

Scalar Curvature and Gauss-Bonnet Theorem for Noncommutative Tori

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Spectral Triples

Noncommutative geometric spaces are described by spectral triples:

$$(\mathcal{A}, \mathcal{H}, D),$$

$$\pi : \mathcal{A} \rightarrow \mathcal{L}(\mathcal{H}) \quad (*\text{-representation}),$$

$$D = D^* : \text{Dom}(D) \subset \mathcal{H} \rightarrow \mathcal{H},$$

$$D \pi(a) - \pi(a) D \in \mathcal{L}(\mathcal{H}).$$

Examples.

$$(C^\infty(M), L^2(M, S), D = \text{Dirac operator}).$$

$$\left(C^\infty(\mathbb{S}^1), L^2(\mathbb{S}^1), \frac{1}{i} \frac{\partial}{\partial x} \right).$$

Noncommutative Local Invariants

The local geometric invariants such as scalar curvature of (A, \mathcal{H}, D) are detected by the high frequency behavior of the spectrum of D and the action of A via heat kernel asymptotic expansions of the form

$$\text{Trace} (a e^{-tD^2}) \sim_{t \searrow 0} \sum_{j=0}^{\infty} a_j(a, D^2) t^{(-n+j)/2}, \quad a \in A.$$

Noncommutative 2-Torus $A_\theta = C(\mathbb{T}_\theta^2)$

It is the universal C^* -algebra generated by U and V s.t.

$$U^* = U^{-1},$$

$$V^* = V^{-1},$$

$$VU = e^{2\pi i\theta}UV,$$

where $\theta \in \mathbb{R}$ is fixed.

The geometry of the Kronecker foliation $dy = \theta dx$ on the ordinary torus $\mathbb{R}^2/\mathbb{Z}^2$ is closely related to the structure of this algebra.

A representation of A_θ :

$$U\xi(x) = e^{2\pi i x}\xi(x), \quad V\xi(x) = \xi(x + \theta), \quad \xi \in L^2(\mathbb{R}).$$

Action of $\mathbb{T}^2 = (\frac{\mathbb{R}}{2\pi\mathbb{Z}})^2$ on A_θ and Smooth Elements

•

$$\begin{aligned}\alpha_s &: A_\theta \rightarrow A_\theta, \quad s \in \mathbb{R}^2, \\ \alpha_s(U^m V^n) &= e^{is \cdot (m,n)} U^m V^n, \quad m, n \in \mathbb{Z}.\end{aligned}$$

•

$$\begin{aligned}A_\theta^\infty &:= \{a \in A_\theta; \quad s \mapsto \alpha_s(a) \text{ is smooth from } \mathbb{R}^2 \text{ to } A_\theta\} \\ &= \left\{ \sum_{m,n \in \mathbb{Z}} a_{m,n} U^m V^n \in A_\theta; \quad (a_{m,n}) \in \mathcal{S}(\mathbb{Z}^2) \right\}.\end{aligned}$$

•

$$\delta_j = \frac{\partial}{\partial s_j} \Big|_{s=0} \alpha_s : A_\theta^\infty \rightarrow A_\theta^\infty.$$

The Derivations δ_1, δ_2 and the Volume Form

- $\delta_1, \delta_2 : A_\theta^\infty \rightarrow A_\theta^\infty$ are defined by:

$$\delta_1(U) = U, \quad \delta_1(V) = 0, \quad \delta_2(U) = 0, \quad \delta_2(V) = V,$$

$$\delta_i(ab) = \delta_i(a)b + a\delta_i(b), \quad a, b \in A_\theta^\infty.$$

- Tracial state $\varphi_0 : A_\theta \rightarrow \mathbb{C}$ (analog of integration):

$$\varphi_0(1) = 1, \quad \varphi_0(U^m V^n) = 0 \quad \text{if } (m, n) \neq (0, 0).$$

Conformal Structure on A_θ (Connes)

The Dolbeault operators associated with $\tau \in \mathbb{C}, \Im(\tau) > 0$ are

$$\partial = \delta_1 + \bar{\tau}\delta_2 : \mathcal{H}_0 \rightarrow \mathcal{H}^{(1,0)},$$

$$\bar{\partial} = \delta_1 + \tau\delta_2 : \mathcal{H}_0 \rightarrow \mathcal{H}^{(0,1)}.$$

The conformal structure represented by τ is encoded in

$$\psi(a, b, c) = -\varphi_0(a \partial(b) \bar{\partial}(c)), \quad a, b, c \in A_\theta^\infty,$$

which is a positive Hochschild cocycle.

