

The Spectral Geometry of Curved Noncommutative Tori

Masoud Khalkhali

COSY 2014

Spectral Triples (Connes)

- ▶ Noncommutative geometric spaces are described by *spectral triples* (first order elliptic PDE's on NC spaces), $(\mathcal{A}, \mathcal{H}, D)$, where

$$\pi : \mathcal{A} \rightarrow \mathcal{L}(\mathcal{H}) \quad (*\text{-representation}),$$

$$D = D^* : \text{Dom}(D) \subset \mathcal{H} \rightarrow \mathcal{H}, \quad \text{s.a.}$$

$$D \pi(a) - \pi(a) D \in \mathcal{L}(\mathcal{H}), \quad \text{bounded commutators,}$$

D has compact resolvent.

- ▶ Example: The Dirac spectral triple $(C^\infty(M), L^2(M, S), D)$, e.g. $D = \frac{1}{i} \frac{d}{dx}$, or the Cauchy-Riemann operator $\frac{\partial}{\partial \bar{z}}$.

The scalar curvature of a spectral triple

- ▶ *Connes' distance formula* recovers the metric from D , but a more difficult issue is how to define and compute the scalar curvature using D .
- ▶ A spectral triple is a NC Riemannian manifold. It is tempting to think that one might be able to define a Levi-Civita type connection for a spectral triple and then define the curvature of this connection. For many reasons this algebraic approach does not work in NCG in general.
- ▶ Instead one needs to import ideas of spectral geometry to NCG.

Spectral geometry: can one hear the shape of a drum?

- ▶ **Weyl's law:** for a compact Riemannian manifold M

$$N(\lambda) \sim \frac{\omega_n \text{Vol}(M)}{(2\pi)^n} \lambda^{\frac{n}{2}} \quad \lambda \rightarrow \infty,$$

where $N(\lambda) = \#\{\lambda_i \leq \lambda\}$ is the eigenvalue counting function for the Laplacian Δ on M .

Spectral geometry: can one hear the shape of a drum?

- ▶ **Weyl's law:** for a compact Riemannian manifold M

$$N(\lambda) \sim \frac{\omega_n \text{Vol}(M)}{(2\pi)^n} \lambda^{\frac{n}{2}} \quad \lambda \rightarrow \infty,$$

where $N(\lambda) = \#\{\lambda_i \leq \lambda\}$ is the eigenvalue counting function for the Laplacian Δ on M .

- ▶ A better way to think of Weyl's law: quantize the classical Hamiltonian $h(x, p) = \frac{p^2}{2m} + V(x)$, to the quantum Hamiltonian $H = -\frac{\hbar^2}{2m} \Delta + V(x)$. Then

$$N(a \leq \lambda \leq b) = \frac{1}{(2\pi\hbar)^d} \text{Vol} \{a \leq h \leq b\} + o(\hbar^{-d})$$

(Physics proof: by Heisenberg uncertainty relation, each quantum state occupies a volume of $\sim (2\pi\hbar)^d$ in phase space. quantized energy levels are approximated by phase space volumes; Bohr's correspondence principle; semiclassical approximation)

- ▶ Weyl's law: One can hear the volume and dimension of a manifold. We shall see one can hear the volume and scalar curvature of curved noncommutative tori too.

Beyond Weyl's law

- ▶ (M, g) = closed Riemannian manifold. Laplacian on forms

$$\Delta = (d + d^*)^2 : \Omega^p(M) \rightarrow \Omega^p(M),$$

has pure point spectrum:

$$0 \leq \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$$

- ▶ Fact: Dimension, volume, total scalar curvature, Betti numbers, and hence the Euler characteristic of M are fully determined by the spectrum of Δ (on all p -forms).

Heat trace asymptotics

- ▶ $N(\lambda) = \text{Tr } P_\lambda$ is too brutal. Mollify it by a smoothing operator like $\text{Tr}(e^{-t\Delta})$ and use Tauberian theorems to obtain information about $N(\lambda)$.
- ▶ $k(t, x, y) = \text{kernel of } e^{-t\Delta}$. Asymptotic expansion near $t = 0$:

$$k(t, x, x) \sim \frac{1}{(4\pi t)^{m/2}} (a_0(x, \Delta) + a_1(x, \Delta)t + a_2(x, \Delta)t^2 + \dots)$$

- ▶ $a_i(x, \Delta)$, Seeley-De Witt-Gilkey coefficients.

