Quadrature-Based Moment Methods for Kinetic Models

Rodney O. Fox

Department of Chemical and Biological Engineering Iowa State University, Ames, Iowa, USA & École Centrale Paris, France

Workshop on Moment Methods in Kinetic Theory II Fields Institute, University of Toronto October 14-17, 2014

R. O. Fox (ISU & ECP) [Quadrature-Based Moment Methods](#page-44-0) WMMKTII 2014 1 / 45

Application: polydisperse multiphase flow

- continuous phase
- disperse phase
- size distribution
- finite particle inertia
- collisions
- variable mass loading
- multiphase turbulence

Bidisperse gas-particle flow (DNS of S. Subramaniam)

4 m + 4 m + 4 m

[Introduction](#page-2-0) [Target Applications](#page-2-0)

Application: polydisperse multiphase flows

Bubble columns

Power stations

Brown-out

Volcanos

Jet break up

Spray flames

 4 O \rightarrow 4 \overline{m} \rightarrow 4 \overline{m} \rightarrow

 \sim 舌

 QQ

Modeling challenges

- Strong coupling between continuous and disperse phases
- Wide range of particle volume fractions (even in same flow!)
- Inertial particles with wide range of Stokes numbers
- Collision-dominated to collision-less regimes in same flow
- Granular temperature can be very small and very large in same flow
- Particle polydispersity (e.g. size, density, shape) is always present

Need a modeling framework that can handle all aspects!

Overview of kinetic modeling approach

Mesoscale model incorporates more microscale physics in closures!

R. O. Fox (ISU & ECP) [Quadrature-Based Moment Methods](#page-0-0) WMMKTII 2014 5 / 45

 Ω

Types of mesoscale transport (kinetic) equations

• Population balance equation (PBE): $n(t, \mathbf{x}, \xi)$

$$
\frac{\partial n}{\partial t} + \frac{\partial}{\partial x_i} [u_i(t, \mathbf{x}, \xi)n] + \frac{\partial}{\partial \xi_i} [G_j(t, \mathbf{x}, \xi)n] = \frac{\partial}{\partial x_i} \left(D(t, \mathbf{x}, \xi) \frac{\partial n}{\partial x_i} \right) + \mathbb{S}
$$

with known velocity u, acceleration G, diffusivity *D* and source S • Kinetic equation (KE): $n(t, x, v)$

$$
\frac{\partial n}{\partial t} + \frac{\partial}{\partial x_i} (v_i n) + \frac{\partial}{\partial v_i} [A_i(t, \mathbf{x}, \mathbf{v}) n] = \mathbb{C}
$$

with known acceleration **A** and collision operator $\mathbb C$

Generalized population balance equation (GPBE): $n(t, x, v, \xi)$

$$
\frac{\partial n}{\partial t} + \frac{\partial}{\partial x_i} (v_i n) + \frac{\partial}{\partial v_i} [A_i(t, \mathbf{x}, \mathbf{v}, \xi) n] + \frac{\partial}{\partial \xi_i} [G_j(t, \mathbf{x}, \mathbf{v}, \xi) n] = \mathbb{C}
$$

with known accelerations **A**, **G** and collision/aggregation operator \mathbb{C}

 Ω

イロト イ押 トイヨ トイヨト

Moment transport equations

• PBE:
$$
M_k = \int \xi^k n d\xi
$$

\n
$$
\frac{\partial M_k}{\partial t} + \frac{\partial}{\partial x} \left(\int \xi^k u n d\xi \right) = k \int \xi^{k-1} G_n d\xi + \frac{\partial}{\partial x} \left(\int \xi^k D \frac{\partial n}{\partial x} d\xi \right) + \int \xi^k S d\xi
$$

• KE:
$$
M_k = \int v^k n \, dv
$$

$$
\frac{\partial M_k}{\partial t} + \frac{\partial M_{k+1}}{\partial x} = k \int v^{k-1} An \, dv + \int v^k C \, dv
$$

• GPBE:
$$
M_{kl} = \int v^k \xi^l n \, dv d\xi
$$

\n
$$
\frac{\partial M_{kl}}{\partial t} + \frac{\partial M_{k+1l}}{\partial x} = k \int v^{k-1} \xi^l An \, dv d\xi + l \int v^k \xi^{l-1} G n \, dv d\xi + \int v^k \xi^l C \, dv d\xi
$$

Terms in red will usually require mathematical closure

R. O. Fox (ISU & ECP) Duadrature-Based Moment Methods WMMKTII 2014 7/45

 QQ

 $\mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}^d$

4 0 8 \sim

Closure with moment methods

Close moment equations by reconstructing density function

R. O. Fox (ISU & ECP) [Quadrature-Based Moment Methods](#page-0-0) WMMKTII 2014 8/45

イロト イ押ト イヨト イヨト

 299

Quadrature-based moment methods (QBMM)

Basic idea: Given a set of transported moments, reconstruct the number density function (NDF)

Things to consider:

- Which moments should we choose?
- What method should we use for reconstruction?
- How can we extend method to multivariate phase space?
- How should we design the numerical solver for the moments?

We must be able to demonstrate *a priori* that numerical algorithm is robust and accurate!