Conformal Perturbation (Connes-Tretkoff)

Let $h = h^* \in A_\theta^\infty$ and replace the trace φ_0 by

$$\varphi : A_\theta \rightarrow \mathbb{C},$$

$$\varphi(a) := \varphi_0(a e^{-h}), \quad a \in A_\theta.$$

φ is a KMS state with the modular group

$$\sigma_t(a) = e^{ith} a e^{-ith}, \quad a \in A_\theta,$$

and the modular automorphism

$$\Delta(a) := \sigma_i(a) = e^{-h} a e^h, \quad a \in A_\theta.$$

$$\varphi(ab) = \varphi(b \Delta(a)), \quad a, b \in A_\theta.$$

A Spectral Triple $(A_\theta^\infty, \mathcal{H}, D)$

$$\mathcal{H} := \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)},$$

$$a \mapsto \begin{pmatrix} a & 0 \\ 0 & a \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H},$$

$$D := \begin{pmatrix} 0 & \partial_\varphi^* \\ \partial_\varphi & 0 \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H},$$

$$\partial_\varphi := \partial = \delta_1 + \bar{\tau}\delta_2 : \mathcal{H}_\varphi \rightarrow \mathcal{H}^{(1,0)}.$$

Anti-Unitary Equivalence of the Laplacians

$$D^2 = \begin{pmatrix} \partial_\varphi^* \partial_\varphi & 0 \\ 0 & \partial_\varphi \partial_\varphi^* \end{pmatrix} : \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)}.$$

Lemma: Let

$$k = e^{h/2}.$$

We have

$$\begin{aligned} \partial_\varphi^* \partial_\varphi : \mathcal{H}_\varphi &\rightarrow \mathcal{H}_\varphi & \sim & k \bar{\partial} \partial k : \mathcal{H}_0 \rightarrow \mathcal{H}_0, \\ \partial_\varphi \partial_\varphi^* : \mathcal{H}^{(1,0)} &\rightarrow \mathcal{H}^{(1,0)} & \sim & \bar{\partial} k^2 \partial : \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}^{(1,0)}. \end{aligned}$$

Derivation of the Asymptotic Expansion

Approximate e^{-tD^2} by pseudodifferential operators:

$$e^{-tD^2} = \frac{1}{2\pi i} \int_C e^{-t\lambda} (D^2 - \lambda)^{-1} d\lambda,$$

$$B_\lambda (D^2 - \lambda) \approx 1,$$

$$\sigma(B_\lambda) = b_0 + b_1 + b_2 + \dots .$$

Connes' pseudodifferential calculus (1980)

- Symbols $\rho : \mathbb{R}^2 \rightarrow A_\theta^\infty \Rightarrow P_\rho : A_\theta^\infty \rightarrow A_\theta^\infty$

$$P_\rho(a) = (2\pi)^{-2} \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} e^{-is \cdot \xi} \rho(\xi) \alpha_s(a) ds d\xi, \quad a \in A_\theta^\infty.$$

- Differential operators:

$$\rho(\xi_1, \xi_2) = \sum a_{ij} \xi_1^i \xi_2^j, \quad a_{ij} \in A_\theta^\infty \quad \Rightarrow \quad P_\rho = \sum a_{ij} \delta_1^i \delta_2^j.$$

- Ψ DO's on A_θ^∞ form an algebra:

$$\sigma(PQ) \sim \sum_{\ell_1, \ell_2 \geq 0} \frac{1}{\ell_1! \ell_2!} \partial_1^{\ell_1} \partial_2^{\ell_2} (\rho(\xi)) \delta_1^{\ell_1} \delta_2^{\ell_2} (\rho'(\xi)).$$