- Theorem: $a_i(x, \Delta)$ are universal polynomials in the curvature tensor $R = R_{jkl}^1$ and its covariant derivatives:

$$a_0(x, \Delta) = 1 \quad \text{Weyl's law}$$

$$a_1(x, \Delta) = \frac{1}{6}S(x) \quad \text{scalar curvature}$$

$$a_2(x, \Delta) = \frac{1}{360}(2|R(x)|^2 - 2|\text{Ric}(x)|^2 + 5|S(x)|^2)$$

$$a_3(x, \Delta) = \dots\dots\dots$$

Noncommutative Local Invariants

- ▶ Local geometric invariants such as scalar curvature of (A, \mathcal{H}, D) are detected by the high frequency behavior of the spectrum of D and the action of A via heat kernel asymptotic expansions of the form

$$\text{Trace} (a e^{-tD^2}) \sim \sum_{j=0}^{\infty} a_j(a, D^2) t^{(-n+j)/2}, \quad t \searrow 0, \quad a \in A.$$

Example: Gauss-Bonnet

- ▶ For surfaces

$$\chi(\Sigma) = \frac{1}{2\pi} \int_{\Sigma} K dA$$

- ▶ Spectral zeta function: Let $\lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \dots$ be the eigenvalues of Δ ,
and

$$\zeta_{\Delta}(s) = \sum \lambda_j^{-s}, \quad \Re(s) > 1.$$

It has a meromorphic extension to \mathbb{C} with a simple pole at $s = \frac{1}{2}$. G-B is equivalent to

$$\zeta_{\Delta}(s) + 1 = 0$$

Curved noncommutative tori

- ▶ A_θ : universal C^* -algebra generated by unitaries U and V

$$VU = e^{2\pi i\theta}UV.$$

- ▶ Smooth structure:

$$A_\theta^\infty = \left\{ \sum_{m,n \in \mathbb{Z}} a_{m,n} U^m V^n : (a_{m,n}) \in \mathcal{S}(\mathbb{Z}^2) \right\}.$$

- ▶ Derivations $\delta_1, \delta_2 : A_\theta^\infty \rightarrow A_\theta^\infty$

$$\delta_1(U) = U, \quad \delta_1(V) = 0, \quad \delta_2(U) = 0, \quad \delta_2(V) = V,$$

- ▶ Canonical trace $\varphi_0 : A_\theta \rightarrow \mathbb{C}$

Complex structure on A_θ

- ▶ Fix $\tau = \tau_1 + i\tau_2$, $\tau_2 = \Im(\tau) > 0$, and define the Dolbeault operators

$$\partial := \delta_1 + \tau\delta_2, \quad \partial^* := \delta_1 + \bar{\tau}\delta_2.$$

- ▶ Let $\mathcal{H}_0 = L^2(A_\theta)$ = GNS completion of A_θ w.r.t. φ_0 .
- ▶ $\mathcal{H}^{(1,0)}$ = Hilbert space of $(1,0)$ -forms: completion of finite sums $\sum a\partial b$, $a, b \in A_\theta^\infty$, under

$$\langle a\partial b, a'\partial b' \rangle := \varphi_0((a'\partial b')^* a\partial b).$$

- ▶ ∂^* is the formal adjoint of $\partial : \mathcal{H}_0 \rightarrow \mathcal{H}^{(1,0)}$.

- ▶ Flat Dolbeault Laplacian:

$$\Delta = \partial^* \partial = \delta_1^2 + 2\tau_1 \delta_1 \delta_2 + |\tau|^2 \delta_2^2.$$

For $\tau = i$, we get

$$\Delta = \delta_1^2 + \delta_2^2.$$

Conformal perturbation of the metric

- ▶ Fix a Weyl factor: $h = h^* \in A_\theta^\infty$. Replace φ_0 by

$$\varphi(a) = \varphi_0(a e^{-h}).$$

- ▶ φ is a KMS state

$$\varphi(ab) = \varphi(b \Delta(a)),$$

with modular automorphism

$$\Delta(a) = \sigma_i(a) = e^{-h} a e^h,$$

and modular group

$$\sigma_t(a) = e^{ith} a e^{-ith}.$$

- ▶ Warning: \triangle and Δ are very different operators!

