# Arnold diffusion for convex nearly integrable systems

V. Kaloshin

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Arnold diffusion

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#### • Motivation: Ergodic and quasiergodic hypothesis.

- Nearly integrable systems and the problem of Arnold diffusion
- Results in 3, 4, and more degrees of freedom
- Indication of Arnold diffusion in the Solar system
- Stochastic aspects of Arnold diffusion

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Let  $H : \mathbb{R}^{2n} \to \mathbb{R}$  be a smooth function,  $(q, p) \in \mathbb{R}^n \times \mathbb{R}^n$ . Let  $X_H$  be the Hamiltonian flow associated to H.

$$\begin{cases} \dot{q} = \partial_{p}H \\ \dot{p} = -\partial_{q}H \end{cases}$$
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Let  $S_E = \{(q, p) \in T^*M : H(q, p) = E\}$  be an energy surface.

**Ergodic Hypothesis** (Boltzmann, Maxwell) Is a generic Hamiltonian flow  $X_H$  on a generic energy surface  $S_E$  ergodic?

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Numerical doubts (Fermi-Pasta-Ulam) Chains of nonlinear springs

$$\ddot{u}_n = k(u_{n+1} - u_n) - k(u_n - u_{n-1}) + \alpha(u_{n+1} - u_n)^2 + \alpha(u_n - u_{n-1})^2$$

the  $\alpha$ -term — nonlinearity. **Most** "small" solutions are **almost periodic**!

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## Quasiergodic Hypothesis

**KAM theory** Each *nearly integrable* systems has collections of invariant tori of positive measure  $\implies$  no ergodicity!



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#### Integrable systems & action-angles coordinates

#### Let $H : \mathbb{R}^{2n} \to \mathbb{R}$ be a Hamiltonian, $\varphi \in \mathbb{T}^n$ be angle, $I \in \mathbb{R}^n$ be action.

A Hamiltonian system is **Arnold-Liouville integrable** if for an open set  $U \subset \mathbb{R}^n$  there exists a symplectic map  $\Phi : \mathbb{T}^n \times U \to \mathbb{R}^{2n}$  s. t.  $H \circ \Phi(\varphi, I)$  depends only on I and

$$\begin{cases} \dot{\varphi} = \partial_l (H \circ \Phi)(I) = \omega(I), \\ \dot{I} = 0. \end{cases} \quad (\varphi, I) \text{-action-angle coordinates}$$

In particular,  $\Phi(\mathbb{T}^n \times U)$  is foliated by invariant *n*-dim'l tori & on each torus  $\mathbb{T}^n$  the flow is linear.

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- Pendulum  $H = \frac{l^2}{2} \cos 2\pi \varphi$ ,  $(\varphi, I) \in T^*\mathbb{T} = \mathbb{T} \times \mathbb{R}$ .
- Harmonic oscillator  $\ddot{q} = -kq$  or  $H = \frac{p^2}{2} + \frac{kq^2}{2}$ .
- Motion in a central force field  $\ddot{q} = F(||q||)q$ .
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- Lagrange's top, Kovaleskaya's top, Euler top.
- Toda lattice: chain ··· < x<sub>0</sub> < x<sub>1</sub> < ... with the neighbor interaction ∑<sub>i</sub> exp(x<sub>i</sub> − x<sub>i+1</sub>)
- Calogero-Moser system: chain of harmonic oscillators with a neighbor repulsive interaction.
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#### Arnold, 63: Let $(\varphi, I) \in T^* \mathbb{T}^n = \mathbb{T}^n \times \mathbb{R}^n, t \in \mathbb{T}$ .

 (weak form) Does there exist a real instability in many-dimensional problems of perturbation theory when the invariant tori do not divide the phase space? More precisely, for a generic perturbation εH<sub>1</sub>(φ, *I*, *t*) the Hamiltonian

$$H_{\varepsilon}(\varphi, I, t) = H_0(I) + \varepsilon H_1(\varphi, I, t)$$

has an orbit whose action component "travels" in action space, in particular,  $\max_t ||I(t) - I(0)|| = O(1)$ .

