langle \sum_{k,l} \langle \nabla_1 g, \pi(\alpha k, \beta l) h \rangle \pi(\alpha k, \beta l) g, \nabla_2 h \rangle \\ &= \sum_{k,l} \langle \pi(-\alpha k, -\beta l) g, \nabla_1 h \rangle \langle \nabla_2 g, \pi(-\alpha k, -\beta l) h \rangle \\ &= \langle \nabla_2 g, \sum_{k,l} \langle \nabla_1 h, \pi(\alpha k, \beta l) g \rangle \pi(\alpha k, \beta l) h \rangle \\ &= \langle \nabla_2 g, \nabla_1 h \rangle \end{split}$$

However, $\nabla_1\nabla_2-\nabla_2\nabla_1=2\pi i l$, canonical commutation relations, gives

$$1 = \langle g, h \rangle = \langle \nabla_2 g, \nabla_1 h \rangle - \langle \nabla_1 g, \nabla_2 h \rangle = 0.$$

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Projections in noncommutative tori

Theorem:

Let $\mathcal{G}(g, \Lambda)$ be a Gabor system on $L^2(\mathbb{R}^d)$. Then $p_g = {}_{\Lambda}\langle g, g \rangle$ is a projection in $C^*(\Lambda, c)$ if and only if $g \cdot \langle g, g \rangle_{\Lambda^\circ} = g$.

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Proposition:

Let g be in $_{\Lambda}V_{\Lambda^{\circ}}$. Then $P_g := _{\Lambda}\langle g, g \rangle$ is a projection in $C^*(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$ if and only if $g \langle g, g \rangle_{\Lambda^{\circ}} = g$. If $g \in M^1_s(\mathbb{R})$ or $\mathscr{S}(\mathbb{R})$, then P_g gives a projection in $\mathcal{A}^1_s(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$ or $\mathcal{A}^{\infty}(\alpha \mathbb{Z} \times \beta \mathbb{Z}, c)$, respectively.

First we assume that $g\langle g,g\rangle_{\Lambda^\circ}=g$ for some g in $_\Lambda V_{\Lambda^\circ}.$ Then we have that

$$\begin{split} P_g^2 &= {}_{\Lambda} \langle g,g \rangle_{\Lambda} \langle g,g \rangle = {}_{\Lambda} \langle \langle g,g \rangle g,g \rangle = {}_{\Lambda} \langle g \langle g,g \rangle_{\Lambda^\circ},g \rangle = {}_{\Lambda} \langle g,g \rangle \\ \text{and } P_g^* &= P_g. \end{split}$$

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The unit sphere of the Hilbert $C^*(\beta^{-1}\mathbb{Z} \times \alpha^{-1}\mathbb{Z}, \overline{c})$ -module V is defined by $S(V) = \{g \in V_{:} \langle g, g \rangle_{\Lambda^{\circ}} = I\}$, that is the set of all tight Gabor frames.

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 $g_1(t) = (2)^{1/4} e^{-\pi t^2}$ a Gaussian, $g_2(t) = (\frac{\pi}{2})^{1/2} \frac{1}{\cosh(\pi t)}$ the hyperbolic secant and $g_3(t) = e^{-\pi |t|}$ the two-sided exponential.

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The case of the Gaussian g_1 is known as **Boca's projection**. In Manin's work $p_{g_1} = {}_{\Lambda}\langle g_1, g_1 \rangle$ appears as **quantum theta functions**.

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- projective representation of phase space ℝ^{2d}, representations of the Heisenberg group
- time-frequency localization modulation spaces

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- Gabor frames
- noncommutative tori
- strong Morita equivalence of operator algebras
- connections on noncommutative tori

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