Curved Laplacian

- ▶ Hilbert space $\mathcal{H}_\varphi = GNS$ completion of A_θ under

$$\varphi(a) = \varphi_0(a e^{-h}).$$

- ▶ Let $\partial_\varphi = \delta_1 + \tau\delta_2 : \mathcal{H}_\varphi \rightarrow \mathcal{H}^{(1,0)}$. It has an adjoint

$$\partial_\varphi^* = R_{k^2}\partial^* : \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}_\varphi.$$

- ▶ Curved Laplacian

$$\Delta' = \partial_\varphi^* \partial_\varphi : \mathcal{H}_\varphi \rightarrow \mathcal{H}_\varphi.$$

A Spectral Triple $(A_\theta^\infty, \mathcal{H}, D)$

$$\mathcal{H} := \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)},$$

$$a \mapsto \begin{pmatrix} a & 0 \\ 0 & a \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H},$$

$$D := \begin{pmatrix} 0 & \partial_\varphi^* \\ \partial_\varphi & 0 \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H},$$

$$\partial_\varphi := \partial = \delta_1 + \bar{\tau}\delta_2 : \mathcal{H}_\varphi \rightarrow \mathcal{H}^{(1,0)}.$$

Anti-Unitary Equivalence of the Laplacians

$$D^2 = \begin{pmatrix} \partial_\varphi^* \partial_\varphi & 0 \\ 0 & \partial_\varphi \partial_\varphi^* \end{pmatrix} : \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)}.$$

Lemma: Let

$$k = e^{h/2}.$$

We have

$$\begin{aligned} \partial_\varphi^* \partial_\varphi : \mathcal{H}_\varphi &\rightarrow \mathcal{H}_\varphi & \sim & k \bar{\partial} \partial k : \mathcal{H}_0 \rightarrow \mathcal{H}_0, \\ \partial_\varphi \partial_\varphi^* : \mathcal{H}^{(1,0)} &\rightarrow \mathcal{H}^{(1,0)} & \sim & \bar{\partial} k^2 \partial : \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}^{(1,0)}. \end{aligned}$$

(The Tomita anti-unitary map J is used.)

**Conformal Geometry of \mathbb{T}_θ^2 with $\tau = i$
(Cohen-Connes, late 80's)**

Let

$\lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \dots$ be the eigenvalues of $\partial_\varphi^* \partial_\varphi$,

and

$$\zeta(s) = \sum \lambda_j^{-s}, \quad \Re(s) > 1.$$

Then

$$\zeta(0) + 1 =$$

$$\varphi(f(\Delta)(\delta_1(e^{h/2})) \delta_1(e^{h/2})) + \varphi(f(\Delta)(\delta_2(e^{h/2})) \delta_2(e^{h/2})),$$

where

$$f(u) = \frac{1}{6}u^{-1/2} - \frac{1}{3} + \mathcal{L}_1(u) - 2(1+u^{1/2})\mathcal{L}_2(u) + (1+u^{1/2})^2\mathcal{L}_3(u),$$

$$\mathcal{L}_m(u) = (-1)^m (u-1)^{-(m+1)} \left(\log u - \sum_{j=1}^m (-1)^{j+1} \frac{(u-1)^j}{j} \right).$$

The Gauss-Bonnet theorem for \mathbb{T}_θ^2

Theorem. (Connes-Tretkoff; Fathizadeh-Kh.) For any $\theta \in \mathbb{R}$, complex parameter $\tau \in \mathbb{C} \setminus \mathbb{R}$ and Weyl conformal factor $e^h, h = h^* \in A_\theta^\infty$, we have