Gauss quadrature in 1-D (real line)

• The formula

$$
\int g(v)n(v) dv = \sum_{\alpha=1}^{N} n_{\alpha}g(v_{\alpha}) + R_{N}(g)
$$

is a Gauss quadrature iff the *N* nodes v_α are roots of an N^{th} -order orthogonal polynomial $P_N(v)$ (\perp with respect to $n(v)$)

• Recursion formula for $P_N(v)$:

$$
P_{\alpha+1}(v) = (v - a_{\alpha})P_{\alpha}(v) - b_{\alpha}P_{\alpha-1}(v), \quad \alpha = 0, 1, 2, \dots
$$

Inversion algorithm (QMOM) for moments $M_k = \int v^k n(v) dv$:

$$
\{M_0, M_1, \ldots, M_{2N-1}\} \stackrel{\text{hard}}{\Longrightarrow} \{a_0, a_1, \ldots, a_{N-1}\}, \{b_1, b_2, \ldots, b_{N-1}\}
$$
\n
$$
\stackrel{\text{easy}}{\Longrightarrow} \{n_1, n_2, \ldots, n_N\}, \{v_1, v_2, \ldots, v_N\}
$$

 \sim

Szegö quadrature on unit circle

• If $n(\phi)$ is periodic on the unit circle:

$$
\int_{-\pi}^{\pi} g(e^{i\phi}) n(\phi) d\phi = \sum_{\alpha=1}^{N} n_{\alpha} g(e^{i\phi_{\alpha}}) + R_N(g)
$$

is a Szegö quadrature iff the *N* nodes $z_{\alpha} = e^{i\phi_{\alpha}}$ are zeros of an *N*th-order para-orthogonal polynomials $B_N(z)$

• Trigonometric moments:

$$
\langle \cos(n\phi) \rangle = \int_{-\pi}^{\pi} \frac{1}{2} (z^n + z^{-n}) n(\phi) d\phi, \langle \sin(n\phi) \rangle = \int_{-\pi}^{\pi} \frac{1}{2} (z^n - z^{-n}) n(\phi) d\phi
$$

are natural choice for reconstruction

• Except for special case $[n(\phi)$ symmetric wrt 0, no fast inversion algorithm is available to find n_{α} and ϕ_{α}

 Ω

イロト イ押 トイヨ トイヨト

1-D quadrature method of moments (QMOM)

Use Gaussian quadrature to approximate unclosed terms in moment equations:

$$
\frac{\mathrm{d}\mathbf{M}}{\mathrm{d}t} = \int \mathbf{S}(v)n(v)\mathrm{d}v \approx \sum_{\alpha=1}^{N} n_{\alpha} \mathbf{S}(v_{\alpha})
$$

where $\mathbf{M} = \{M_0, M_1, \ldots, M_{2N-1}\}\$ and **S** is "source term"

- Exact if S is polynomial of order $\leq 2N 1$
- Provides good approximation for most other cases with small $N \approx 4$
- Complications arise in particular cases (e.g. spatial fluxes)
- In all cases, moments M must remain **realizable** for moment inversion

N.B. equivalent to reconstructed *N*-point distribution function:

$$
n^*(v) = \sum_{\alpha=1}^N n_\alpha \delta(v - v_\alpha)
$$

 \implies **realizable** if $n_{\alpha} \geq 0$ for all α

 Ω

(ロトス例) スミトスミト

Quadrature in multiple dimensions

No method equivalent to Gaussian quadrature for multiple dimensions!

Given a realizable moment set $\mathbf{M} = \{M_{ijk} : i, j, k \in 0, 1, \dots\}$, find n_{α} and \mathbf{v}_{α} such that

$$
M_{ijk} = \int v_1^i v_2^j v_3^k n(\mathbf{v}) \mathrm{d}\mathbf{v} = \sum_{\alpha=1}^N n_\alpha v_{1\alpha}^i v_{2\alpha}^j v_{3\alpha}^k
$$

What moment set to use?

- If M corresponds to an *N*-point distribution, then method should be exact
- Avoid brute-force nonlinear iterative solver (poor convergence, ill-conditioned, too slow, . . .)
- Algorithm must be realizable (i.e. non-negative weights, ...)
- Strategy: choose an optimal moment set to av[oid](#page-11-0) [il](#page-13-0)[l-](#page-11-0)[co](#page-12-0)[n](#page-7-0)[d](#page-29-0)[i](#page-12-0)[ti](#page-17-0)[o](#page-18-0)n[e](#page-8-0)d [s](#page-30-0)[ys](#page-0-0)[tem](#page-44-0)s

Brute-force QMOM (2-D phase space)

Given $3n^2$ bivariate optimal moments ($n = 2$):

$$
M_{00} \t M_{01} \t M_{02} \t M_{03} M_{10} \t M_{11} \t M_{12} \t M_{13} M_{20} \t M_{21} M_{30} \t M_{31}
$$

• Solve 12 moment equations:

$$
\sum_{\alpha=1}^4 |n_{\alpha}| u_{\alpha}^i v_{\alpha}^j = M_{ij}
$$

to find $\{n_1, \ldots, n_4; u_1, \ldots, u_4; v_1, \ldots, v_4\}$

- Problem: iterative solver converges slowly (or not at all)
- Problem: system is singular for (nearly) degenerate cases

 Ω

イロト イ押ト イヨト イヨト

Conditional QMOM (2-D phase space)

