



# Measuring and controlling quantum transport of heat in trapped-ion crystals

August 14th 2013

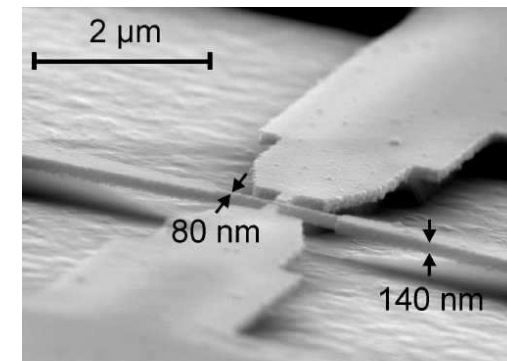
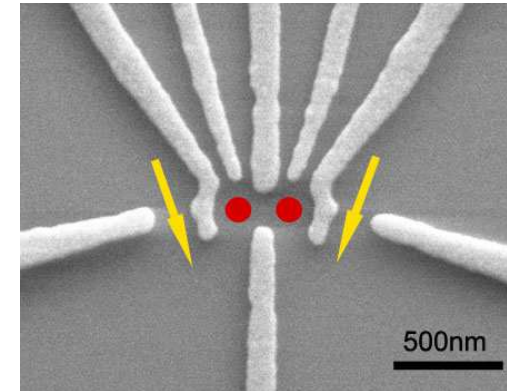
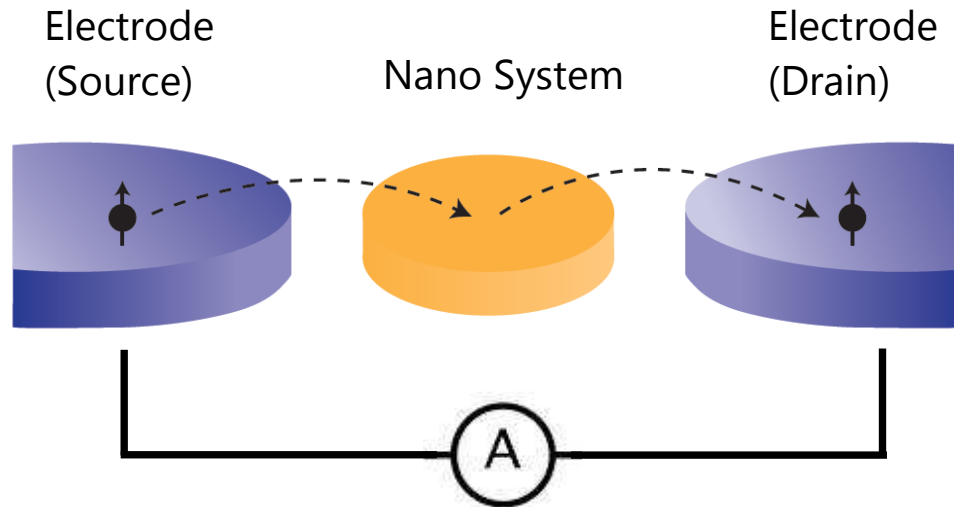
Alejandro Bermudez

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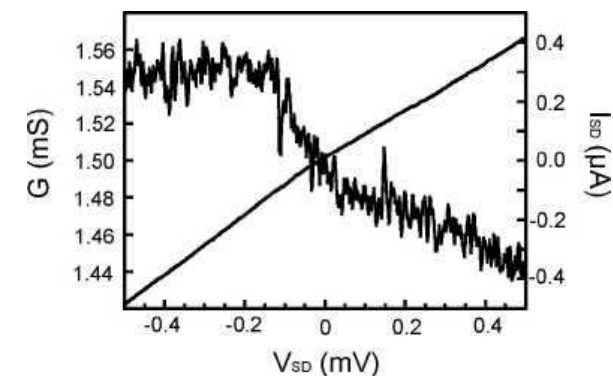


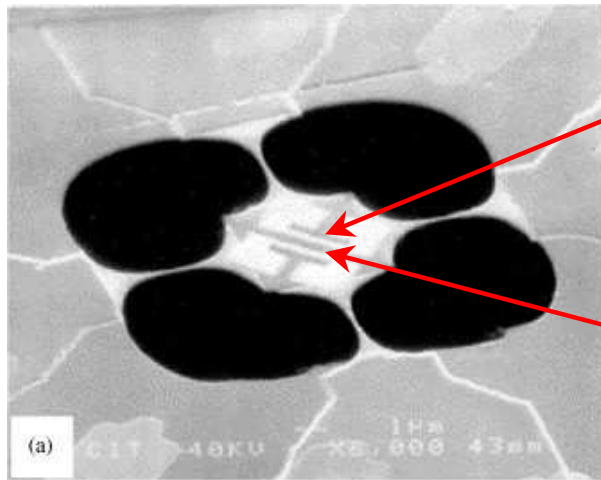
**Alexander von Humboldt**  
Stiftung/Foundation



## State-of-the-art in electronic transport

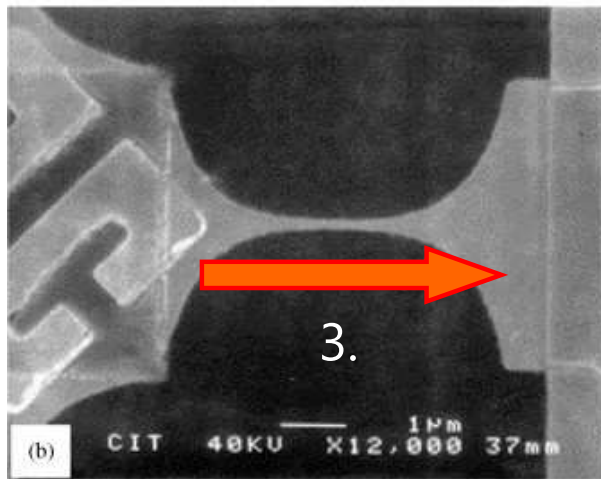
- Control current via bias & gate voltage
- Measure electronic currents & fluctuations using amperemeter (pA)
- Charge gives a handle on electrons