$$\zeta(0) + 1 = 0.$$

Final Part of the Proof

$$\zeta(0) + 1 =$$

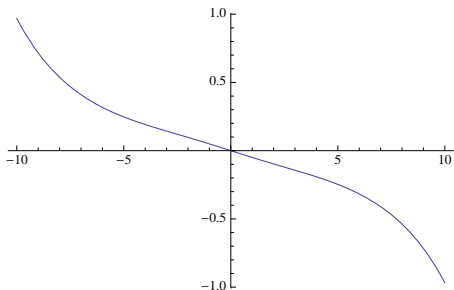
$$\begin{aligned} & \frac{2\pi}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_1 \left(\frac{h}{2} \right) \right) \delta_1 \left(\frac{h}{2} \right) \right) + \frac{2\pi|\tau|^2}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_2 \left(\frac{h}{2} \right) \right) \delta_2 \left(\frac{h}{2} \right) \right) \\ & + \frac{2\pi\Re(\tau)}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_1 \left(\frac{h}{2} \right) \right) \delta_2 \left(\frac{h}{2} \right) \right) + \frac{2\pi\Re(\tau)}{\Im(\tau)} \varphi_0 \left(K(\nabla) \left(\delta_2 \left(\frac{h}{2} \right) \right) \delta_1 \left(\frac{h}{2} \right) \right), \end{aligned}$$

where

$$K(x) = - \frac{(3x - 3 \sinh \left(\frac{x}{2} \right) - 3 \sinh(x) + \sinh \left(\frac{3x}{2} \right)) \operatorname{csch}^2 \left(\frac{x}{2} \right)}{3x^2}$$

is an odd entire function, and $\nabla = \log \Delta$.

$$K(x) = -\frac{x}{20} + \frac{x^3}{2240} - \frac{23x^5}{806400} + O(x^6).$$



Scalar curvature for A_θ

- ▶ The scalar curvature of the curved nc torus $(\mathbb{T}_\theta^2, \tau, k)$ is the unique element $R \in A_\theta^\infty$ satisfying

$$\text{Trace}(a\Delta^{-s})|_{s=0} + \text{Trace}(aP) = \mathfrak{t}(aR), \quad \forall a \in A_\theta^\infty,$$

where P is the projection onto the kernel of Δ .

- ▶ In practice this is done by finding an asymptotic expansion for the kernel of the operator $ae^{-t\Delta}$,

$$\text{Trace}(ae^{-tD^2}) \sim \sum_{n \geq 0} B_n(a, D^2) t^{\frac{n-2}{2}}, \quad a \in A_\theta^\infty.$$

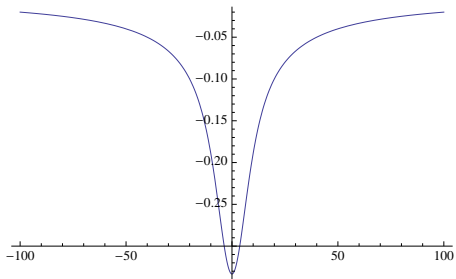
using Connes' **pseudodifferential calculus** for nc tori. A good pseudo diff calculus for general nc spaces is still illusive.

Final Formula for the Scalar Curvature of \mathbb{T}_θ^2

Theorem. (Connes-Moscovici; Fathizadeh-Kh.) Up to an overall factor of $\frac{-\pi}{\Im(\tau)}$, R is equal to

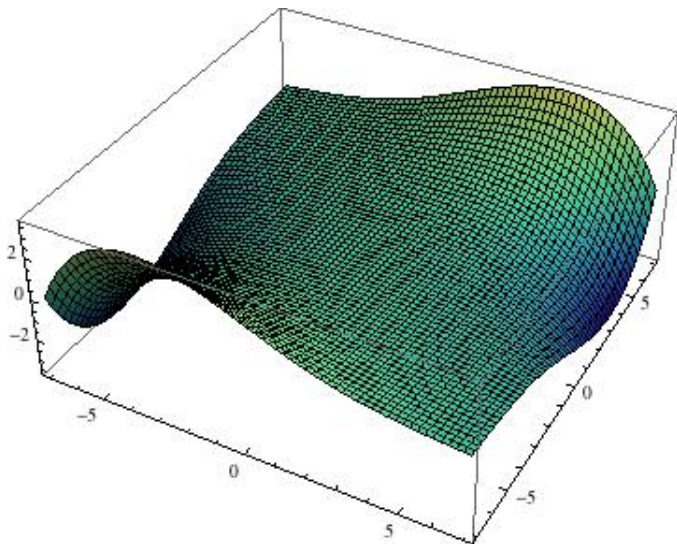
$$\begin{aligned} & R_1(\nabla)\left(\delta_1^2\left(\frac{\hbar}{2}\right) + 2\tau_1\delta_1\delta_2\left(\frac{\hbar}{2}\right) + |\tau|^2\delta_2^2\left(\frac{\hbar}{2}\right)\right) \\ & + R_2(\nabla, \nabla)\left(\delta_1\left(\frac{\hbar}{2}\right)^2 + |\tau|^2\delta_2\left(\frac{\hbar}{2}\right)^2 + \Re(\tau)\left\{\delta_1\left(\frac{\hbar}{2}\right), \delta_2\left(\frac{\hbar}{2}\right)\right\}\right) \\ & + iW(\nabla, \nabla)\left(\Im(\tau)\left[\delta_1\left(\frac{\hbar}{2}\right), \delta_2\left(\frac{\hbar}{2}\right)\right]\right). \end{aligned}$$

