Relaxed Lagrangian solutions for the Semi-Geostrophic Shallow Water system in physical space

A. Tudorascu, West Virginia U.

Based on joint work with M. Feldman (UW-Madison)

Outline

The SGSW system

- SGSW in physical space
- SGSW in a readable form
- Existence unknown
- SGSW in dual space

2 Lagrangian solutions

- Yet another system
- Weak Lagrangian solutions in physical space
- Case of singular measures in dual space
- Renormalized solutions
- Relaxed renormalized Lagrangian solution

3 Main results

- Weak stability
- Existence
- Comments

Setting

• SGSW models the motion of a fluid rapidly rotating around the vertical axis x_3 , contained within the evolving 3-dimensional region $\mathcal{D}(t)$ which has the structure:

$$\mathcal{D}(t) = \{ (x_1, x_2, x_3) \in \mathbb{R}^3 : (x_1, x_2) \in \Omega, \ 0 \le x_3 \le h(t, x_1, x_2) \},\$$

where the region Ω of the (x_1, x_2) -plane is given and fixed, but the height h above the reference level is unknown and can evolve in time. The pressure on the top boundary of the fluid is a given constant p_0 , and

$$p(t, x_1, x_2, x_3) = [h(t, x_1, x_2) - x_3] + p_0.$$

• The horizontal components of the velocity $\mathbf{v} = (\mathbf{u}, u_3)$ of the fluid are independent of x_3 .

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SGSW in terms of the modified pressure

• Introduce

$$P(t,x) = p(t,x) + (x_1^2 + x_2^2)/2,$$

to write SGSW as

$$\begin{split} D_t X &= J(X - x), \\ \partial_t h + \nabla \cdot (h\mathbf{u}) &= 0 \\ X &= \nabla P, \ P &= h + \frac{1}{2} |\mathrm{Id}_{\Omega}|^2 \quad \mathrm{in} \ [0, T) \times \Omega; \\ \mathbf{u} \cdot \nu &= 0 & & \mathrm{on} \ [0, T) \times \partial \Omega, \\ P(0, \cdot) &= P_0 & & \mathrm{in} \ \Omega, \end{split}$$

• Here $J = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$.

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- Let (P, \mathbf{u}) be a solution and set $\mathcal{X}(t) := \{(Y, \rho) : Y : \Omega \to \mathbb{R}^2 \text{ Borel}, \ \rho \in \mathcal{P}^{ac}(\Omega), \ Y_{\#}\rho = \nabla P(t, \cdot)_{\#}h(t, \cdot)\}.$
- (Cullen & Shutts) Then $(\nabla P(t,\cdot),h(t,\cdot))$ is a critical point for

$$I(Y,\rho) = \int_{\Omega} |Y(x) - x|^2 \rho(x) dx + \int_{\Omega} \rho^2(x) dx$$

over $\mathcal{X}(t)$, for all $t \in (0,T)$.

- (Cullen & Purser) Only minimizers are stable, in the sense that SG accurately describes their evolution.
- In the language of Optimal Transport, this means $X(t, \cdot)$ must be the gradient of a convex function, i.e. $P(t, \cdot)$ must be convex for all $t \in (0, T)$.
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• No existence results for weak Eulerian solutions.

Theorem (Feldman & T.)

Let (P, \mathbf{u}) be a distributional solution for SG in the physical space such that $\nabla P \in H^1(0, T; L^2(\Omega; \mathbb{R}^3))$. Then $\alpha_t := \nabla P_{t\#}\chi$ is atom-free for \mathcal{L}^1 -a.e. $t \in (0, T)$.

 (Cullen) The model must accommodate solutions for which ∇P_t is locally constant. Observations show that the atmosphere contains significant regions where the potential temperature and absolute momentum of the atmosphere are well-mixed, representing a state of neutrality to parcel displacements (yielding "flat spots" in ∇P_t). Such states commonly arise as a result of atmospheric forcing either through surface heating or latent heat release.

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A change of variable and a new system

• Let $\nabla P_{t\#}h_t =: \alpha_t$. If $\int_{\Omega} h_0(x)dx = 1$, then α_t is a Borel probability. • Let $X = \nabla P(t, x)$. Then α_t solves:

 $\begin{aligned} \partial_t \alpha + \nabla \cdot (U\alpha) &= 0 & \text{in } [0,T) \times \mathbb{R}^2; \\ \nabla P(t,\cdot)_{\#} h(t,\cdot) &= \alpha(t,\cdot) & \text{for any } t \in [0,T); \\ P(t,x) &= h(t,x) + |x|^2/2, & \text{with } P(t,\cdot) \text{ convex for all } t \in [0,T); \\ U(t,X) &= J[X - \bar{\gamma}(t,X)], \\ \alpha(0,X) &= \alpha_0(X) & \text{for a e } X \in \mathbb{R}^2 \end{aligned}$

where $\bar{\gamma}(t, X)$ denotes the barycentric projection onto α_t of the optimal transfer plan between α_t and h_t , i.e.

$$\int_{\mathbb{R}^2} \xi(X) \cdot \bar{\gamma}(X) \, \alpha(dX) = \iint_{\mathbb{R}^2 \times \Omega} \xi(X) \cdot y \, \gamma(dX, dy)$$

for all continuous $\xi : \mathbb{R}^2 \to \mathbb{R}^2$ of at most quadratic growth, where γ is the (unique) optimal plan between α and h.