• Conditional density function and conditional moments (2-D)

$$
n(u, v) = f(v|u)n(u) \quad \Longrightarrow \quad \langle V^k | U = u \rangle = \int v^k f(v|u) \, dv
$$

- 1-D adaptive quadrature for *U* direction $(n = 2)$ $\langle U^k \rangle = M_{k0}, k \in \{0, 1, 2, 3\} \Longrightarrow$ find weights ρ_i , abscissas u_i
- Solve linear systems for conditional moments $\langle V^k | u_i \rangle$:

$$
\begin{bmatrix} \rho_1 & \rho_2 \\ \rho_1 u_1 & \rho_2 u_2 \end{bmatrix} \begin{bmatrix} \langle V^k | u_1 \rangle \\ \langle V^k | u_2 \rangle \end{bmatrix} = \begin{bmatrix} \langle V^k \rangle \\ \langle UV^k \rangle \end{bmatrix} = \begin{bmatrix} M_{0k} \\ M_{1k} \end{bmatrix} \text{ for } k \in \{1, 2, 3\}
$$

• In principle, CQMOM controls 10 of 12 optimal moments:

R. O. Fox $(ISU & ECP)$

$$
M_{00} \quad M_{01} \quad M_{02} \quad M_{03}
$$
\n
$$
M_{10} \quad M_{11} \quad M_{12} \quad M_{13}
$$
\n
$$
M_{20}
$$
\n
$$
M_{30}
$$
\n
$$
\text{Quadrature-Based Moment Methods}
$$
\n
$$
WMMKTII 2014 \quad 15/45
$$

Conditional QMOM (cont.)

1-D adaptive quadrature in *V* direction for each *i*:

 $\langle V^k | u_i \rangle, k \in \{0, 1, 2, 3\} \Longrightarrow$ find weights ρ_{ij} , abscissas v_{ij}

- Adaptive quadrature sets some $\rho_{ii} = 0$ if subset of conditional moments are not realizable
- Reconstructed density: $n^*(u, v) = \sum_i \sum_j \rho_i \rho_{ij} \delta(u u_i) \delta(v v_{ij})$
- Conditioning on $V = v_i$ uses 10 of 12 optimal moments:

Union of two sets \implies optimal moment set

• Extension to higher-dimensional phase space [is](#page-14-0) s[tr](#page-16-0)[ai](#page-14-0)[gh](#page-15-0)[t](#page-16-0)[fo](#page-11-0)[r](#page-12-0)[w](#page-17-0)[ar](#page-7-0)[d](#page-8-0)

 2990

Optimal moment set

Moments needed for all COMOM permutations \implies Optimal moment set

27 moments

 $N = 9$ nodes in 2-D

Only optimal moment set is transported

R. O. Fox (ISU & ECP) Duadrature-Based Moment Methods WMMKTII 2014 17/45

Examples of 2-D quadrature

QBMM approximations for bivariate Gaussian with $\rho = 0$ (top) and $\rho = 0.5$ (bottom) for $N = 4$ (left) and $N = 9$ (right).

Brute-force QMOM (green diamond) Tensor-product QMOM (blue circle) CQMOM (red square)

÷.

4 0 8 4

 2990

CQMOM on unit circle

- For $n(\phi)$ on the unit circle, define $x = \cos \phi$ and $y = \sin \phi$
- Conditional pdf is known exactly $n(x, y) = n(x)f(y|x)$

$$
f(y|x) = w_1(x)\delta\left(y - \sqrt{1 - x^2}\right) + w_2(x)\delta\left(y + \sqrt{1 - x^2}\right)
$$

with one unknown $w_1(x)$ ($w_2 = 1 - w_1$)

- Apply QMOM with 2*N* moments $\langle \cos^n \phi \rangle$ to find n_α and $x_\alpha = \cos \phi_\alpha$
- COMOM requires conditional moment $\langle y|x_{\alpha}\rangle$ to find $w_1(x_{\alpha})$
- Apply CQMOM with *N* moments $\langle \sin \phi \cos^n \phi \rangle$ to find $w_{1\alpha} = w_1(x_\alpha)$
- Gauss/Swegö quadrature for symmetric ndf, otherwise fast, realizable reconstruction

KOD KARD KED KE DA GARA

Example of CQMOM on unit circle

R. O. Fox (ISU & ECP) [Quadrature-Based Moment Methods](#page-0-0) WMMKTII 2014 20 / 45

Extended quadrature method of moments (EQMOM)

Can we improve reconstructed distribution using kernel density functions?

$$
n(v) = \sum_{i=1}^{N} n_i \delta_{\sigma}(v, v_i)
$$

with *N* weights $n_i > 0$, *N* abscissas v_i but only one spread parameter $\sigma > 0$

Gaussian $(-\infty < v < +\infty)$:

$$
\delta_{\sigma}(v, v_i) \equiv \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left(-\frac{(v - v_i)^2}{2\sigma^2}\right)
$$

• Beta $(0 < v < 1)$: with $\lambda_i = v_i/\sigma$ and $\mu_i = (1 - v_i)/\sigma$

$$
\delta_{\sigma}(\nu,\nu_i) \equiv \frac{\nu^{\lambda_i - 1} (1 - \nu)^{\mu_i - 1}}{B(\lambda_i, \mu_i)}
$$