$$R_1(x) = \frac{\frac{1}{2} - \frac{\sinh(x/2)}{x}}{\sinh^2(x/4)}.$$



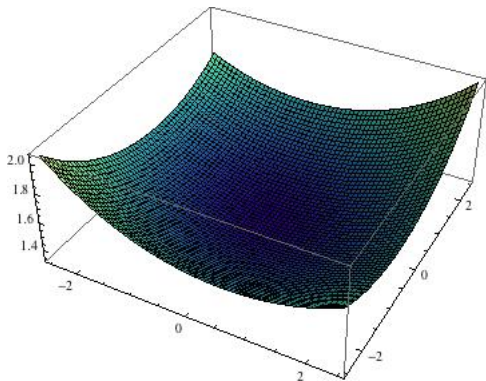
$$R_2(s, t) =$$

$$\frac{(1 + \cosh((s+t)/2))(-t(s+t) \cosh s + s(s+t) \cosh t - (s-t)(s+t + \sinh s + \sinh t - \sinh(s+t)))}{st(s+t) \sinh(s/2) \sinh(t/2) \sinh^2((s+t)/2)}$$



$$W(s, t) =$$

$$\frac{(-s - t + t \cosh s + s \cosh t + \sinh s + \sinh t - \sinh(s + t))}{st \sinh(s/2) \sinh(t/2) \sinh((s + t)/2)}.$$



Noncommutative 4-Torus \mathbb{T}_θ^4

- ▶ Complex Structure on \mathbb{T}_θ^4

$$\begin{aligned}\partial &= \partial_1 \oplus \partial_2, & \bar{\partial} &= \bar{\partial}_1 \oplus \bar{\partial}_2, \\ \partial_1 &= \frac{1}{2} (\delta_1 - i\delta_3), & \partial_2 &= \frac{1}{2} (\delta_2 - i\delta_4), \\ \bar{\partial}_1 &= \frac{1}{2} (\delta_1 + i\delta_3), & \bar{\partial}_2 &= \frac{1}{2} (\delta_2 + i\delta_4).\end{aligned}$$

Conformal perturbation of the metric

Let $h = h^* \in C^\infty(\mathbb{T}_\theta^4)$ and replace the trace φ_0 by

$$\varphi : C(\mathbb{T}_\theta^4) \rightarrow \mathbb{C},$$

$$\varphi(a) := \varphi_0(a e^{-2h}), \quad a \in C(\mathbb{T}_\theta^4).$$

φ is a KMS state with the modular group

$$\sigma_t(a) = e^{2ith} a e^{-2ith}, \quad a \in C(\mathbb{T}_\theta^4),$$

and the modular automorphism

$$\Delta(a) := \sigma_i(a) = e^{-2h} a e^{2h}, \quad a \in C(\mathbb{T}_\theta^4).$$

$$\varphi(ab) = \varphi(b \Delta(a)), \quad a, b \in C(\mathbb{T}_\theta^4).$$

Perturbed Laplacian on \mathbb{T}_θ^4

$$d = \partial \oplus \bar{\partial} : \mathcal{H}_\varphi \rightarrow \mathcal{H}_\varphi^{(1,0)} \oplus \mathcal{H}_\varphi^{(0,1)},$$

$$\Delta_\varphi := d^*d.$$

Remark. If $h = 0$ then $\varphi = \varphi_0$ and

$$\Delta_{\varphi_0} = \delta_1^2 + \delta_2^2 + \delta_3^2 + \delta_4^2 = \partial^*\partial$$

(the underlying manifold is Kähler).