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 $\begin{aligned} \partial_t \alpha + \nabla \cdot (U\alpha) &= 0 & \text{in } [0,T) \times \mathbb{R}^2; \\ \nabla P(t,\cdot)_{\#} h(t,\cdot) &= \alpha(t,\cdot) & \text{for any } t \in [0,T); \\ P(t,x) &= h(t,x) + |x|^2/2, & \text{with } P(t,\cdot) \text{ convex for all } t \in [0,T); \\ U(t,X) &= J[X - \bar{\gamma}(t,X)], \\ \alpha(0,X) &= \alpha_0(X) & \text{for a e } X \in \mathbb{R}^2 \end{aligned}$

where $\bar{\gamma}(t, X)$ denotes the barycentric projection onto α_t of the optimal transfer plan between α_t and h_t , i.e.

$$\int_{\mathbb{R}^2} \xi(X) \cdot \bar{\gamma}(X) \, \alpha(dX) = \iint_{\mathbb{R}^2 \times \Omega} \xi(X) \cdot y \, \gamma(dX, dy)$$

for all continuous $\xi : \mathbb{R}^2 \to \mathbb{R}^2$ of at most quadratic growth, where γ is the (unique) optimal plan between α and h.

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Fields Institute (Toronto)

SG in dual variables

- Existence of solutions in dual variables: Cullen & Gangbo in the case $\alpha_t \ll \mathcal{L}^2$.
- This is a Hamiltonian system (Gangbo & Pacini, Ambrosio & Gangbo) for

$$H(\mu) = -\frac{1}{2} \inf_{h \in \mathcal{P}^{ac}(\Omega)} \left\{ W_2^2(\mu, h) + \|h\|_{L^2(\Omega)}^2 \right\}.$$

• Uniqueness is open!

Assume one has a classical solution to SGSW in physical space.
If u has a flow F : [0,T] × Ω → Ω defined by

 $\dot{F}(t,x) = \mathbf{u}(t,F(t,x)), \ \ F(0,x) = x \text{ for } h_0-\text{a.e. } x \in \Omega,$

then F can replace u as an unknown. Let $Z(t,x) := \nabla P(t,F(t,x))$. • The system for (P,F) becomes

 $\partial_t Z(t,x) = J \begin{bmatrix} Z(t,x) - F(t,x) \end{bmatrix} \text{ for } t \in [0,T) \text{ and } h_0 - \text{a.e. } x \in \Omega,$ $Z(0,x) = \nabla P_0(x) \qquad \qquad h_0 - \text{a.e. in } \Omega.$ (LE)

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Let $P_0 \in L^{\infty}(\Omega)$ be convex such that $h_0 := P_0 - |\mathrm{Id}_{\Omega}|^2/2 \in \mathcal{P}(\Omega)$, and let $p \in [1, \infty)$. Let $P : [0, T) \times \Omega \to \mathbb{R}$ such that

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Let $F \in C([0,T); L^p(h_0; \mathbb{R}^2))$ be a Borel map. The pair (P,F) is called a weak Lagrangian solution of SGSW in physical space if

• $F(0, \cdot) = \operatorname{Id}_{\Omega} h_0$ - a.e. in Ω , $P(0, x) = P_0(x)$ for a.e. $x \in \Omega$,

• for any t > 0 we have $F_{t\#}h_0 = h_t$;

• There exists a Borel map $F^*: [0,T) \times \Omega \to \Omega$ such that for every $t \in (0,T)$ we have $F^*_{t\#}h_t = h_0$, and $F^*_t \circ F_t(x) = x$ for h_0 -a.e. $x \in \Omega$ and $F_t \circ F^*_t(x) = x$ for h_t -a.e. $x \in \Omega$;

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- Cullen & Feldman used Ambrosio's theory of regular Lagrangian flows to construct weak Lagrangian solutions in physical space in the case $\alpha_0 =: \nabla P_{0\#} h_0 \in L^p(\nabla P_0(\Omega))$ for some p > 1, where P_0 is bounded and convex in an open ball containing $\overline{\Omega}$.
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- If $\alpha_0 = 1.25 \pi \delta_{z_0}$ for some $z_0 \in \partial B(0, 1)$, one can readily check that if $\dot{z}(t) = 0.8 Jz(t), \ z(0) = z_0$, then $\alpha_t := 1.25 \pi \delta_{z(t)}$ solves SGSW in dual space.
- Here, $h_t(x) = 2 |x z(t)|^2/2$, $P_t(x) = x \cdot z(t) + 3/2$.
- There are no maps F_t as in the definition of weak Lagrangian solutions.
- It is natural to ask if one can appropriately weaken the notion so as to accommodate this case.
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- It is natural to ask if one can appropriately weaken the notion so as to accommodate this case.
- Desirable to accommodate the general case $\alpha_0 \in \mathcal{P}_2(\mathbb{R}^2)$.