Symbol of the first Laplacian

$$\sigma(k\bar{\partial}\partial k) = a_2(\xi) + a_1(\xi) + a_0(\xi),$$

where

$$a_2(\xi) = \xi_1^2 k^2 + |\tau|^2 \xi_2^2 k^2 + 2\Re(\tau)\xi_1\xi_2 k^2,$$

$$a_1(\xi) = 2\xi_1 k \delta_1(k) + 2|\tau|^2 \xi_2 k \delta_2(k) + 2\Re(\tau)\xi_1 k \delta_2(k) + 2\Re(\tau)\xi_2 k \delta_1(k),$$

$$a_0(\xi) = k\delta_1^2(k) + |\tau|^2 k\delta_2^2(k) + 2\Re(\tau)k\delta_1\delta_2(k).$$

$$b_n = - \sum_{\substack{2+j+\ell_1+\ell_2-k=n, \\ 0 \leq j < n, 0 \leq k \leq 2}} \frac{1}{\ell_1! \ell_2!} \partial_1^{\ell_1} \partial_2^{\ell_2} (b_j) \delta_1^{\ell_1} \delta_2^{\ell_2} (a_k) b_0, \quad n > 0.$$

$$b_0 = a_2'^{-1} = (\xi_1^2 k^2 + |\tau|^2 \xi_2^2 k^2 + 2\Re(\tau)\xi_1\xi_2 k^2 - \lambda)^{-1}.$$

Weyl's law for \mathbb{T}_θ^2

Theorem. (Khalkhali-F.) Let

$$N(\lambda) = \#\{\lambda_j \leq \lambda\}$$

be the eigenvalue counting function of D^2 . We have

$$N(\lambda) \sim \frac{\pi}{\mathfrak{S}(\tau)} \varphi_0(e^{-h}) \lambda \quad (\lambda \rightarrow \infty).$$

Equivalently:

$$\lambda_j \sim \frac{\mathfrak{S}(\tau)}{\pi \varphi(1)} j \quad (j \rightarrow \infty).$$

Connes' trace theorem for \mathbb{T}_θ^2

Classical symbols: $\rho : \mathbb{R}^2 \rightarrow A_\theta^\infty$

$$\rho(\xi) \sim \sum_{i=-\infty}^{\infty} \rho_{m-i}(\xi) \quad (\xi \rightarrow \infty),$$

$$\rho_{m-i}(t\xi) = t^{m-i} \rho_{m-i}(\xi), \quad t > 0, \quad \xi \in \mathbb{R}^2.$$

Theorem. (Khalkhali-F.) For any classical symbol ρ of order -2 on A_θ , we have

$$P_\rho \in \mathcal{L}^{1,\infty}(\mathcal{H}_0),$$

and

$$\mathrm{Tr}_\omega(P_\rho) = \frac{1}{2} \int_{\mathbb{S}^1} \varphi_0(\rho_{-2}(\xi)) d\xi.$$

$$\begin{aligned}
b_1 &= -(b_0 a_1 b_0 + \partial_1(b_0)\delta_1(a_2)b_0 + \partial_2(b_0)\delta_2(a_2)b_0), \\
b_2 &= -(b_0 a_0 b_0 + b_1 a_1 b_0 + \partial_1(b_0)\delta_1(a_1)b_0 + \\
&\quad \partial_2(b_0)\delta_2(a_1)b_0 + \partial_1(b_1)\delta_1(a_2)b_0 + \partial_2(b_1)\delta_2(a_2)b_0 + \\
&\quad (1/2)\partial_{11}(b_0)\delta_1^2(a_2)b_0 + (1/2)\partial_{22}(b_0)\delta_2^2(a_2)b_0 + \partial_{12}(b_0)\delta_{12}(a_2)b_0).
\end{aligned}$$

Connes' Rearrangement Lemma

For any $m = (m_0, m_1, \dots, m_\ell) \in \mathbb{Z}_{>0}^{\ell+1}$ and $\rho_1, \dots, \rho_\ell \in A_\theta^\infty$

$$\begin{aligned} & \int_0^\infty u^{|m|-2} (e^h u + 1)^{-m_0} \prod_1^\ell \rho_j (e^h u + 1)^{-m_j} du \\ &= e^{-(|m|-1)h} F_m(\Delta, \dots, \Delta) \left(\prod_1^\ell \rho_j \right), \end{aligned}$$

where

$$F_m(u_1, \dots, u_\ell) = \int_0^\infty \frac{x^{|m|-2}}{(x+1)^{m_0}} \prod_1^\ell \left(x \prod_1^j u_k + 1 \right)^{-m_j} dx.$$