• (strong form) For any two open sets  $U, U' \subset B^n$  the Hamiltonian  $H_{\varepsilon}(\varphi, I, t)$  has an orbit whose action component "travels" from U to U', i.e.  $I(0) \in U$  and  $I(T) \in U'$  for some T > 0.

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## Nearly integrable systems in dimension 2

## Let $H_0(I) = \frac{I^2}{2}$ . Time one map $(\varphi, I) \to (\varphi + I, I) \pmod{1}$ .

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Arnold diffusion

#### KAM Theorem, obstacles to instability

Let  $H_0(I)$  have non-degenerate Hessian, e.g.  $H_0(I) = \sum I_i^2/2$ .



**KAM Theorem** Let  $H_{\varepsilon}(\varphi, I, t) = H_0(I) + \varepsilon H_1(\varphi, I, t)$  be a smooth perturbation. Then with probability  $1 - O(\sqrt{\varepsilon})$  has an initial condition in  $\mathbb{T}^n \times B^n \times \mathbb{T}$  having a quasiperiodic orbit. Moreover,  $\mathbb{T}^n \times B^n \times \mathbb{T}$  with **certain neighborhood of rational lines deleted** is laminated by invariant (n + 1)-dimensional tori, one for each diophantine  $\omega$ .

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## The heuristic picture



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In (2n + 1)-dimensional space there are (n + 1)-dimensional tori. For n = 1 they confine orbits! For n > 1 they do not!

The First Main Result For any  $\gamma > 0$  & a generic smooth perturbation  $\varepsilon H_1(\phi, I, t)$  the Hamiltonian

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Fig. 5. The number of asteroids as a function of the semi-major axis a. The  $\alpha$ -values corresponding to certain fractions of Jupiter's period are marked below. Some of these 'resonances' have produced gaps in the asteroid distribution.

J. Wisdom,83, Chaotic Behavior & the Origin of the 3/1 Kirkwood Gaps

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1.1.

3.0 LUL\_L 4.5

## Diffusion conjecture for Arnold's example

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 $f_0: (\varphi, I) \to (\varphi + I + \varepsilon \cos \varphi, I + \varepsilon \cos \varphi),$ 

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be a pair of standard maps.

Consider random composition of these maps

$$f_{\omega_n} \circ f_{\omega_{n-1}} \circ \cdots \circ f_{\omega_1}(\varphi_0, I_0) = (\varphi_n, I_n).$$

**Theorem** (joint work with O. Castejon) For  $n \sim \varepsilon^{-2}$  such compositions satisfy the Central Limit Theorem, i.e.

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## Preprints contributing to the talk

- P. Bernard., V. Kaloshin, K. Zhang, Arnold diffusion in arbitrary degrees of freedom and crumpled 3-dimensional normally hyperbolic invariant cylinders, arXiv:1112.2773v Dec 2011, 58pp.
- V. Kaloshin, K. Zhang, A strong form of Arnold diffusion for two and a half degrees of freedom, arXiv:1212.1150 [math.DS] Dec 2012, 208pp.;
- J. Fejoz, M. Guardia, V. Kaloshin, P. Roldan, Kirkwood gaps and diffusion along mean motion resonances in the restricted planar three-body problem, to appear in J of European Math Soc.
- V. Kaloshin, K. Zhang, Arnold diffusion for three and a half degrees of freedom, April 2014, 25pp.;
- V. Kaloshin, K. Zhang, Dynamics of the dominant Hamiltonian, with applications to Arnold diffusion, October 2014, 75pp.;
- M. Guardia, V. Kaloshin, Orbits of nearly integrable systems accumulating to KAM tori, preprint, 2014, 112pp. .

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Arnold diffusion