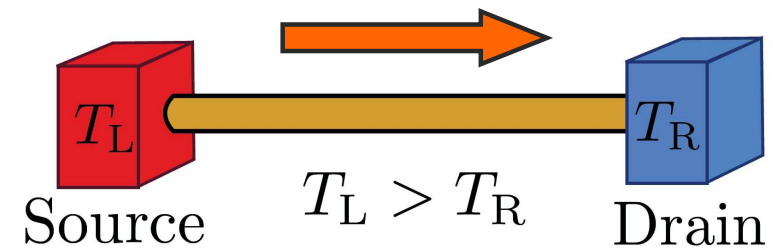
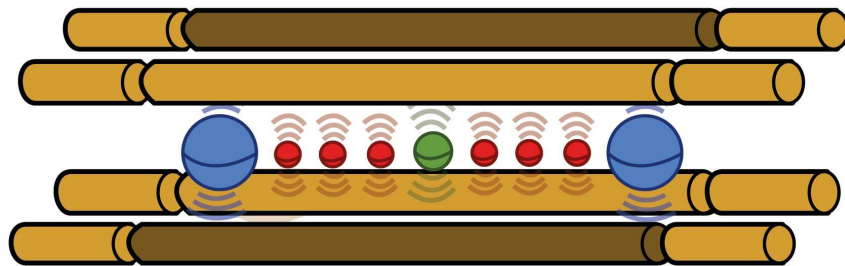


Measuring heat currents in nanoscale systems

1. Controlled heating
2. Temperature measurement
3. Infer heat current from 1. & 2.

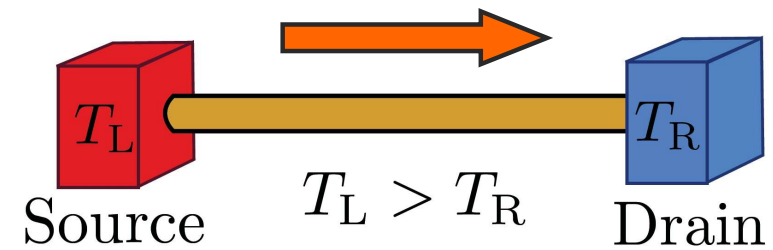
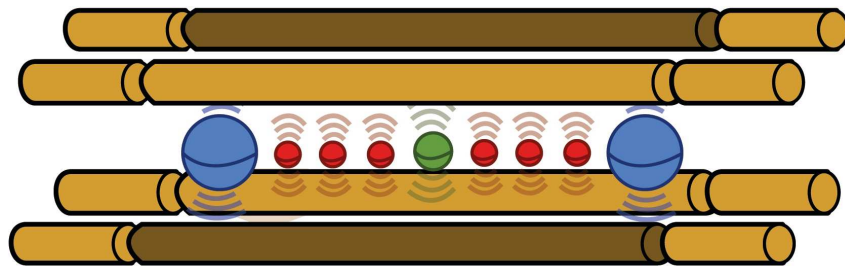


- No 'ampere meter' for heat currents
- No charge for heat
- Difficult to control heat currents



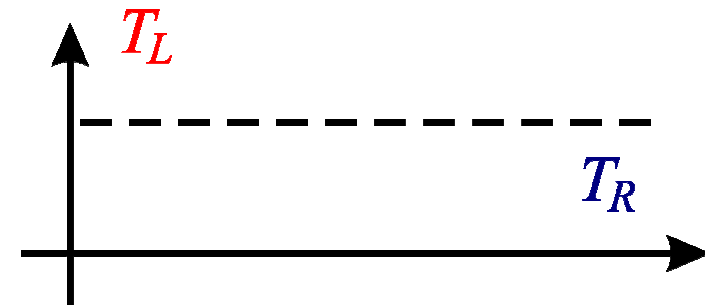
Ion crystal is made of

- **Bulk ions**  
Transverse vibrations (vibrons)  
‣ coupled harmonic oscillators
- **Reservoir ions (source & drain)**  
Cooled to different temperatures  
‣ induce thermal currents
- **Probe ions**  
Vibrons coupled to internal states  
‣ 'ampere meter' for heat current