Conformal Geometry of \mathbb{T}_θ^2 with $\tau = i$ (Cohen-Connes)

Let

$\lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \dots$ be the eigenvalues of $\partial_\varphi^* \partial_\varphi$,

and

$$\zeta(s) = \sum \lambda_j^{-s}, \quad \Re(s) > 1.$$

Then

$$\zeta(0) + 1 =$$

$$\varphi(f(\Delta)(\delta_1(e^{h/2})) \delta_1(e^{h/2})) + \varphi(f(\Delta)(\delta_2(e^{h/2})) \delta_2(e^{h/2})),$$

where

$$f(u) = \frac{1}{6}u^{-1/2} - \frac{1}{3} + \mathcal{L}_1(u) - 2(1+u^{1/2})\mathcal{L}_2(u) + (1+u^{1/2})^2\mathcal{L}_3(u),$$

$$\mathcal{L}_m(u) = (-1)^m (u-1)^{-(m+1)} \left(\log u - \sum_{j=1}^m (-1)^{j+1} \frac{(u-1)^j}{j} \right).$$

The Gauss-Bonnet theorem for \mathbb{T}_θ^2

Theorem. (Connes-Tretkoff; Khalkhali-F.) For any $\theta \in \mathbb{R}$, complex parameter $\tau \in \mathbb{C} \setminus \mathbb{R}$ and Weyl conformal factor $e^h, h = h^* \in A_\theta^\infty$, we have

$$\zeta(0) + 1 = 0.$$

Final Part of the Proof

$$\zeta(0) + 1 =$$

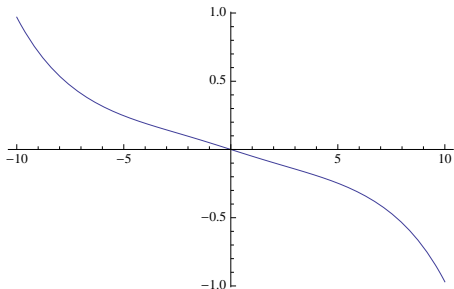
$$\begin{aligned} & \frac{2\pi}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_1 \left(\frac{h}{2} \right) \right) \delta_1 \left(\frac{h}{2} \right) \right) + \frac{2\pi|\tau|^2}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_2 \left(\frac{h}{2} \right) \right) \delta_2 \left(\frac{h}{2} \right) \right) \\ & + \frac{2\pi\Re(\tau)}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_1 \left(\frac{h}{2} \right) \right) \delta_2 \left(\frac{h}{2} \right) \right) + \frac{2\pi\Re(\tau)}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_2 \left(\frac{h}{2} \right) \right) \delta_1 \left(\frac{h}{2} \right) \right), \end{aligned}$$

where

$$K(x) = - \frac{\left(3x - 3 \sinh \left(\frac{x}{2} \right) - 3 \sinh(x) + \sinh \left(\frac{3x}{2} \right) \right) \operatorname{csch}^2 \left(\frac{x}{2} \right)}{3x^2}$$

is an odd entire function, and $\nabla = \log \Delta$.

$$K(x) = -\frac{x}{20} + \frac{x^3}{2240} - \frac{23x^5}{806400} + O(x^6).$$



Scalar Curvature for $(A_\theta^\infty, \mathcal{H}, D)$

It is the unique element $R \in A_\theta^\infty$ such that

$$\zeta_a(0) + \varphi_0(a) = \varphi_0(a R), \quad a \in A_\theta^\infty,$$

where

$$\zeta_a(s) := \text{Trace}(a |D|^{-2s}), \quad \text{Re}(s) \gg 0.$$

Equivalently, consider small-time heat kernel expansions:

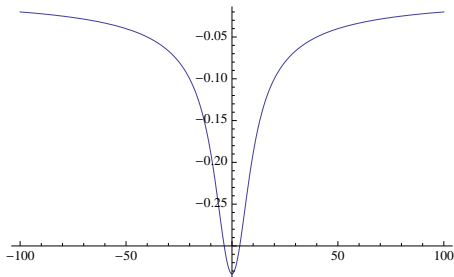
$$\text{Trace}(a e^{-tD^2}) \sim \sum_{n \geq 0} B_n(a, D^2) t^{\frac{n-2}{2}}, \quad a \in A_\theta^\infty.$$