 Ω

 $A \cup B \cup A \cap B \cup A \subseteq B \cup A \subseteq B \cup B$

EQMOM algorithm

 $2N + 1$ moments of $n(v)$ (denote by m_k) for beta-EQMOM:

$$
m_0 = m_0^*
$$

\n
$$
m_1 = m_1^*
$$

\n
$$
m_2 = \frac{1}{1+\sigma} (\sigma m_1^* + m_2^*)
$$

\n
$$
m_3 = \frac{1}{(1+2\sigma)(1+\sigma)} (2\sigma^2 m_1^* + 3\sigma m_2^* + m_3^*)
$$

\n
$$
m_4 = \frac{1}{(1+3\sigma)(1+2\sigma)(1+\sigma)} (6\sigma^3 m_1^* + 11\sigma^2 m_2^* + 6\sigma m_3^* + m_4^*) \equiv m_{2N}^{\dagger}(\sigma)
$$

\nwith $m_k^* \equiv \sum_{i=1}^N n_i v_i^k$ (i.e. QMOM moments)
\nGiven m_k for $k = 0, ..., 2N$
\n**O** Guess σ
\n**O** Sales σ

2 Solve for m^* for $k = 0, ..., 2N - 1$

Solve for n_i and v_i using 1-D quadrature with m_k^* for $k = 0, \ldots, 2N - 1$

• Compute m_{2N}^* and resulting estimate m_2^{\dagger} 2*N*

5 Iterate on σ until $m_{2N} = m_2^{\dagger}$ 2*N*

Example: beta-EQMOM with $N = 2$

 $n_1 = n_2 = 1/2$, $\xi_1 = 1/3$, $\xi_2 = 2/3$ for different values of σ

First 2*N* moments always exact with max $\sigma : m_{2N} \ge m_2^{\dagger}$ $_{2N}^{\intercal}(\sigma)$ Converges to exact NDF as $N \to \infty$ (Gavriliadis and Athanassoulis 2002)

Closure with EQMOM

Unclosed integrals (given n_i , v_i and σ):

 $\ddot{}$

$$
\int g(v)n(v)dv = \sum_{i=1}^{N} n_i \int g(v)\delta_{\sigma}(v, v_i)dv
$$

Use Gaussian quadrature with known weights w_{ii} and abscissas v_{ii} :

$$
\int g(v)\delta_{\sigma}(v,v_i)dv = \sum_{j=1}^{M_i} w_{ij}g(v_{ij})
$$

where M_i can be chosen arbitrarily large to control error Dual-quadrature representation of EQMOM:

$$
n(v) = \sum_{i=1}^{N} \sum_{j=1}^{M_i} n_i w_{ij} \delta(v - v_{ij}) \quad (M_i = 1 \text{ when } \sigma = 0)
$$

=⇒ exact for polynomials of order ≤ 2*N*

 Ω

イロト イ押ト イヨト イヨト

EQMOM on unit circle

• For $n(\phi)$ on the unit circle, define $x = \cos \phi$ and $y = \sin \phi$, and

$$
n(x, y) = \sum_{\alpha=1}^{N} n_{\alpha} \delta_{\sigma_{\phi}}(x, x_{\alpha}) f_{\alpha}(y|x)
$$

where the kernel density $\delta_{\sigma_\phi}(x, x_\alpha)$ is periodic wrt ϕ

• Conditional pdf is known exactly, but with constant weights:

$$
f_{\alpha}(y|x) = w_{1\alpha}\delta\left(y - \sqrt{1 - x^2}\right) + w_{2\alpha}\delta\left(y + \sqrt{1 - x^2}\right)
$$

with $w_{2\alpha} = 1 - w_{1\alpha}$

- Apply EQMOM with $2N + 1$ moments $\langle \cos^n \phi \rangle$ to find n_α , x_α and σ_ϕ
- Apply CQMOM with *N* moments $\langle \sin \phi \cos^n \phi \rangle$ to find $w_{1\alpha}$

 Ω

 $A \cup B \cup A \cap B \cup A \subseteq B \cup A \subseteq B \cup B$

Multivariate EQMOM

Example: 2-D case \Rightarrow Extended CQMOM

$$
n(u, v) = n(u)f(v|u) = \sum_{\alpha=1}^{N} n_{\alpha} \delta_{\sigma_u}(u, u_{\alpha}) \left(\sum_{\beta=1}^{N_{\alpha}} n_{\alpha\beta} \delta_{\sigma_{v, \alpha}}(v, v_{\alpha\beta}) \right)
$$

with *N* abscissas u_{α} , $\mathcal{N} = \sum_{\alpha=1}^{N} N_{\alpha}$ weights $w_{\alpha\beta} = n_{\alpha} n_{\alpha\beta} \ge 0$ and \mathcal{N} abscissas $v_{\alpha\beta}$, but only one parameter $\sigma_u \geq 0$ and *N* parameters $\sigma_{v,\alpha}$ Define moments:

$$
M_{ij} = \int u^i v^j n(u, v) \, \mathrm{d}u \, \mathrm{d}v = \sum_{\alpha=1}^N \sum_{\beta=1}^{N_\alpha} w_{\alpha\beta} m_{1,i}^{(\alpha)} m_{2,j}^{(\alpha\beta)}
$$

where

$$
m_{1,i}^{(\alpha)} \equiv \int u^i \delta_{\sigma_u}(u, u_\alpha) \, \mathrm{d}u \qquad m_{2,j}^{(\alpha \beta)} \equiv \int v^j \delta_{\sigma_{v,\alpha}}(v, v_{\alpha \beta}) \, \mathrm{d}v
$$

are known functions of the EQMOM parameters

 Ω

 $A \cup B \rightarrow A \oplus B \rightarrow A \oplus B \rightarrow A \oplus B \rightarrow A \oplus B$