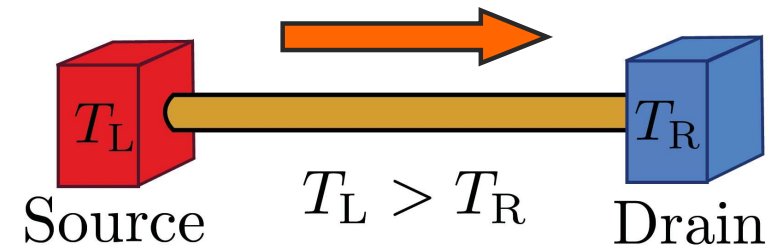
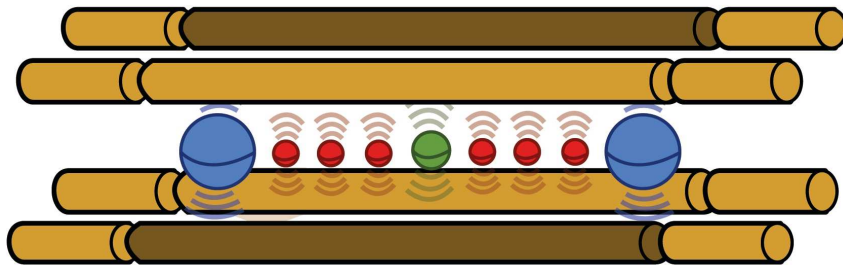
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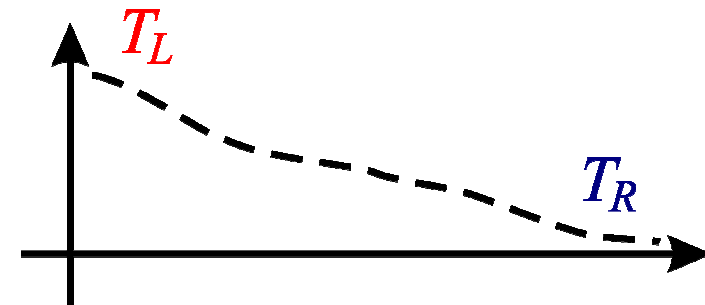
Transition from ballistic to diffusive transport

- Onset of Fourier's law



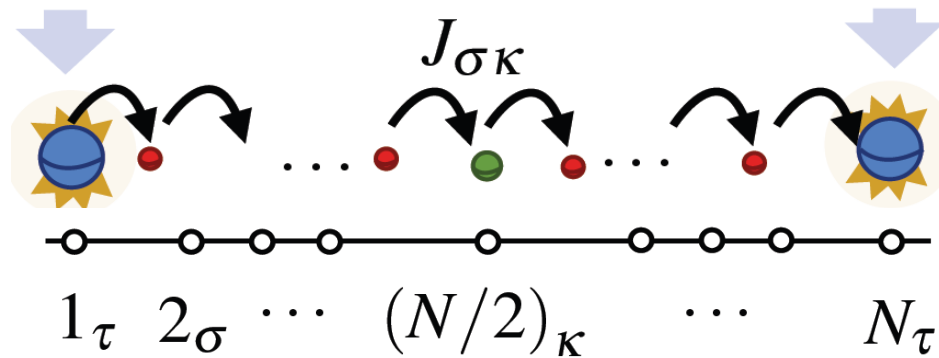
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Transition from ballistic to diffusive transport

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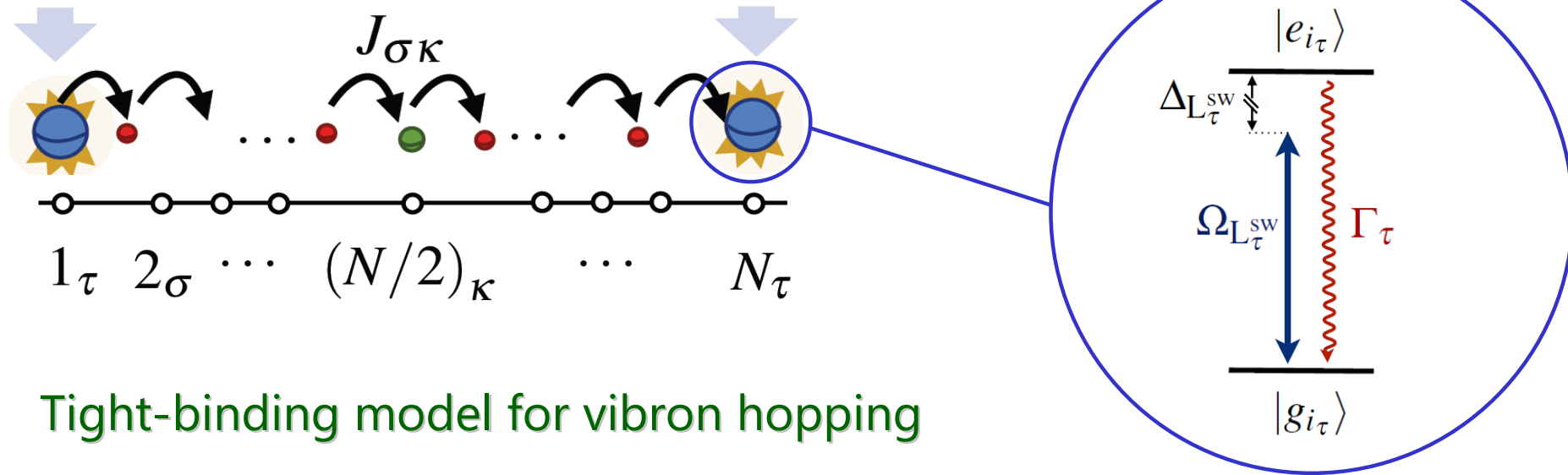
Additional functionality  
by using internal states

- slow dynamics
- fully controllable

Tight-binding model for vibron hopping

$$H_{\text{tb}} = \sum_{\alpha, i_\alpha} \omega_{i_\alpha} a_{i_\alpha}^\dagger a_{i_\alpha} + \sum_{\alpha, \beta} \sum_{i_\alpha \neq j_\beta} (J_{i_\alpha j_\beta} a_{i_\alpha}^\dagger a_{j_\beta} + \text{H.c.})$$

- small oscillations above ground state
- coupling via dipole-dipole interaction  $J \sim 1/d^3$
- energy (heat) transport by vibron hopping



