

Propagation of Long-crested Water Waves

Jerry Bona

University of Illinois at Chicago

Collaborators

- **Barbara Boczar**
- **Min Chen**
- **Colette Guillope**
- **David Lannes**
- **Wojtek Romanczyk**
- **Shuming Sun**
- **Bing-Yu Zhang**
- **Hongqiu Chen**
- **Thierry Colin**
- **Angus Jackson**
- **Juan Restepo**
- **Jean-Claude Saut**
- **Ed Thornton**

Examples of Long-crested Waves

















Evidence of the Persistence of Long-crested Waves

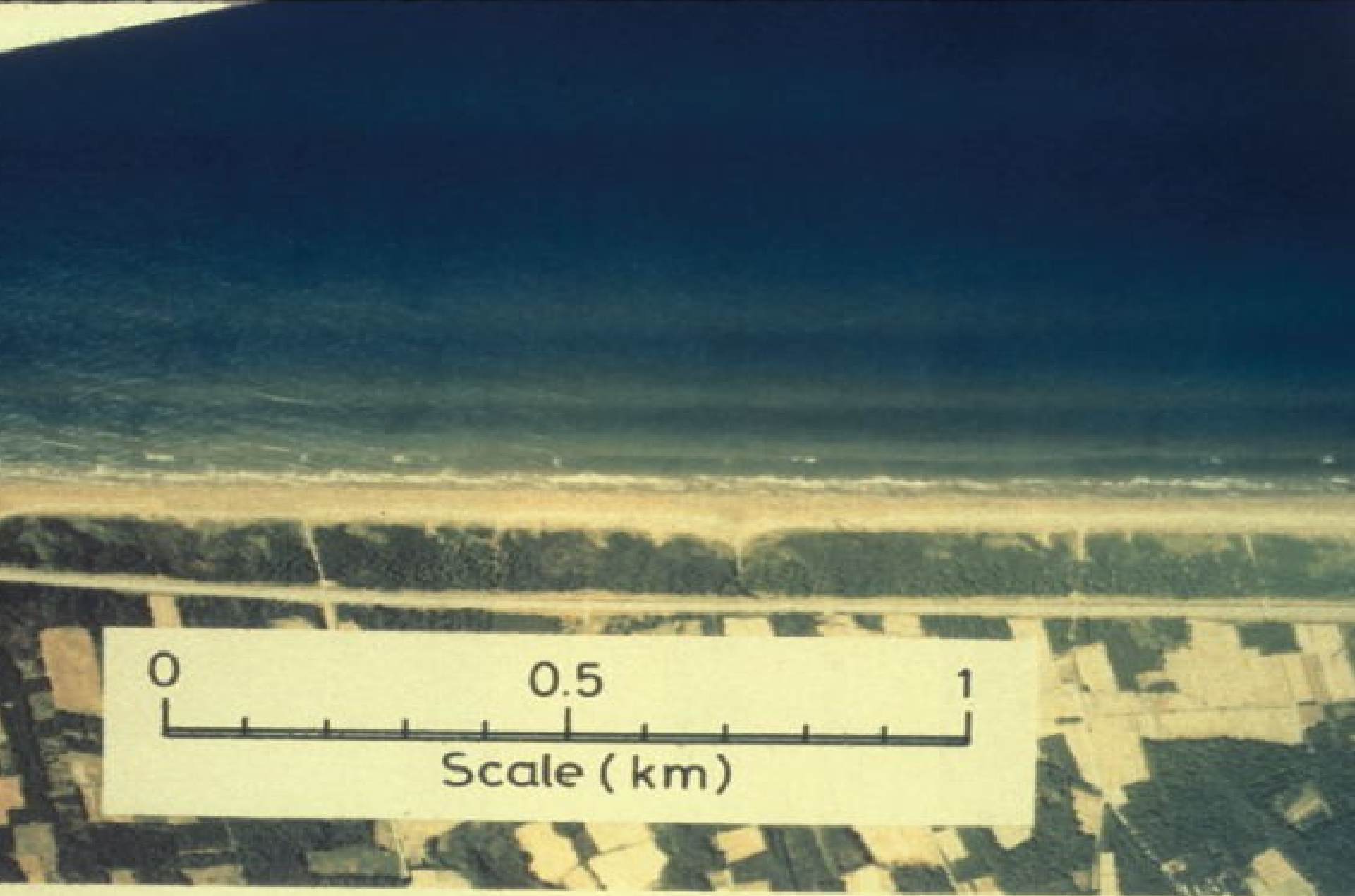












(a) Sept. 6, 1975 (11h. 29m.). Reference point No.42 to No.54.