Algorithm for 2-D ECQMOM

1-D EQMOM for moments in *u*:

$$
M_{i0} = \sum_{\alpha=1}^{N} n_{\alpha} m_{1,i}^{(\alpha)} \quad \text{for } i = 0, \ldots, 2N \Longrightarrow n_{\alpha}, u_{\alpha} \text{ and } \sigma_{u}
$$

Use CQMOM to find conditional moments $\langle V^j \rangle_{\alpha} \equiv \sum_{\beta=1}^{N_{\alpha}} n_{\alpha\beta} m_{2,j}^{(\alpha\beta)}$ $\sum_{i=1}^{(\alpha,\beta)}$ from the bivariate moments (i.e. solve linear system):

$$
\sum_{\alpha=1}^N n_{\alpha} m_{1,i}^{(\alpha)} \langle V^j \rangle_{\alpha} = M_{ij} \quad \text{for } i = 0, \ldots, N-1
$$

• For each α , apply 1-D EQMOM to conditional moments:

$$
\{1, \langle V \rangle_{\alpha}, \ldots, \langle V^{2N_{\alpha}} \rangle_{\alpha}\} \implies n_{\alpha\beta}, v_{\alpha\beta} \text{ and } \sigma_{v,\alpha}
$$

Uses the extended optimal moment set

 Ω

イロト イ押 トイヨ トイヨト

Extended optimal moment set

All CQMOM permutations \Longrightarrow Extended optimal moment set

 $\mathcal{N} = 4$ nodes in 2-D

*M*₀₀ *M*₁₀ *M*₂₀ *M*₃₀ *M*₄₀ *M*⁰¹ *M*¹¹ *M*²¹ *M*³¹ *M*⁴¹ M_{02} M_{12} *M*⁰³ *M*¹³ *M*⁰⁴ *M*¹⁴

*M*⁰⁰ *M*¹⁰ *M*²⁰ *M*³⁰ *M*⁴⁰ *M*⁵⁰ *M*⁶⁰ *M*⁰¹ *M*¹¹ *M*²¹ *M*³¹ *M*⁴¹ *M*⁵¹ *M*⁶¹ *M*⁰² *M*¹² *M*²² *M*³² *M*⁴² *M*⁵² *M*⁶² *M*⁰³ *M*¹³ *M*²³ M_{04} M_{14} M_{24} *M*⁰⁵ *M*¹⁵ *M*²⁵ *M*⁰⁶ *M*¹⁶ *M*²⁶

 $\mathcal{N} = 9$ nodes in 2-D

16 moments

33 moments

Only extended optimal moment set is transported

KOD KARD KED KE DAGA

ECQMOM on unit sphere

• For $n(\phi, \theta)$ on the unit sphere, ECOMOM reconstruction is

$$
n(\phi,\theta) = \sum_{\alpha=1}^{N} n_{\alpha} \delta_{\sigma_{\theta}}(\theta,\theta_{\alpha}) \left(\sum_{\beta=1}^{N_{\alpha}} n_{\alpha\beta} \delta_{\sigma_{\phi,\alpha}}(\phi,\phi_{\alpha\beta}) \right)
$$

with periodic kernel density functions for $\theta \in [0, \pi]$ and $\phi \in [-\pi, \pi]$

- Define $z = \cos \theta$ and conditional pdf $f(\phi|z)$
- Apply EQMOM for $2N + 1$ moments $\langle z^n \rangle$ to find n_α , θ_α and σ_θ
- Use CQMOM to find trigonometric moments involving ϕ conditioned on $z_\alpha = \cos \theta_\alpha$
- For each α , apply EQMOM on unit circle to conditional moments to find $n_{\alpha\beta}$, $\phi_{\alpha\beta}$, $\sigma_{\phi\alpha}$ and $w_{1\alpha\beta}$

KOD KARD KED KE DA GARA

Summary of QBMM

- Extended optimal moment set \implies reconstruct NDF with fast, robust algorithm
- NDF must be realizable and moment-inversion algorithm must be robust
	- CQMOM is always realizable by construction
	- EQMOM gives a smooth NDF with low computational cost
- Current "best" moment-inversion algorithms:
	- 1-D phase space \Longrightarrow EQMOM
	- Multivariate phase space \implies multivariate ECQMOM
- Dual-quadrature representation used to close source terms
- Given smooth NDF, high-order kinetic-based transport solvers can be derived to ensure realizability

 QQ

Kinetic-based finite-volume methods (KBFVM)