## Tight-binding model for vibron hopping

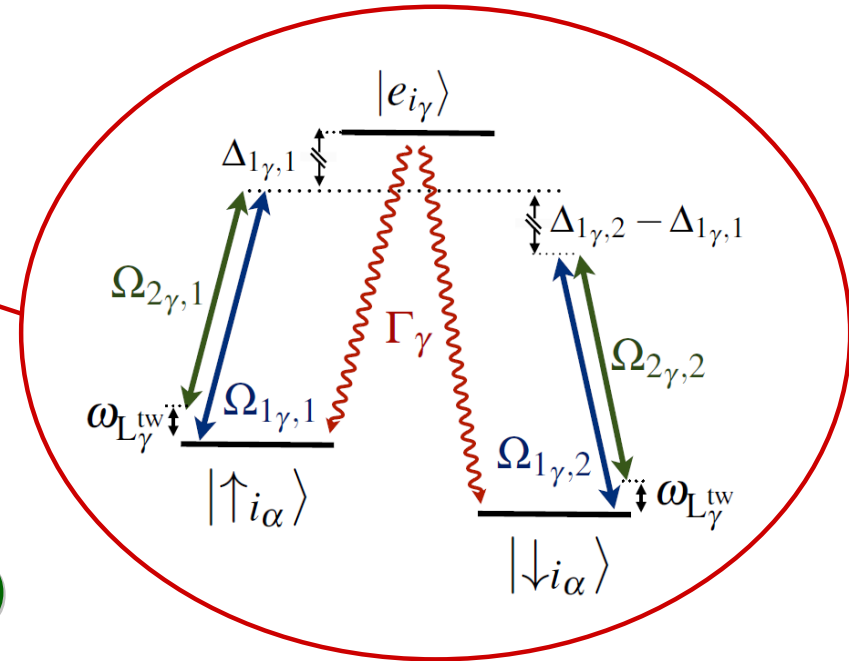
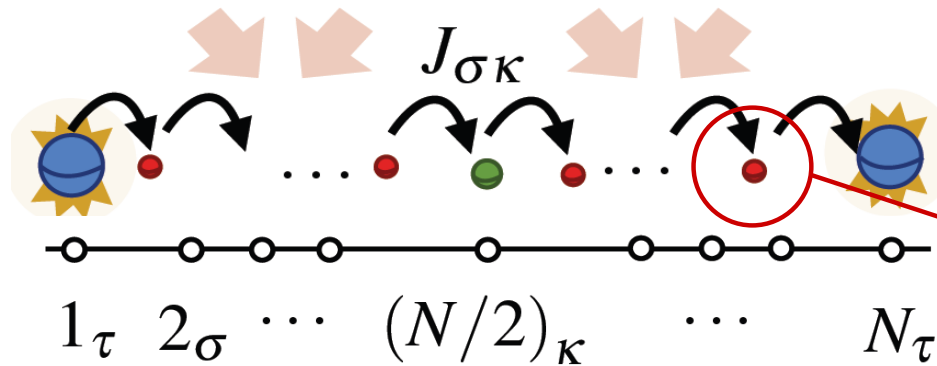
$$H_{\text{tb}} = \sum_{\alpha, i_{\alpha}} \omega_{i_{\alpha}} a_{i_{\alpha}}^{\dagger} a_{i_{\alpha}} + \sum_{\alpha, \beta} \sum_{i_{\alpha} \neq j_{\beta}} (J_{i_{\alpha} j_{\beta}} a_{i_{\alpha}}^{\dagger} a_{j_{\beta}} + \text{H.c.})$$

## Doppler cooling of edge ions at rate

$$\mathcal{D}^{i_{\alpha}}(\mu) = \mathcal{D}[\Lambda_{i_{\alpha}}^{+}, a_{i_{\alpha}}^{\dagger}, a_{i_{\alpha}}](\mu) + \mathcal{D}[\Lambda_{i_{\alpha}}^{-}, a_{i_{\alpha}}, a_{i_{\alpha}}^{\dagger}](\mu)$$

- ▶ Effective cooling rate  $\gamma_{i_{\alpha}} = \text{Re}\{(\Lambda_{i_{\alpha}}^{-})^* - \Lambda_{i_{\alpha}}^{+}\}$
- ▶ Doppler cooling much faster than hopping  $\gamma_{i_{\alpha}} \gg J_{i_{\alpha}, j_{\beta}}$





Couple internal states to vibrons (heat)

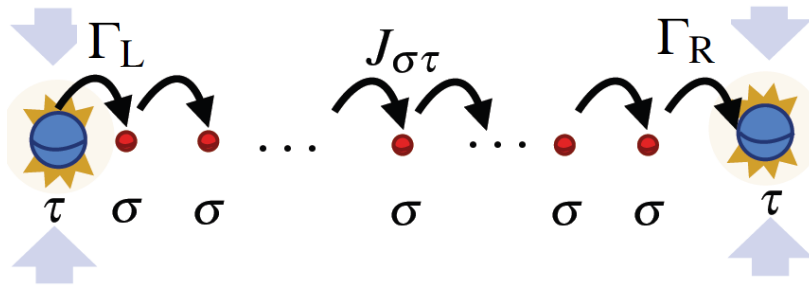
$$H_{SV}^{i_\alpha}(t) = \frac{1}{2}(\Delta\omega_\alpha^+ + \Delta\omega_\alpha^- \sigma_{i_\alpha}^z) \cos(\nu_\alpha t - \varphi_\alpha) a_{i_\alpha}^\dagger a_{i_\alpha}$$

- ▷ internal states (spins)  $|\uparrow\rangle, |\downarrow\rangle$
- ▷ two-photon transition

(1)  $H_{SV}^{i_\alpha}(t) = \frac{1}{2}\Delta\omega_\alpha^+ \cos(\nu_\alpha t) a_{i_\alpha}^\dagger a_{i_\alpha}$       ▷ Photon-assisted tunneling

(2)  $H_{SV}^{i_\alpha} = \frac{1}{2}\Delta\omega_\alpha^- \sigma_{i_\alpha}^z a_{i_\alpha}^\dagger a_{i_\alpha}$       ▷ Probing & Disorder

What physics can we do?