(1)

• For $\xi \in C^1(\mathbb{R}^2) \cap \operatorname{Lip}(\mathbb{R}^2)$ the map $t \mapsto \xi(Z(t,x))$ is absolutely continuous for \mathcal{L}^2 -a.e. $x \in \Omega$ and

$$\frac{\mathrm{d}}{\mathrm{d}t}\xi(Z(t,x))h_0(x) = \nabla\xi(Z(t,x)) \cdot J[Z(t,x) - F(t,x)]h_0(x)$$

for a.e. $t \in [0,T]$. Consequently, a more general, "renormalized" version of (LE) is available in the form

$$\begin{split} &\int_0^T \!\!\!\int_{\Omega} \Big\{ \xi(Z(t,x)) \partial_t \zeta(t,x) + \nabla \xi(Z(t,x)) \cdot J \big[F(t,x) \\ &- Z(t,x) \big] \zeta(t,x) \Big\} h_0(x) \, dx \, dt + \int_{\Omega} \xi(\nabla P_0(x)) \zeta(0,x) h_0(x) \, dx = 0, \\ &\text{for any } \xi \in C^1(\mathbb{R}^2) \cap \operatorname{Lip}(\mathbb{R}^2), \ \zeta \in C^1_c([0,T) \times \Omega). \end{split}$$

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for all $\xi \in C_b((0,T) \times \Omega \times \Omega)$. We notice first that the property $F(t,\cdot)_{\#}h_0 = h_t$ shows that σ disintegrates as

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Relaxed notion; definition

• Let $P_0 : \mathbb{R}^2 \to \mathbb{R} \cup \{\infty\}$ be convex such that $P_0|_{\Omega} \in L^2(\Omega)$, $P_0|_{\Omega} - \frac{1}{2}|\mathrm{Id}_{\Omega}|^2 =: h_0 \in \mathcal{P}(\Omega)$ and $\nabla P_0|_{\Omega} \in L^2(h_0; \mathbb{R}^2)$

Definition of Relaxed (Renormalized) Lagrangian Solutions

Consider a Borel function $P: [0,T) \times \mathbb{R}^2 \to \mathbb{R} \cup \{\infty\}$ such that $P(t, \cdot)$ is convex for all $t \in [0,T)$ and a Borel family of probability measures $[0,T) \ni t \mapsto \sigma_t \in \mathcal{P}(\Omega \times \Omega)$. Let σ be given by $d\sigma = d\sigma_t dt$. We say that (P,σ) is a *relaxed Lagrangian solution* for the SGSW system with initial data P_0 if

- $P(0,\cdot)|_{\Omega} \equiv P_0|_{\Omega};$
- (1) $P_t |\mathrm{Id}_{\Omega}|^2/2 =: h_t \in \mathcal{P}(\Omega)$ for all $t \in [0,T);$

• $\nabla P \in L^2(h; \mathbb{R}^2)$ (as functions of both variables);

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Weak stability

Theorem

Let P_0 , P_0^n satisfy (C) for all positive integers n with respect to Ω . Assume (P^n, σ^n) are relaxed solutions for SGSW in physical space corresponding to the initial data P_0^n . Then, possibly up to a subsequence, (P^n, σ^n) converges to a relaxed solution (P, σ) corresponding to the initial datum P_0 . The convergence is in the following sense:

(i) $P_t^n
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(ii) $\nabla P_t^n \to \nabla P_t$ a.e. in Ω for all $t \in [0, T]$, and $\{\nabla P_t\}_n$ is locally bounded uniformly with respect to $t \in [0, T]$ and $n \ge 1$;

(iii) σ^n converges narrowly to σ .

Furthermore, the corresponding dual space solutions satisfy $W_p(\alpha_t^n, \alpha_t) \to 0$ for all $t \in [0, T]$ and all $1 \le p < 2$.

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Corollary

Let $\Omega \subset \mathbb{R}^2$ be open, bounded and connected. Let P_0 satisfy (C) with respect to Ω . Then there exists a relaxed Lagrangian solution for SGSW corresponding to the initial data P_0 .

- Renormalization is used to show that the relaxed solutions give rise to dual-space solutions.
- The variational characterization of P, h, α was extended (with stability) from the case $\alpha \ll \mathcal{L}^2$ (Cullen & Gangbo) to $\alpha \in \mathcal{P}_2(\mathbb{R}^2)$: h minimizes $\mathcal{P}^{ac}(\Omega) \ni \rho \mapsto W_2^2(\rho, \alpha) + \|\rho\|_{L^2(\Omega)}^2$ iff $P = h + |\mathrm{Id}_{\Omega}|^2/2$ satisfies (C) and $\alpha = \nabla P_{\#}h$.
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Thank you!

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