Given a set of extended optimal moments, solve

$$
\frac{\partial M_{kl}}{\partial t} + \frac{\partial M_{k+1l}}{\partial x} = k \int v^{k-1} \xi^l An \, dv d\xi + l \int v^k \xi^{l-1} G n \, dv d\xi + \int v^k \xi^l \mathbb{C} \, dv d\xi
$$

where RHS is closed using QBMM:

$$
\frac{\partial M_{kl}}{\partial t} + \frac{\partial M_{k+1l}}{\partial x} = \sum_{\alpha=1}^{N} n_{\alpha} \left\{ k v_{\alpha}^{k-1} \xi_{\alpha}^{l} A_{\alpha} + l v_{\alpha}^{k} \xi_{\alpha}^{l-1} G_{\alpha} + v_{\alpha}^{k} \xi_{\alpha}^{l} \mathbb{C}_{\alpha} \right\}
$$

Things to consider:

- How do we discretize the spatial fluxes?
- How do we update the moments in time?
- How can we ensure that the moments are always **realizable**?

Kinetic-based spatial fluxes

Spatial fluxes can use kinetic formulation: e.g. $\partial_t M_{00} + \partial_x M_{10} = 0$

$$
M_{10} = Q_{10}^- + Q_{10}^+
$$

=
$$
\int_{-\infty}^0 u \left(\int n^*(u, v) dv \right) du + \int_0^\infty u \left(\int n^*(u, v) dv \right) du
$$

Using reconstructed *n*[∗], downwind and upwind flux components are

$$
Q_{10}^- = \sum_{\alpha=1}^N n_{\alpha} u_{\alpha} I_{(-\infty,0)} (u_{\alpha}) \qquad Q_{10}^+ = \sum_{\alpha=1}^N n_{\alpha} u_{\alpha} I_{(0,\infty)} (u_{\alpha})
$$

where $I_{\mathcal{S}}(x)$ is the indicator function for the interval \mathcal{S}

Kinetic-based fluxes are always hyperbolic

 Ω

4 ロ ト 4 何 ト 4 ヨ ト 4 ヨ ト

Finite-volume method: definitions

• 1-D advection problem:

$$
\frac{\partial \mathbf{M}}{\partial t} + \frac{\partial \mathbf{F}(\mathbf{M})}{\partial x} = 0
$$

where $\mathbf{M} = \int \mathbf{K}(v) n(v) dv$ and $\mathbf{F}(\mathbf{M}) = \int v \mathbf{K}(v) n(v) dv$

Finite-volume representation of moment vector:

$$
\mathbf{M}_{i}^{n} \equiv \frac{1}{\Delta x} \int_{x_{i}}^{x_{i+1}} \mathbf{M}(t_{n}, x) \mathrm{d}x
$$

• Finite-volume formula:

$$
\mathbf{M}^{n+1}_i = \mathbf{M}^n_i - \lambda \left[\mathbf{G}\left(\mathbf{M}^n_{i+\frac{1}{2},l},\mathbf{M}^n_{i+\frac{1}{2},r}\right) - \mathbf{G}\left(\mathbf{M}^n_{i-\frac{1}{2},l},\mathbf{M}^n_{i-\frac{1}{2},r}\right)\right]
$$

where $\mathbf{G}(\mathbf{M}_l, \mathbf{M}_r) = \int v^+ \mathbf{K}(v) n_l(v) dv + \int v^- \mathbf{K}(v) n_r(v) dv$

Realizability and spatial fluxes

- Flux functions: given M_i^n define $G(M_l, M_r)$ to achieve high-order spatial accuracy but keep M_i^{n+1} realizable!
- Discrete distribution function: Define

$$
\mathbf{M}_i^{n+1} \equiv \int \mathbf{K}(v) h_i(v) \mathrm{d}v
$$

and finite-volume formula can be written as

$$
h_i(v) = \lambda |v^-| n_{i+\frac{1}{2},r}^n + \lambda v^+ n_{i-\frac{1}{2},l}^n + n_i^n - \lambda |v^-| n_{i-\frac{1}{2},r}^n - \lambda v^+ n_{i+\frac{1}{2},l}^n
$$

(black part ≥ 0 , red part can be negative)

• Sufficient condition for realizable moments: $h_i(v) \geq 0$ for all *v* and *i*

 Ω

(ロ) (何) (ヨ) (ヨ)

Realizable, high-order, spatial fluxes

• First order:
$$
n_{i-\frac{1}{2},r}^n = n_{i+\frac{1}{2},l}^n = n_i^n
$$
 so that
\n
$$
h = \lambda |v^-| n_{i+1}^n + \lambda v^+ n_{i-1}^n + (1 - \lambda |v^-| - \lambda v^+) n_i^n \implies \frac{1}{|v^-| + v^+} \ge \lambda
$$
\nMoments are realizable, but scheme is diffusive ...