Assume  $\gamma_{l\tau} \gg J_{i\alpha j\beta}$  and project dynamics onto state  $\mu_{1\tau}^{\text{th}} \otimes \mu_{\text{bulk}} \otimes \mu_{N\tau}^{\text{th}}$

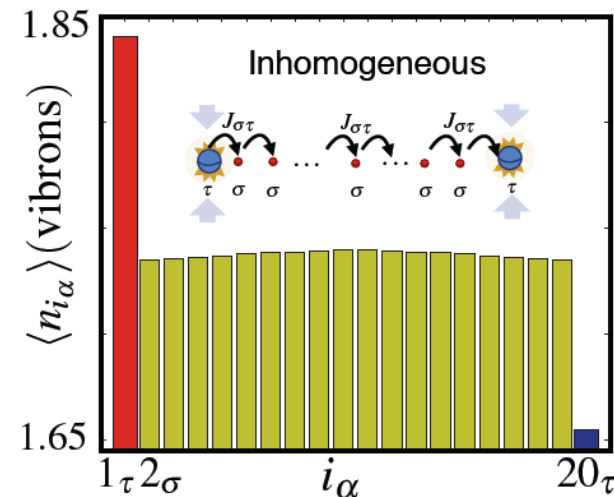
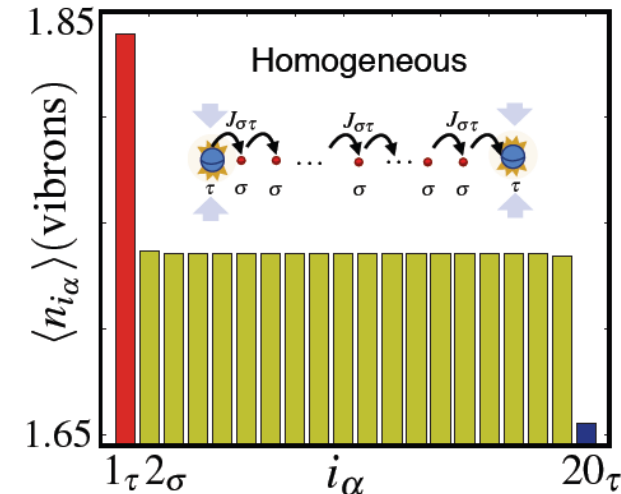
$$\langle n_{i\alpha} \rangle_{\text{ss}} = \frac{\Gamma_L \bar{n}_L + \Gamma_R \bar{n}_R}{\Gamma_L + \Gamma_R}$$

$$\langle I_{\rightarrow i\alpha}^{\text{vib}} \rangle_{\text{ss}} = \frac{\Gamma_L \Gamma_R}{\Gamma_L + \Gamma_R} (\bar{n}_L - \bar{n}_R)$$

$$\Gamma_{i\alpha, j\beta}^{l\tau} = 2\pi J_{i\alpha, l\tau} \rho_{l\tau}(\omega_{i\alpha}) J_{l\tau, j\beta}$$

Ballistic transport of vibrons across TQC

## Vibron occupations in TQC



Strong spin-vibron coupling

$$H_{SV}^{i\alpha}(t) = \frac{1}{2} \Delta\omega_{\alpha}^{-} \sigma_{i\alpha}^z a_{i\alpha}^{\dagger} a_{i\alpha}$$

Spins in superposition

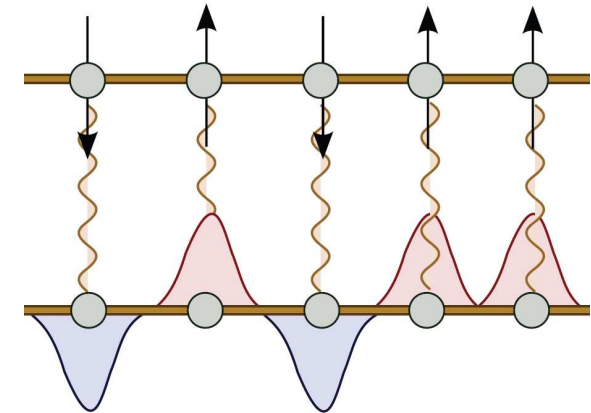
$$|+i_{\alpha}\rangle = (|\uparrow i_{\alpha}\rangle + |\downarrow i_{\alpha}\rangle) / \sqrt{2}$$

Tight-binding model with disorder

$$H_{\text{stb}} = \sum_{\alpha, i_{\alpha}} \epsilon_{i_{\alpha}} a_{i_{\alpha}}^{\dagger} a_{i_{\alpha}} + \sum_{\alpha, \beta} \sum_{i_{\alpha} \neq j_{\beta}} (J_{i_{\alpha} j_{\beta}} a_{i_{\alpha}}^{\dagger} a_{j_{\beta}} + \text{H.c.})$$

Binary diagonal disorder

$$\epsilon_{i_{\alpha}} \in \left\{ \omega_{\alpha} - \frac{1}{2} \Delta\omega_{\alpha}^{-}, \omega_{\alpha} + \frac{1}{2} \Delta\omega_{\alpha}^{-} \right\}$$