Final Formula for the Scalar Curvature of \mathbb{T}_θ^2

Theorem. (Connes-Moscovici; Khalkhali-F.) Up to an overall factor of $\frac{-\pi}{\Im(\tau)}$, R is equal to

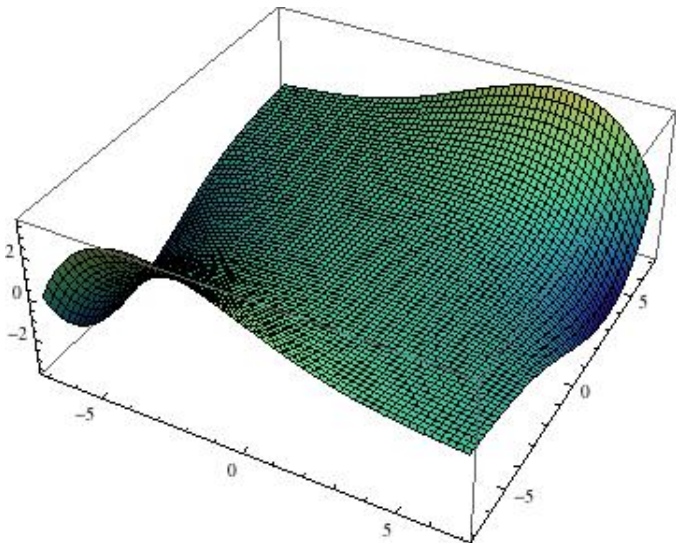
$$\begin{aligned} & R_1(\nabla)\left(\delta_1^2\left(\frac{h}{2}\right) + 2\tau_1\delta_1\delta_2\left(\frac{h}{2}\right) + |\tau|^2\delta_2^2\left(\frac{h}{2}\right)\right) \\ & + R_2(\nabla, \nabla)\left(\delta_1\left(\frac{h}{2}\right)^2 + |\tau|^2\delta_2\left(\frac{h}{2}\right)^2 + \Re(\tau)\left\{\delta_1\left(\frac{h}{2}\right), \delta_2\left(\frac{h}{2}\right)\right\}\right) \\ & + iW(\nabla, \nabla)\left(\Im(\tau)\left[\delta_1\left(\frac{h}{2}\right), \delta_2\left(\frac{h}{2}\right)\right]\right). \end{aligned}$$

$$R_1(x) = \frac{\frac{1}{2} - \frac{\sinh(x/2)}{x}}{\sinh^2(x/4)}.$$



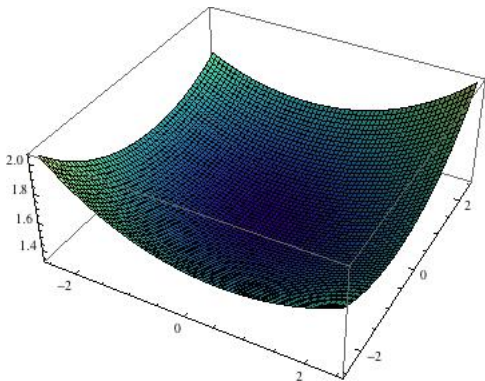
$$R_2(s, t) =$$

$$\frac{(1 + \cosh((s+t)/2))(-t(s+t) \cosh s + s(s+t) \cosh t - (s-t)(s+t + \sinh s + \sinh t - \sinh(s+t)))}{st(s+t) \sinh(s/2) \sinh(t/2) \sinh^2((s+t)/2)}$$



$$W(s, t) =$$

$$\frac{(-s - t + t \cosh s + s \cosh t + \sinh s + \sinh t - \sinh(s + t))}{st \sinh(s/2) \sinh(t/2) \sinh((s + t)/2)}.$$