Quasi-higher order: Let $n_i^n = \sum_{\alpha} \rho_{\alpha,i}^n \delta(v - v_{\alpha,i}^n)$ and define

$$
n_{i-\frac{1}{2},r}^{n} = \sum_{\alpha} \rho_{\alpha,i-\frac{1}{2},r}^{n} \delta(\nu - \nu_{\alpha,i}^{n})
$$

$$
n_{i+\frac{1}{2},l}^{n} = \sum_{\alpha} \rho_{\alpha,i+\frac{1}{2},l}^{n} \delta(\nu - \nu_{\alpha,i}^{n})
$$

so that

$$
h = \lambda |v^-| n_{i + \frac{1}{2},r}^n + \lambda v^+ n_{i - \frac{1}{2},l}^n + \sum_{\alpha} \left(\rho_{\alpha,i}^n - \lambda |v^-| \rho_{\alpha,i - \frac{1}{2},r}^n - \lambda v^+ \rho_{\alpha,i + \frac{1}{2},l}^n \right) \delta(v - v_{\alpha,i}^n)
$$

\n
$$
\implies \min_{\alpha} \left(\frac{\rho_{\alpha,i}^n}{|v_{\alpha,i}| \rho_{\alpha,i - \frac{1}{2},r}^n + v_{\alpha,i}^+ \rho_{\alpha,i + \frac{1}{2},l}^n} \right) \ge \lambda
$$

Use high-order, finite-volume schemes only for the weights

 QQ

 $A \cup B \rightarrow A \oplus B \rightarrow A \oplus B \rightarrow A \oplus B \rightarrow A \oplus B$

Realizable time-stepping schemes

• First-order explicit:

$$
\mathbf{M}^{n+1}_i = \mathbf{M}^n_i - \lambda \left[\mathbf{G}\left(\mathbf{M}^n_{i+\frac{1}{2},l},\mathbf{M}^n_{i+\frac{1}{2},r}\right) - \mathbf{G}\left(\mathbf{M}^n_{i-\frac{1}{2},l},\mathbf{M}^n_{i-\frac{1}{2},r}\right)\right]
$$

is realizable

- Second-order Runga-Kutta (RK2) is not realizable
- **RK2SSP:**

$$
\begin{aligned} & \mathbf{M}_{i}^{*}=\mathbf{M}_{i}^{n}-\lambda\left[\mathbf{G}\left(\mathbf{M}_{i+\frac{1}{2},l}^{n},\mathbf{M}_{i+\frac{1}{2},r}^{n}\right)-\mathbf{G}\left(\mathbf{M}_{i-\frac{1}{2},l}^{n},\mathbf{M}_{i-\frac{1}{2},r}^{n}\right)\right] \\ & \mathbf{M}_{i}^{**}=\mathbf{M}_{i}^{*}-\lambda\left[\mathbf{G}\left(\mathbf{M}_{i+\frac{1}{2},l}^{*},\mathbf{M}_{i+\frac{1}{2},r}^{*}\right)-\mathbf{G}\left(\mathbf{M}_{i-\frac{1}{2},l}^{*},\mathbf{M}_{i-\frac{1}{2},r}^{*}\right)\right] \\ & \mathbf{M}_{i}^{n+1}=\frac{1}{2}\left(\mathbf{M}_{i}^{n}+\mathbf{M}_{i}^{**}\right) \end{aligned}
$$

is realizable

Achieve second order in space and time on unstructured grids

4 0 8 4

 Ω

 \triangleright 4 \equiv \triangleright 4 \equiv \triangleright

Bubbly flow

Loading movie. . .

Quasi-secon[d](#page-35-0)-order *[re](#page-37-0)alizable* finite-volume scheme [on u](#page-35-0)[nst](#page-37-0)[ru](#page-35-0)[ctu](#page-36-0)red [me](#page-36-0)[s](#page-37-0)[h](#page-29-0) $\epsilon \geq 1$ Ğ. 299

R. O. Fox (ISU & ECP) Ouadrature-Based Moment Methods WMMKTII 2014 37/45

Summary of KBFVM

- When solving moment transport equations, we must guarantee realizability
- First-order FV methods are realizable, but too diffusive
- Standard high-order FV methods lead to unrealizable moments
- Kinetic-based flux functions can be designed to be realizable
- Use dual-quadrature representation with high-order spatial reconstruction
- High-order time-stepping schemes are also possible
- KBFVM provide robust treatment of shocks/discontinuous solutions

 Ω

(ロトス例) スミトスミト

Final remarks

- Mesoscopic models have direct link with underlying physics and result in a kinetic equation
- QBMM solves kinetic equation by reconstructing distribution function from moments
- Reconstruction requires realizable moments
- Numerical schemes must ensure that moments are always realizable
- QBMM on unit sphere can be used for radiation transport

 Ω

ADAKED KED

Principal collaborators and funding

- École Centrale Paris: C. Chalons, F. Laurent, M. Massot, A. Vié
- Iowa State University: V. Vikas, Z. J. Wang, C. Yuan
- Oakridge National Lab: C. Hauck
- US Department of Energy (Grant DE-FC26-07NT43098)
- US National Science Foundation (Grant CFF-0830214)
- Marie-Curie Senior Fellowship

Computational Models for Polydisperse Particulate and Multiphase Systems

R. O. Fox (ISU & ECP) [Quadrature-Based Moment Methods](#page-0-0) WMMKTII 2014 40 / 45

Thanks for your attention!

Questions?

4. 0. 8

Related references I

CHALONS, C., FOX, R. O. & MASSOT, M. 2010 A multi-Gaussian quadrature method of moments for gas-particle flows in a LES framework. *Proceedings of the Summer Program 2010*. Center for Turbulence Research, Stanford, pgs. 347–358.

CHALONS, C., FOX, R. O., LAURENT, F., MASSOT, M. & VIÉ, M. 2011 A multi-Gaussian quadrature method of moments for simulating high-Stokes-number turbulent two-phase flows. *Annual Research Briefs 2011*. Center for Turbulence Research, Stanford, pgs. 309–320.