Scalar Curvature for \mathbb{T}_θ^4

It is the unique element $R \in C^\infty(\mathbb{T}_\theta^4)$ such that

$$\operatorname{Res}_{s=1} \zeta_a(s) = \varphi_0(a R), \quad a \in C^\infty(\mathbb{T}_\theta^4),$$

where

$$\zeta_a(s) := \operatorname{Trace}(a \Delta_\varphi^{-s}), \quad \Re(s) \gg 0.$$

Final Formula for the Scalar Curvature of \mathbb{T}_θ^4

Theorem. (Fathizadeh-Kh.) We have

$$R = e^{-h} k(\nabla) \left(\sum_{i=1}^4 \delta_i^2(h) \right) + e^{-h} H(\nabla, \nabla) \left(\sum_{i=1}^4 \delta_i(h)^2 \right),$$

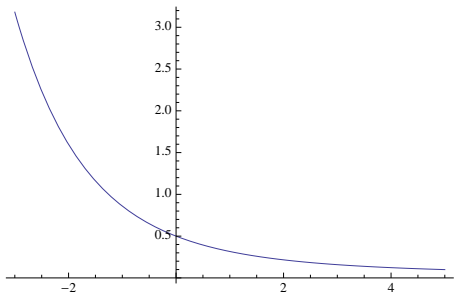
where

$$\nabla(a) = [-h, a], \quad a \in C(\mathbb{T}_\theta^4),$$

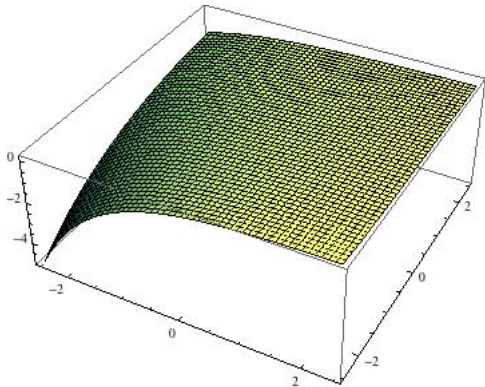
$$k(s) = \frac{1 - e^{-s}}{2s},$$

$$H(s, t) = -\frac{e^{-s-t} ((-e^s - 3) s (e^t - 1) + (e^s - 1) (3e^t + 1) t)}{4 s t (s + t)}.$$

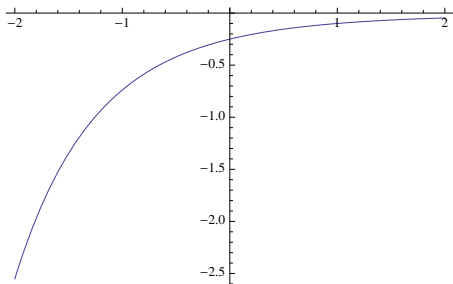
$$k(s) = \frac{1}{2} - \frac{s}{4} + \frac{s^2}{12} - \frac{s^3}{48} + \frac{s^4}{240} - \frac{s^5}{1440} + O(s^6).$$



$$H(s, t) = \left(-\frac{1}{4} + \frac{t}{24} + O(t^3) \right) + s \left(\frac{5}{24} - \frac{t}{16} + \frac{t^2}{80} + O(t^3) \right) \\ + s^2 \left(-\frac{1}{12} + \frac{7t}{240} - \frac{t^2}{144} + O(t^3) \right) + O(s^3).$$



$$\begin{aligned} H(s, s) &= -\frac{e^{-2s}(e^s - 1)^2}{4s^2} \\ &= -\frac{1}{4} + \frac{s}{4} - \frac{7s^2}{48} + \frac{s^3}{16} - \frac{31s^4}{1440} + \frac{s^5}{160} + O(s^6). \end{aligned}$$



$$\begin{aligned}
 G(s) &:= H(s, -s) = \frac{-4s - 3e^{-s} + e^s + 2}{4s^2} \\
 &= -\frac{1}{4} + \frac{s}{6} - \frac{s^2}{48} + \frac{s^3}{120} - \frac{s^4}{1440} + \frac{s^5}{5040} + O(s^6).
 \end{aligned}$$