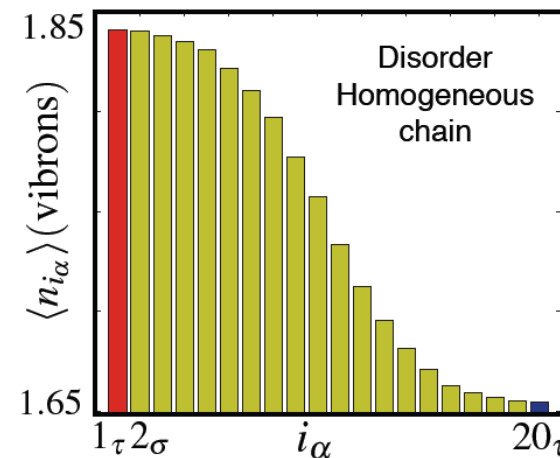
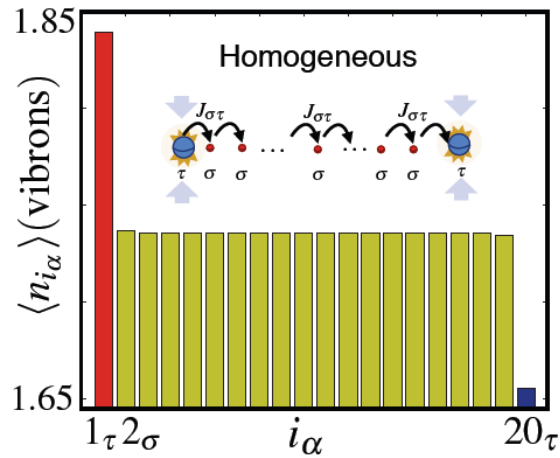
Exploit "quantum parallelism"

B. Paredes, F. Verstraete & J. I. Cirac, PRL **95**, 140501 (2005)

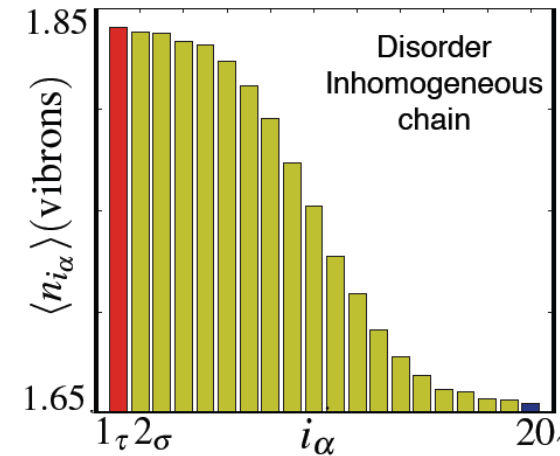
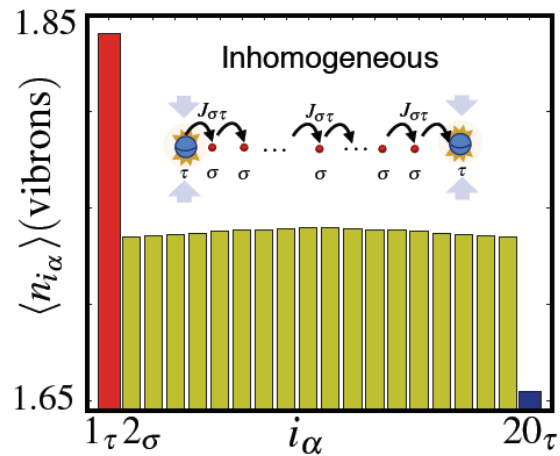
A. Bermudez, M. A. Martin-Delgado & D. Porras, New J. Phys. **12**, 123016 (2010)