Symbol of the second Laplacian

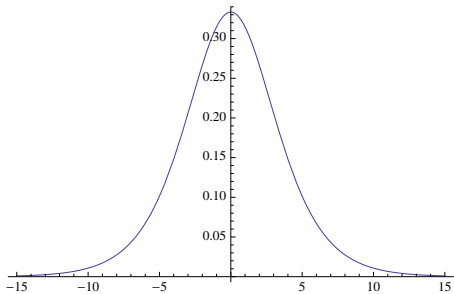
$$\sigma(\partial^* k^2 \partial) = c_2(\xi) + c_1(\xi),$$

where

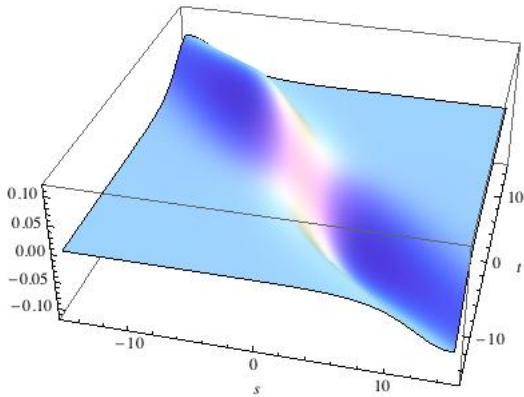
$$c_2(\xi) = \xi_1^2 k^2 + 2\tau_1 \xi_1 \xi_2 k^2 + |\tau|^2 \xi_2^2 k^2,$$

$$c_1(\xi) = (\delta_1(k^2) + \tau \delta_2(k^2)) \xi_1 + (\bar{\tau} \delta_1(k^2) + |\tau|^2 \delta_2(k^2)) \xi_2.$$

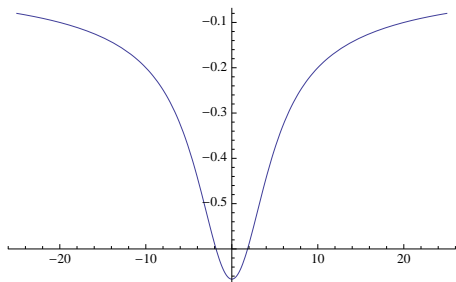
$$K_1(x) = \frac{2e^{x/2} (e^x(x-2) + x + 2)}{(e^x - 1)^2 x}$$



$$H_1(s, t) = \frac{\operatorname{csch}\left(\frac{s}{2}\right)\operatorname{csch}\left(\frac{t}{2}\right)\operatorname{csch}^2\left(\frac{s+t}{2}\right)\left(-s-t\right)\left(-\sinh(s+t)+s+\sinh(s)+t+\sinh(t)\right)-t(s+t)\cosh(s)+s(s+t)}{st(s+t)}$$