E

Ħ

F

CHENG, J. C. & FOX, R. O. 2010 Kinetic modeling of nanoprecipitation using CFD coupled with a population balance. *Industrial & Engineering Chemistry Research* 49, 10651–10662.

CHENG, J. C., VIGIL, R. D. & FOX, R. O. 2010 A competitive aggregation model for Flash NanoPrecipitation. *Journal of Colloid and Interface Science* 351, 330–342.

FOX, R. O. 2008 A quadrature-based third-order moment method for dilute gas-particle flows. *Journal of Computational Physics* 227, 6313–6350.

FOX, R. O. 2009 Higher-order quadrature-based moment methods for kinetic equations. *Journal of Computational Physics* 228, 7771–7791.

 Ω

イロト イ押 トイヨ トイヨト

Related references II

FOX, R. O. 2009 Optimal moment sets for the multivariate direct quadrature method of moments. *Industrial & Engineering Chemistry Research* 48, 9686–9696.

F

Ħ

R

R

FOX, R. O. 2012 Large-eddy-simulation tools for multiphase flows. *Annual Review of Fluid Mechanics* 44, 47–76.

FOX, R. O. & VEDULA, P. 2010 Quadrature-based moment model for moderately dense polydisperse gas-particle flows. *Industrial & Engineering Chemistry Research* 49, 5174–5187.

ICARDI, M., ASINARI, P., MARCHISIO, D. L., IZQUIERDO, S. & FOX, R. O. 2012 Quadrature-based moment closures for non-equilibrium flows: hard-sphere collisions and approach to equilibrium. *Journal of Computational Physics*, (in press).

MARCHISIO, D. L. & FOX, R. O. 2005 Solution of population balance equations using the direct quadrature method of moments. *Journal of Aerosol Science* 36, 43–73.

MEHTA, M., RAMAN, V. & FOX, R. O. 2012 On the role of gas-phase chemistry in the production of titania nanoparticles in turbulent flames. *Chemical Engineering Science*, (submitted).

MEHTA, M., SUNG, Y., RAMAN, V. & FOX, R. O. 2010 Multiscale modeling of TiO₂ nanoparticle production in flame reactors: Effect of chemical mechanism. *Industrial & Engineering Chemistry Research* 49, 10663–10673.

PASSALACQUA, A. & FOX, R. O. 2011 Advanced continuum modeling of gas-particle flows beyond the hydrodynamic limit. *Applied Mathematical Modelling* 35, 1616–1627.

 Ω

 $\left\{ \begin{array}{ccc} 1 & 0 & 0 \\ 0 & 1 & 0 \end{array} \right.$

Related references III

PASSALACQUA, A. & FOX, R. O. 2011 An iterative solution procedure for multi-fluid gas-particle flow models on unstructured grids. *Powder Technology* 213, 174–187.

PASSALACQUA, A. & FOX, R. O. 2012 Simulation of mono- and bidisperse gas-particle flow in a riser with a third-order quadrature-based moment method. *Industrial & Engineering Chemistry Research* 52, 187-198.

PASSALACQUA, A., FOX, R. O., GARG, R. & SUBRAMANIAM, S. 2010 A fully coupled quadrature-based moment method for dilute to moderately dilute fluid-particle flows. *Chemical Engineering Science* 65, 2267–2283.

PASSALACQUA, A., GALVIN, J. E., VEDULA, P., HRENYA, C. M. & FOX, R. O. 2011 A quadrature-based kinetic model for dilute non-isothermal granular flows. *Communications in Computational Physics* 10, 216–252.

SUNG, Y., RAMAN, V. & FOX, R. O. 2011 Large-eddy simulation based multiscale modeling of TiO₂ nanoparticle synthesis in turbulent flame reactors using detailed nucleation chemistry. *Chemical Engineering Science* 66, 4370–4381.

VIKAS, V., WANG, Z. J., PASSALACQUA, A. & FOX, R. O. 2011 Realizable high-order finite-volume schemes for quadrature-based moment methods. *Journal of Computational Physics* 230, 5328–5352.

VIKAS, V., HAUCK, C. D., WANG, Z J. & FOX, R. O. 2013 Radiation transport modeling using extended quadrature methods of moments. *Journal of Computational Physics* 246, 221–241.

VIKAS, V., WANG, Z. J. & FOX, R. O. 2013 Realizable high-order finite-volume schemes for quadrature-based moment methods applied to diffusion population balance equations. *Journal of Computational Physics* 249, 162–179.

VIKAS, V., YUAN, C. WANG, Z. J. & FOX, R. O. 2011 Modeling of bubble-column flows with quadrature-based moment methods. *Chemical Engineering Science* 66, 3058–3070.

 Ω

イロト イ押 トイヨ トイヨト

Related references IV

YUAN, C. & FOX, R. O. 2011 Conditional quadrature method of moments for kinetic equations. *Journal of Computational Physics* 230, 8216–8246.

YUAN, C., LAURENT, F. & FOX, R. O. 2012 An extended quadrature method of moments for population balance equations. *Journal of Aerosol Science* 51, 1–23.

イロト イ母 トイヨ トイヨト

 2990