## Ballistic

## Diffusive

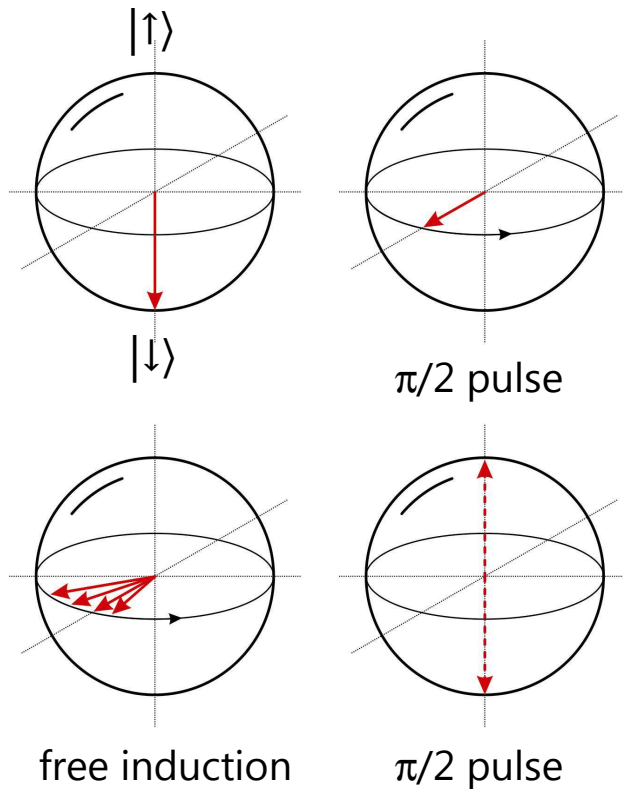


Homogeneous



Inhomogeneous

Clear signature for the onset of Fourier's law



Operator couples weakly to spin

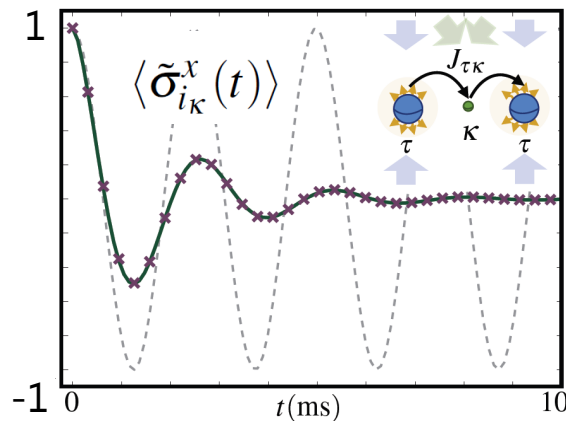
$$H_{SV} = \frac{1}{2} \lambda O_{i\alpha} \sigma_{i\alpha}^z$$

Spin evolution

$$\langle \tilde{\sigma}_{i\alpha}^x(t) \rangle = \cos(\lambda \langle O_{i\alpha} \rangle_{ss} t) e^{-\lambda^2 \text{Re}\{S_{O_{i\alpha} O_{i\alpha}}(0)\} t}$$

$$S_{O_{i\alpha} O_{i\alpha}}(\omega) = \int_0^\infty dt \langle \tilde{O}_{i\alpha}(t) \tilde{O}_{i\alpha}(0) \rangle_{ss} e^{-i\omega t}$$

$$\tilde{O}_{i\alpha} = O_{i\alpha} - \langle O_{i\alpha} \rangle_{ss}$$



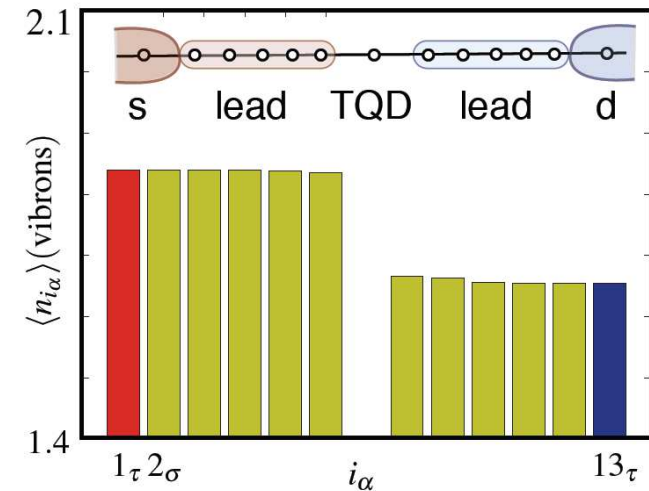
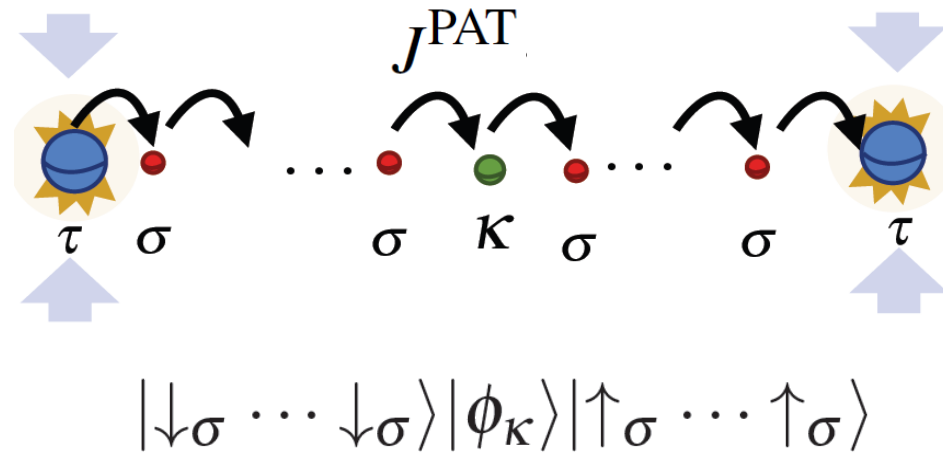
- ▷ Oscillations with frequency  $\sim \langle O \rangle$
- ▷ Damping by fluctuations  $\sim \langle \delta O^2 \rangle$

Measure occupations and thermal currents.

MB and D. Jaksch, New J. Phys. **8**, 87 (2006)

G. B. Lesovik, F. Hassler & G. Blatter, PRL **96**, 106801 (2006)

## Single site & thermal leads



## Photon-assisted tunneling

$$H_{SV}^{i\alpha}(t) = \frac{1}{2} \Delta \omega_{\alpha}^{+} \cos(\nu_{\alpha} t) a_{i_{\alpha}}^{\dagger} a_{i_{\alpha}}$$

$$H_{LKR}^{PAT} = \sum_{i_{\sigma}} (\hat{J}_{i_{\sigma}, p_K}^{PAT} a_{i_{\sigma}}^{\dagger} a_{p_K} + \text{H.c.})$$

- Full control of coupling to leads

Energy mismatch left/right

$$H_{SV}^{i\alpha}(t) = \frac{1}{2} \Delta \omega_{\alpha}^{-} \sigma_{i_{\alpha}}^z a_{i_{\alpha}}^{\dagger} a_{i_{\alpha}}$$