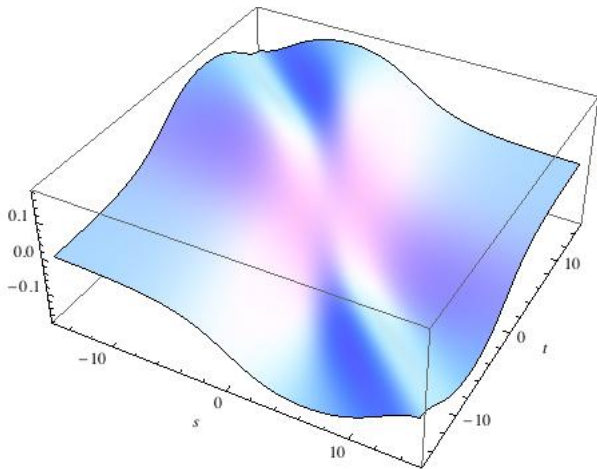


$$K_2(x) = -\frac{4e^x(\sinh(x) - x)}{(e^{x/2} - 1)^2 (e^{x/2} + 1)^2 x}$$

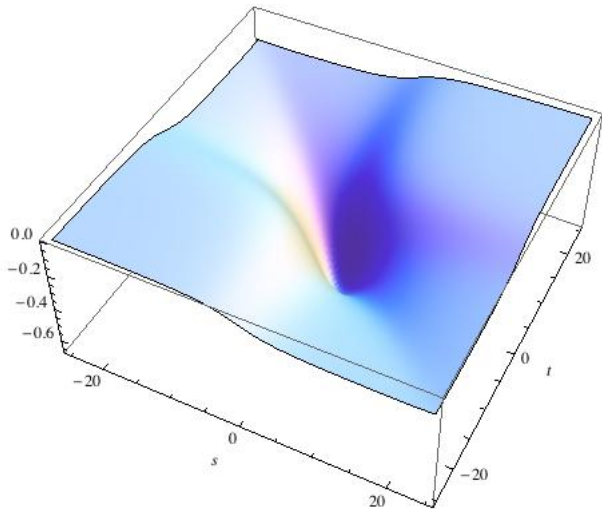


$$H_2(s, t) = \cosh\left(\frac{s+t}{2}\right) \times$$

$$\frac{\operatorname{csch}\left(\frac{s}{2}\right) \operatorname{csch}\left(\frac{t}{2}\right) \operatorname{csch}^2\left(\frac{s+t}{2}\right) (-s-t)(-\sinh(s+t)+s+\sinh(s)+t+\sinh(t))-t(s+t) \cosh(s)+s(s+t)}{st(s+t)}$$



$$W(s, t) = -\frac{4(-\sinh(s+t)+s \cosh(t)+t \cosh(s)-s+\sinh(s)-t+\sinh(t))}{st(-\sinh(s+t)+\sinh(s)+\sinh(t))}$$



Noncommutative 4-Torus \mathbb{T}_θ^4

$C(\mathbb{T}_\theta^4)$ is the universal C^* -algebra generated by 4 unitaries

$$U_1, U_2, U_3, U_4,$$

satisfying

$$U_k U_\ell = e^{2\pi i \theta_{k\ell}} U_\ell U_k,$$

for a skew symmetric matrix

$$\theta = (\theta_{k\ell}) \in M_4(\mathbb{R}).$$

Perturbed Laplacian on \mathbb{T}_θ^4

$$d = \partial \oplus \bar{\partial} : \mathcal{H}_\varphi \rightarrow \mathcal{H}_\varphi^{(1,0)} \oplus \mathcal{H}_\varphi^{(0,1)},$$

$$\Delta_\varphi := d^*d.$$

Lemma. (Khalkhali-F.) Up to an anti-unitary equivalence Δ_φ is given by

$$e^h \bar{\partial}_1 e^{-h} \partial_1 e^h + e^h \partial_1 e^{-h} \bar{\partial}_1 e^h + e^h \bar{\partial}_2 e^{-h} \partial_2 e^h + e^h \partial_2 e^{-h} \bar{\partial}_2 e^h,$$

where $\partial_1, \bar{\partial}_2$ are analogues of the Dolbeault operators.

Scalar Curvature for \mathbb{T}_θ^4

It is the unique element $R \in C^\infty(\mathbb{T}_\theta^4)$ such that

$$\operatorname{Res}_{s=1} \zeta_a(s) = \varphi_0(a R), \quad a \in C^\infty(\mathbb{T}_\theta^4),$$

where

$$\zeta_a(s) := \operatorname{Trace}(a \Delta_\varphi^{-s}), \quad \Re(s) \gg 0.$$

Final Formula for the Scalar Curvature of \mathbb{T}_θ^4

Theorem. (Khalkhali-F.) We have

$$R = e^{-h} k(\nabla) \left(\sum_{i=1}^4 \delta_i^2(h) \right) + e^{-h} H(\nabla, \nabla) \left(\sum_{i=1}^4 \delta_i(h)^2 \right),$$

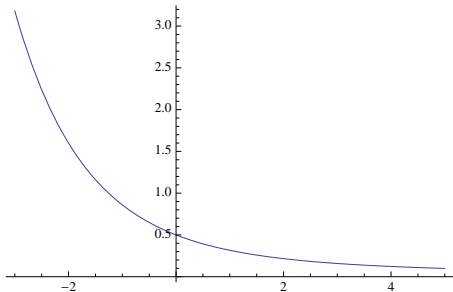
where

$$\nabla(a) = [-h, a], \quad a \in C(\mathbb{T}_\theta^4),$$

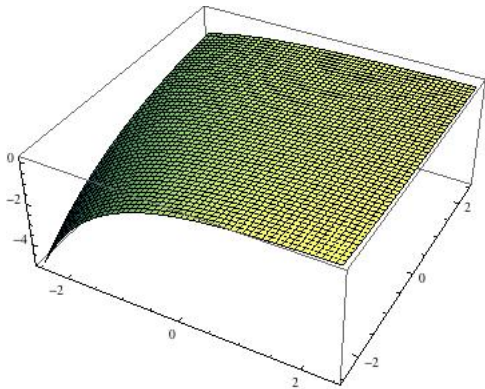
$$k(s) = \frac{1 - e^{-s}}{2s},$$

$$H(s, t) = -\frac{e^{-s-t} ((-e^s - 3)s(e^t - 1) + (e^s - 1)(3e^t + 1)t)}{4st(s+t)}.$$