Photograph 1 Multiple longshore bars at Hakui Beach, Japan taken by the Geographical Survey Institute, Ministry of Construction. Number of reference points

Analysis of an Example Problem

Three-dimensional Euler Equations

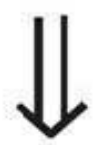
$\alpha \ll 1$
 $\beta \ll 1$
 $S \sim 1$



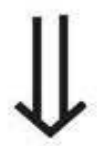
\Rightarrow
y-independent

Two-dimensional Euler Equations

$\alpha \ll 1$
 $\beta \ll 1$
 $S \sim 1$

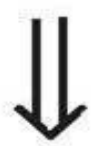


Three-dimensional Boussinesq System



one-way propagation

Two-dimensional Boussinesq system



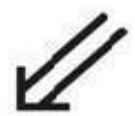
one-way propagation

KP-like equations



wave equation
(hydraulics)

KdV or BBM



$\alpha \ll \ll 1$
 $\beta \ll \ll 1$

The a-b-c-d systems

$$V_t + \nabla \eta + \epsilon \left(\frac{1}{2} \nabla |V|^2 + a \nabla \Delta \eta - b \Delta V_t \right) = 0$$

$$\eta_t + \nabla \cdot V + \epsilon (\nabla \cdot (\eta V) + c \Delta \nabla V - d \Delta \eta_t) = 0$$

where

$$a = \frac{1 - \theta^2}{2} \mu \qquad b = \frac{1 - \theta^2}{2} (1 - \mu)$$
$$c = \frac{\theta^2 - \frac{1}{3}}{2} \lambda \qquad d = \frac{\theta^2 - \frac{1}{3}}{2} (1 - \lambda)$$

and $\theta \in [0, 1]$, $\lambda, \mu \in \mathcal{R}$ present themselves as possible models for the surface waves.

Two-dimensional Versions

(negligible variation in the y-direction)

$$\eta_t + u_x + \epsilon((u\eta)_x + au_{xxx} - b\eta_{xxt}) = 0$$

$$u_t + \eta_x + \epsilon(uu_x + c\eta_{xxx} - du_{xxt}) = 0$$

where

$$a = \frac{1}{2} \left(\theta^2 - \frac{1}{3} \right) \lambda$$
$$b = \frac{1}{2} \left(\theta^2 - \frac{1}{3} \right) (1 - \lambda)$$
$$c = \frac{1}{2} (1 - \theta^2) \mu$$
$$d = \frac{1}{2} (1 - \theta^2) (1 - \mu)$$

$$0 \leq \theta \leq 1, \lambda, \mu \in \mathcal{R}, \quad a + b + c + d = \frac{1}{3}$$

coupled KdV system :

$$\eta_t + u_x + (u\eta)_x + \frac{1}{6}u_{xxx} = 0$$

$$u_t + \eta_x + uu_x + \frac{1}{6}\eta_{xxx} = 0$$

regularized Boussinesq system :

$$\eta_t + u_x + (u\eta)_x - \frac{1}{6}\eta_{xxt} = 0$$

$$u_t + \eta_x + uu_x - \frac{1}{6}u_{xxt} = 0$$

Boussinesq's 'original' system:

$$\eta_t + u_x + (u\eta)_x = 0$$

$$u_t + \eta_x + uu_x - u_{xxt} = 0$$

Bona-Smith system:

$$\eta_t + u_x + (u\eta)_x - b\eta_{xxt} = 0$$

$$u_t + \eta_x + uu_x + c\eta_{xxx} - bu_{xxt} = 0$$

The a-b-c-d systems

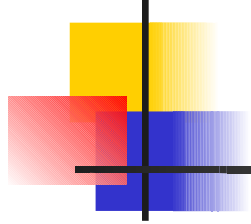
$$V_t + \nabla \eta + \epsilon \left(\frac{1}{2} \nabla |V|^2 + a \nabla \Delta \eta - b \Delta V_t \right) = 0$$

$$\eta_t + \nabla \cdot V + \epsilon (\nabla \cdot (\eta V) + c \Delta \nabla V - d \Delta \eta_t) = 0$$

where

$$a = \frac{1 - \theta^2}{2} \mu \qquad b = \frac{1 - \theta^2}{2} (1 - \mu)$$
$$c = \frac{\theta^2 - \frac{1}{3}}{2} \lambda \qquad d = \frac{\theta^2 - \frac{1}{3}}{2} (1 - \lambda)$$

and $\theta \in [0, 1]$, $\lambda, \mu \in \mathcal{R}$ present themselves as possible models for the surface waves.



If we set $\theta = \sqrt{\frac{2}{3}}$ and $\mu = \lambda = 0$, the Boussinesq equations become the coupled system

$$\begin{aligned} V_t + \nabla \eta + \epsilon \left(\frac{1}{2} \nabla |V|^2 - \frac{1}{6} \Delta V_t \right) &= 0, \\ \eta_t + \nabla \cdot V + \epsilon \left(\nabla \cdot (\eta V) - \frac{1}{6} \Delta \eta_t \right) &= 0. \end{aligned}$$



A simple rescaling of the independent and dependent variables gives the tidy form

$$\begin{aligned}V_t + \nabla \eta + \nabla |V|^2 - \Delta V_t &= 0, \\ \eta_t + \nabla \cdot V + \nabla \cdot (\eta V) - \Delta \eta_t &= 0.\end{aligned}$$

Happy Birthday, Walter!!
And best wishes for many,
many more!!