## Same averages for fermions and bosons

$$\langle n_{i\alpha} \rangle_{ss} = \frac{\Gamma_L \bar{n}_L + \Gamma_R \bar{n}_R}{\Gamma_L + \Gamma_R}$$

$$\langle I_{\rightarrow i\alpha}^{\text{vib}} \rangle_{ss} = \frac{\Gamma_L \Gamma_R}{\Gamma_L + \Gamma_R} (\bar{n}_L - \bar{n}_R)$$

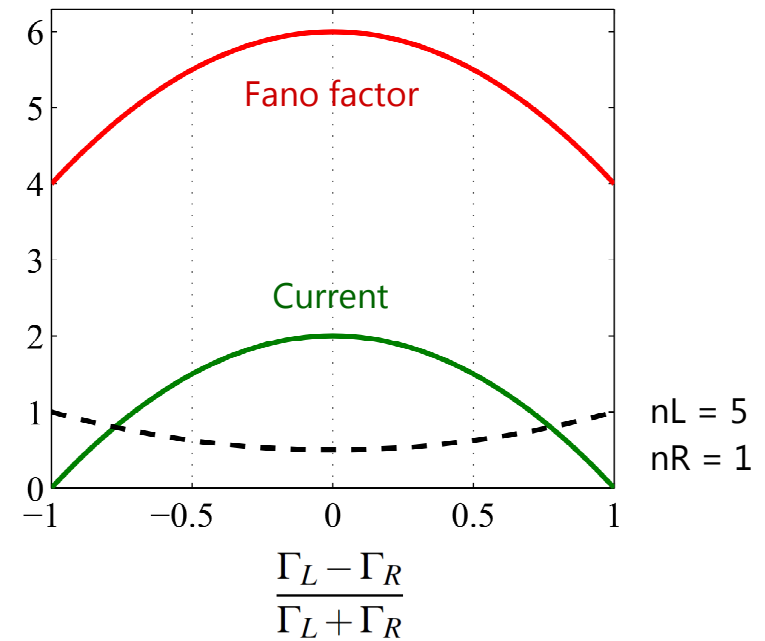
## Fluctuations distinguish bosons/fermions

$$F = \frac{\langle O^2 \rangle - \langle O \rangle^2}{\langle O \rangle} \equiv \frac{\langle \delta O^2 \rangle}{\langle O \rangle}$$

Fano factor

Poissonian  $F = 1$

Mandel  $Q = F - 1$

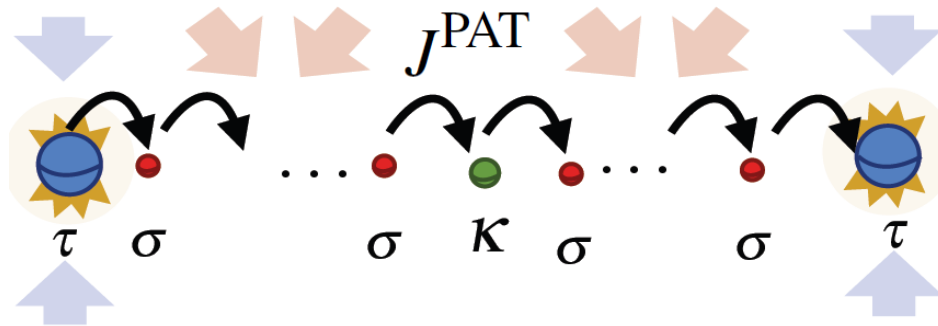


	Fermions	Bosons	Classical
Occupation	$1 - \langle n \rangle$	$1 + \langle n \rangle$	$\langle n \rangle$
Current*	$1 - \frac{1}{2}n_L$	$1 + \frac{1}{2}n_L$	1

\*  $n_R = 0$  and symmetric coupling

Fluctuations reveal bosonic nature of thermal currents.





Thermal currents more resilient to decoherence than electronic currents!

## Single-spin heat switch

- Spin of TQD controls current

$$J_{p_K i_\sigma}^{\text{PAT}}(\sigma_{p_K}^z)|\downarrow_{p_K}\rangle = 0$$

$$J_{p_K i_\sigma}^{\text{PAT}}(\sigma_{p_K}^z)|\uparrow_{p_K}\rangle \neq 0$$

- Superposition of heat current on/off



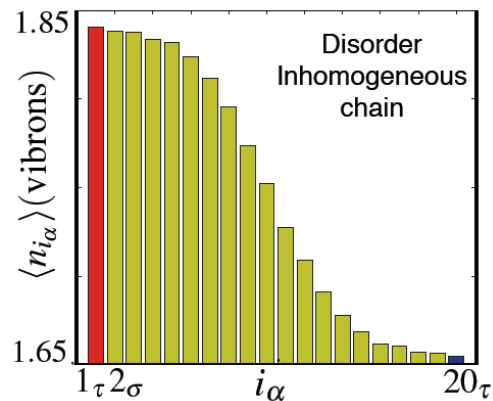
Thermal reservoirs  
Doppler cooling



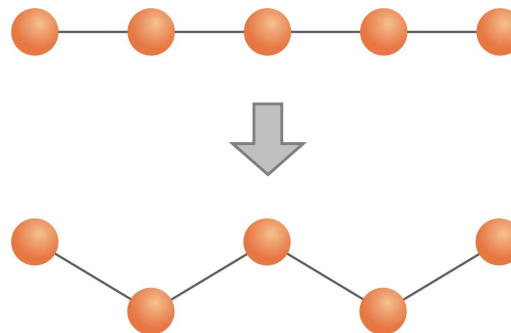
Ampere meter & thermometer  
Spin-vibron coupling



Fourier's law



Zig-zag crossover



Thermal diode  
& quantum dot