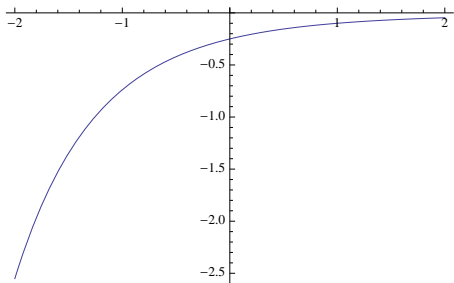
$$k(s) = \frac{1}{2} - \frac{s}{4} + \frac{s^2}{12} - \frac{s^3}{48} + \frac{s^4}{240} - \frac{s^5}{1440} + O(s^6).$$



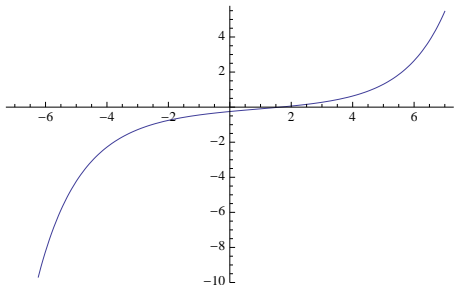
$$H(s, t) = \left(-\frac{1}{4} + \frac{t}{24} + O(t^3) \right) + s \left(\frac{5}{24} - \frac{t}{16} + \frac{t^2}{80} + O(t^3) \right) + s^2 \left(-\frac{1}{12} + \frac{7t}{240} - \frac{t^2}{144} + O(t^3) \right) + O(s^3).$$



$$\begin{aligned}
 H(s, s) &= -\frac{e^{-2s}(e^s - 1)^2}{4s^2} \\
 &= -\frac{1}{4} + \frac{s}{4} - \frac{7s^2}{48} + \frac{s^3}{16} - \frac{31s^4}{1440} + \frac{s^5}{160} + O(s^6).
 \end{aligned}$$



$$\begin{aligned}
 G(s) &:= H(s, -s) = \frac{-4s - 3e^{-s} + e^s + 2}{4s^2} \\
 &= -\frac{1}{4} + \frac{s}{6} - \frac{s^2}{48} + \frac{s^3}{120} - \frac{s^4}{1440} + \frac{s^5}{5040} + O(s^6).
 \end{aligned}$$



Einstein-Hilbert Action for \mathbb{T}_θ^4

Theorem. (Khalkhali-F.) We have the local expression (up to a factor of π^2)

$$\begin{aligned}\varphi_0(R) &= \frac{1}{2} \sum_{i=1}^4 \varphi_0\left(e^{-h} \delta_i^2(h)\right) \\ &\quad + \sum_{i=1}^4 \varphi_0\left(G(\nabla)(e^{-h} \delta_i(h)) \delta_i(h)\right).\end{aligned}$$

Extremum of the Einstein-Hilbert Action

Theorem. (Khalkhali-F.) For any Weyl factor $e^{-h} \in C^\infty(\mathbb{T}_\theta^4)$

$$\varphi_0(R) \leq 0,$$

and the equality happens if and only if h is a constant.

Proof.

$$\varphi_0(R) = \sum_{i=1}^4 \varphi_0(e^{-h} T(\nabla)(\delta_i(h)) \delta_i(h)),$$

where

$$T(s) = \frac{1}{2} \frac{e^{-s} - 1}{-s} + G(s) = \frac{-2s + e^s - e^{-s}(2s + 3) + 2}{4s^2}.$$

$$T(s) = \frac{1}{4} - \frac{s}{12} + \frac{s^2}{16} - \frac{s^3}{80} + \frac{s^4}{288} - \frac{s^5}{2016} + O(s^6).$$

